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## **EMPIR Project – JRP-n08: Sulf-Norm**

### **WP2 A2.2.12**

#### **Modelling of SO<sub>2</sub> losses in conditioned sampling**

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# 1 Introduction

This report presents the modeling of SO<sub>2</sub> diffusion into water droplet and water film which arise in P-AMs during the process of drying. The motivation is an attempt to quantify the SO<sub>2</sub> losses in dryers during sampling of emissions.

Two basic phenomena need to be taken into account in modeling of removing gaseous substances during drying process. Since the gas is dissolved in liquid water condensing in the dryer, the condensation processes are needed to be described quantitatively in dependence on the physical and geometrical conditions. Once the microscopic liquid phase of water occurs in the dryer (in form of water film or water droplet), the transfer of gaseous soluble species into water is driven by diffusion processes simultaneously with the ongoing condensation of water. Thus, the condensation and diffusion are the key physical phenomena needed to model. However, in more detailed view, the mechanism of dissolution of gaseous substances into liquid consists of four partial processes:

- a) transport of molecules in the gas phase towards the liquid surface (convection, diffusion, Stefan's flow)
- b) transport of the molecules across the gas-liquid interface (mass accommodation effect, Henry's law)
- c) transport of molecules in the liquid phase (convection, diffusion)
- d) chemical reactions in the liquid phase

Since the condensation of water is supposed, the studied situation is moreover complicated by

- e) growth of the water droplet or liquid film formation

All of these phenomena are needed to be implemented in the mathematical model. On top of that, since the processes are running simultaneously, the coupled solution is needed. The mathematical description of each phenomena will be outlined in section 2 and the resulting coupled solution of the comprehensive model will be presented in section 3.

## 2 Theoretical backgrounds

### 2.1 Initial assumptions

For the theoretical findings outlined in the next paragraphs, the single spherical droplet is supposed to be located in mixture of humid air and gaseous SO<sub>2</sub>. For simplicity, the following other assumptions are taken into account for description of condensing droplet:

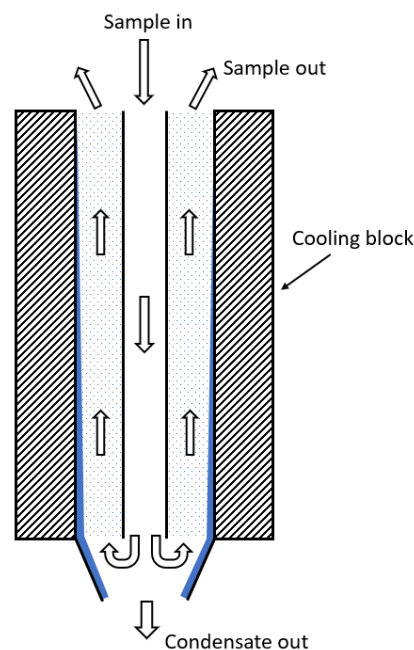
- i. The convection in the droplet and Stefan's flow near the droplet surface due to vapor condensation are not considered.
- ii. No other soluble species than SO<sub>2</sub> are supposed. Note that especially in case of salt water vapor the influence on droplet growth should be assessed.

- iii. The heat transfer is not modeled in gas phase nor in liquid phase. According to [1] under the conditions of slow growth, the droplet temperature is approximately the same as the ambient temperature. In case of more rapid growth empirical relation for droplet temperature based on ambient temperature and relative humidity can be used, see below in paragraph 2.2.
- iv. The molar volume of the  $\text{SO}_2$  dissolved in water is negligible, thus the droplet growth is caused only by condensation of water vapor.
- v. At the droplet surface the equilibrium state of water vapor is supposed, thus the partial water vapor pressure is given by the saturation pressure at the droplet (surface) temperature.
- vi. Since the volume concentration of  $\text{SO}_2$  is low, the latent heat of  $\text{SO}_2$  is negligible compared with the heat released by water vapor condensation.

Analogical assumptions are valid for condensing of liquid film and two others are supposed:

- vii. The liquid film can be considered as planar.
- viii. The laminar convection in the liquid film and the gas phase is taken into account.

On top of assumptions just mentioned, the scheme of the dryer is supposed to be simplified according to Figure 1.



**Figure 1:** Scheme of the condensation in dryer.

In the following text, the theoretical findings for phenomena a) – e) assuming i. – viii. are presented according to cited references.

## 2.2 Droplet growth by condensation

For the description of the droplet growth, the book [1] can serve as a good source text and support the following paragraphs.

The condensation processes in humid air depends on the saturation ratio (multiplied by 100% known as relative humidity) of the partial pressure of water vapor to the saturation vapor pressure

$$S_R = \frac{p}{p_s} \quad (1)$$

The saturation water vapor pressure for a plane liquid surface is given by empirical formula

$$p_s = \exp\left(16.7 - \frac{4060}{T - 37}\right) \quad (2)$$

where  $T$  is the absolute temperature in K and  $p_s$  is the saturation pressure in kPa. In case of the plate surface water begin to condense if  $S_R = 1$ . However, the saturation ratio, called the Kelvin ratio  $K_R$ , in case of microscopic water droplets is given by the Kelvin or Thomson-Gibbs equation

$$K_R = \frac{p_d}{p_s} = \exp\left(\frac{4\gamma M}{\rho R T d}\right) \quad (3)$$

where  $\gamma$ ,  $M$ ,  $\rho$  and  $R$  are the surface tension, molecular weight, density of the liquid and universal gas constant, respectively,  $p_d$  is partial pressure at the surface of the droplet with diameter  $d$ .

Once a stable nucleus of droplet is established, i.e. the diameter of nucleus is bigger than  $d$  given by formula (3) for a given saturation ratio (or reversely saturation ratio is bigger than  $K_R$  given by formula (3) for a given nucleus diameter), the droplet begins to growth. In [1] two formula in dependence on the current diameter are derived. When the droplet diameter is less than gas mean free path,  $\lambda$ , the formula has the form

$$\frac{d(d_p)}{dt} = \frac{2M\alpha(p_\infty - p_d)}{\rho N_A \sqrt{2\pi m k T}} \quad \text{for } d_p < \lambda \quad (4)$$

where  $\alpha$  is accommodation coefficient, see the section 2.4,  $p_\infty$  is partial pressure far from the droplet surface,  $p_d$  is partial pressure of vapor at the droplet surface given by Kelvin equation (3),  $N_A$  is Avogadro constant,  $m = M/N_A$  is mass of vapor molecule and  $k = 1.3806485 \cdot 10^{-23}$  is Boltzmann constant.

In case the droplet is bigger than mean free path, the formula becomes

$$\frac{d(d_p)}{dt} = \frac{4DM}{R\rho d_p} \left(\frac{p_\infty}{T_\infty} - \frac{p_d}{T_d}\right) \phi \quad \text{for } d_p > \lambda \quad (5)$$

where  $D$  is diffusion coefficient of water vapor,  $\phi$  is the Fuchs correction factor significant for particles less than  $1\mu\text{m}$ ,  $p_d$  is vapor partial pressure near the droplet surface which according to assumption v. can be calculated by formula (2) and finally the droplet temperature  $T_d$  can be computed using ambient temperature  $T_\infty$  by

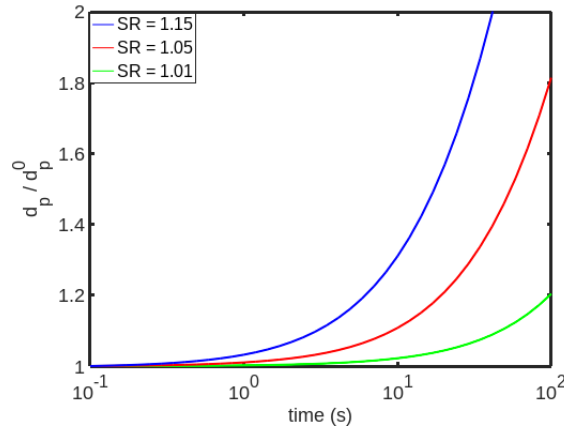
$$T_d = T_\infty + \frac{(6.65 + 0.345T_\infty + 0.0031T_\infty^2)(S_R - 1)}{1 + (0.082 + 0.00782T_\infty)SR} \quad (6)$$

where  $T_\infty$  in ratio needs to be in °C.

Since gas mean free path is quite small in comparison to final droplet diameters, the amount of dissolved  $SO_2$  is not affected by the initial droplet growth described by (4). Moreover, since the expected final droplet diameters are much bigger than 1  $\mu m$  (due to filter resolution), the droplet growth can be approximated by equation (5) for  $\phi = 1$ . After integration of (5) we get the prescription for droplet diameter in time by

$$d_p(t) = \sqrt{\frac{8DM}{R\rho} \left( \frac{p_\infty}{T_\infty} - \frac{p_d}{T_d} \right) (t - t^0) + (d_p^0)^2} \quad (7)$$

where  $d_p^0$  is initial droplet diameter in time  $t^0$ .



**Figure 2:** Temporal evolution of the droplet diameter in dependence on the saturation ratio.

### 2.3 Liquid film condensation

The classical analysis of laminar film on inclined or vertical wall is known from Nusselt (1916). Since the film covers whole surface, the condensation process is pushed by the heat transfer between the vapor and wall. Following from the Fourier's Law and supposing some assumptions (laminar flow, stagnant liquid vapor, smooth liquid film surface, etc. see [2]), the following formula (8) for liquid film thickness,  $\delta$ , can be derived.

$$\delta(x) = \left( \frac{4k_l\mu_l x \Delta T}{\rho_l(\rho_l - \rho_v)gh_{lv}} \right)^{1/4} \quad (8)$$

where  $k_l$  is thermal conductivity of the liquid,  $\mu_l$  and  $\rho_l$  are dynamic liquid viscosity and liquid density,  $\rho_v$  is vapor density,  $x$  is coordinate in the wall direction and  $h_{lv}$  is latent heat of the liquid-vapor phase change.

As follows from the design of P-AMs, the film is supposed to flow downward due to gravity, while stack gas is flowing upward through the cooler. We speak about countercurrent vapor flow and the motion of gas needs to be taken into account. Authors in [2] show that a nonlinear system of govern-

ing equations is derived if a shear stress from vapor motion is supposed on the liquid film interface. The corresponding mechanistic models are outlined in [3], [4] or [5]. Numerical implementation using finite volume methods can be found in [6] or [7]. However, since the models of countercurrent are rather complex, the film thickness given by Nusselt theory is taken as the acceptable estimate.

## 2.4 Mass transfer through gas-liquid interface

### Henry's law

The basic form of Henry's law says that the concentration of a species in the aqueous phase,  $n_l$ , is proportional to the partial pressure of this matter in the gas phase,  $p$ , according to formula

$$n_l = H^{cp} p \quad (9)$$

where  $H^{cp}$  is the Henry's law constant in  $\text{mol}\cdot\text{m}^{-3}\cdot\text{Pa}^{-1}$ . Note that there are several other definitions of Henry's law constant. The dimensionless form is given by

$$H^{cc} = \frac{n_l}{n_g} \quad (10)$$

The conversion between  $H^{cp}$  and  $H^{cc}$  for ideal gas is

$$H^{cc} = H^{cp} R T \quad (11)$$

where  $R$  is the universal gas constant.

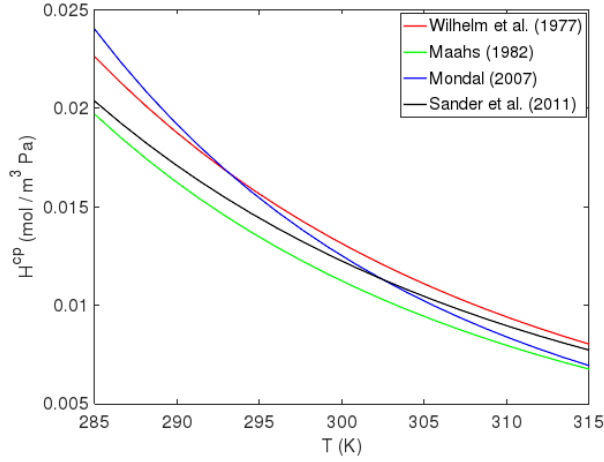
The dependence of Henry's law constant on temperature can be extrapolated from a single data point applying van't Hoff equation by formula

$$H(T) = H^0 \exp\left[-\frac{\Delta h}{R}\left(\frac{1}{T} - \frac{1}{T^0}\right)\right] \quad (12)$$

where  $H^0$  is Henry's law constant in temperature  $T^0$  and  $\Delta h$  is enthalpy change due to transport of soluble gas substance into liquid. The dependence of Henry's law constant on temperature and independence on pressure have been confirmed in [8]. The values of Henry's law constant for different substances and water as solvent can be found in [9]. Selected values of  $H^0$  and  $-\Delta h/R$  can be found in Table 1 for  $T^0=298.15$  K. The corresponding Henry's constants in dependence on temperature are depicted in Figure 3.

**Table 1:** Parameters for determination of Henry's Law constant according to formula (8).

$H^{cp}$ at $T^0$ ( $\text{mol m}^{-3} \text{Pa}^{-1}$ )	$\frac{d \ln H^{cp}}{d(1/T)}$ (K)	Reference
$1.4 \cdot 10^{-2}$	2800	[10]
$1.2 \cdot 10^{-2}$	3200	[11]
$1.3512 \cdot 10^{-2}$	3715.2	[12]
$1.3 \cdot 10^{-2}$	2900	[13]



**Figure 3:** Henry's Law constant vs. temperature.

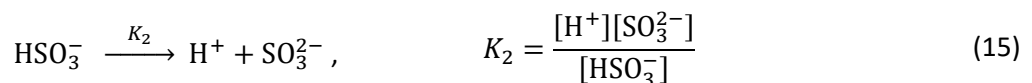
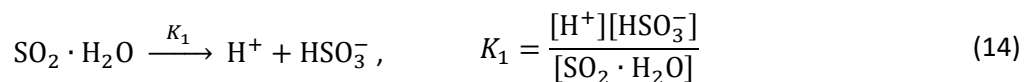
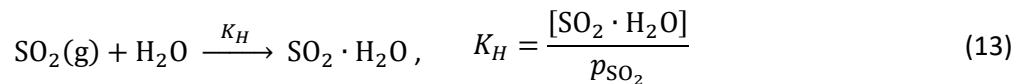
### Mass accommodation coefficient (sticking coefficient)

The mass accommodation coefficient,  $\alpha$ , is defined as the ratio of molecules absorbed through the gas-liquid interface to the number of molecules which hit the liquid surface. As has been shown in [14] using the comparison of characteristic times of different processes, the droplet surface is saturated faster than equilibrium state can be established due to diffusion. This leads to the re-evaporation of some molecules from interface and the resulting ratio of absorbed molecules can be described using the mass accommodation coefficient. The experimental measurements of mass accommodation coefficient of  $\text{SO}_2$  at the air-water interface are presented in [14] and [15]. The authors report measurements of  $\alpha = (6.0 \pm 0.8) \cdot 10^{-2}$  at 298 K and  $\alpha = (5.4 \pm 0.6) \cdot 10^{-2}$  at 295 K, respectively.

## 2.5 Physical and chemical processes in liquid phase

### Chemical reactions

Experimental observations in [14] show that the mass accommodation coefficient is strongly dependent on pH and real solubility at the droplet surface. For considering of these dependencies, the chemical processes in droplet need to be take into account. According to [16], the following chemical reactions occur if  $\text{SO}_2$  is dissolved in water



where  $[\cdot]$  represents the concentration,  $K_H$  is the Henry's constant  $H^{cp}$ ,  $K_1$  and  $K_2$  are the first and the second dissociation constants. The total concentration of the dissolved sulfur with oxidation number four can be calculated as the sum

$$[S(IV)] = [SO_2 \cdot H_2O] + [HSO_3^-] + [SO_3^{2-}] \quad (16)$$

Expressing (16) using dissociation constants defined in (13-15), the total concentration has form

$$[S(IV)] = K_H p_{SO_2} \left( 1 + \frac{K_1}{[H^+]} + \frac{K_1 K_2}{[H^+]^2} \right) \quad (17)$$

### Effective Henry's constant

Following the considerations in previous paragraph, the equilibrium state of the sulfur dioxide concentration inside the droplet can be expressed by

$$n_l = H_{S(IV)}^{cp,*} p_{SO_2}, \quad (18)$$

where  $H_{SO_2}^{cp,*}$  is the effective Henry's constant. Using (17) the effective Henry's constant has the form

$$H_{S(IV)}^{cp,*} = K_H \left( 1 + \frac{K_1}{[H^+]} + \frac{K_1 K_2}{[H^+]^2} \right), \quad (19)$$

where concentration  $[H^+]$  is connected with pH scale by

$$pH = -\log[H^+], \quad (20)$$

where  $[H^+]$  need to be in mol/dm<sup>3</sup>.

Note that neglecting the second dissociation term  $[SO_3^{2-}]$  which has very low concentration and using equality  $[H^+] = [HSO_3^-]$ , see [12], the effective Henry's law constant can be expressed by

$$H_{S(IV)}^{cp,*} = K_H + \sqrt{\frac{K_H K_1}{p_{SO_2}}} \quad (21)$$

As well as the Henry's law constant also the dissociation constants are dependent on temperature. According to experimental data between 0 and 50°C author in [11] proposed correlations

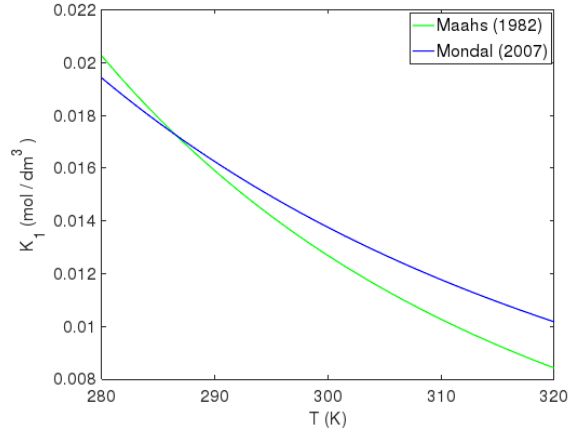
$$\log K_1 = \frac{853}{T} - 4.74 \quad (22)$$

$$\log K_2 = \frac{621.9}{T} - 9.278 \quad (23)$$

Measurements based formula for  $K_1$  in mol/dm<sup>3</sup> presents [12] using

$$\ln K_1 = \frac{1447.1}{T} - 9.11 \quad (24)$$



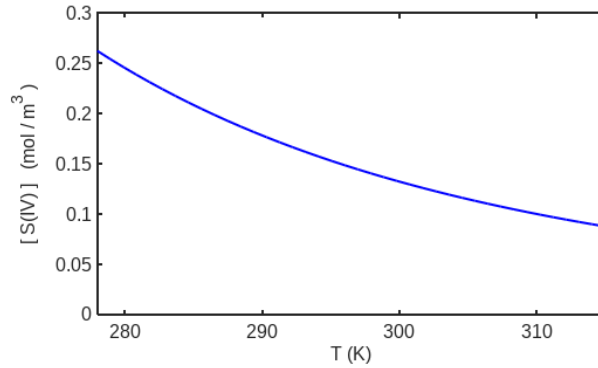


**Figure 4:** First dissociation constant  $K_1$  vs. temperature.

Using equations (18) and (21), prescription for Henry's constant  $K_H$  from [12], see Table 1, and prescription for  $K_1$  by formula (24) authors in [12] derived equation for total concentration of sulfur dioxide absorbed in water in dependence on temperature

$$[S(IV)] = 2.407 \cdot 10^{-6} \cdot \left[ 0.0218 \cdot \exp\left(\frac{3715.2}{T}\right) p + \exp\left(\frac{2581.1}{T}\right) p^{0.5} \right] \quad (25)$$

where  $[S(IV)]$  is in  $\text{mol/dm}^3$ ,  $T$  in K and  $p$  in kPa. The corresponding concentration of  $\text{SO}_2$  which can be absorbed from air with initial concentration of 0.01 ppm in dependence on temperature can be seen in Figure 5.



**Figure 5:** Total  $\text{SO}_2$  absorbed in water in dependence on temperature. Initial  $\text{SO}_2$  concentration in air 0.01ppm.

## 2.6 Diffusion

### Governing equations

The diffusion processes are described by equation

$$\frac{\partial n(x, t)}{\partial t} = \frac{\partial}{\partial x} \left( D \frac{\partial n(x, t)}{\partial x} \right) \quad (26)$$

where  $n(x, t)$  is concentration in position  $x$  and time  $t$  and  $D$  is the diffusion coefficient. As was already mentioned in introduction, the studied phenomenon need to be solved as coupled system

for both phases, therefore the system of two coupled diffusion equation for constant diffusion coefficients takes the form

$$\frac{\partial n_g(x, t)}{\partial t} = D_g \frac{\partial^2 n_g(x, t)}{\partial x^2} \quad (27)$$

$$\frac{\partial n_l(x, t)}{\partial t} = D_l \frac{\partial^2 n_l(x, t)}{\partial x^2} \quad (28)$$

The dependence of diffusion coefficient of SO<sub>2</sub> in water on temperature and pH is presented in [17]. The influence on temperature is given by

$$D_l = -1.21 \cdot 1010^{-3} + 4.33 \cdot 10^{-6} \cdot T \quad (29)$$

Initial conditions define the concentrations inside and outside of the droplet:

$$n_l(x, 0) = n_{l,0} \quad \text{for } x < d_p/2 \quad (30)$$

$$n_g(x, 0) = n_{g,0} \quad \text{for } x > d_p/2 \quad (31)$$

Since the concentration far from the surface is supposed to be constant, the mass flux due to diffusion vanishes. Similarly, the mass flux can be assumed to be zero at the centre of droplet due to symmetry of the process. The corresponding boundary conditions can be defined as

$$-D_l \frac{\partial}{\partial x} n_l(0, t) = 0 \quad (32)$$

$$-D_g \frac{\partial}{\partial x} n_g(\infty, t) = 0 \quad (33)$$

Analytical solutions of equations (27) and (28) for different geometrical configurations and different boundary conditions are given in [18]. However, the presented suitable solutions do not take into account the Henry's law which limits the maximal concentration of dissolved substance in liquid and the mass accommodation coefficient which defines the fraction of molecules entering through the gas-liquid interface.

### Boundary conditions at the interface

The boundary conditions which take into account the Henry's Law and accommodation coefficient are presented in [19] by forms

$$-D_g \frac{\partial}{\partial x} n_g(d_p/2, t) = \frac{\alpha \bar{v}}{4} \left( n_g(d_p/2, t) - \frac{n_l(d_p/2, t)}{H^{cc}} \right), \quad (34)$$

$$-D_l \frac{\partial}{\partial x} n_l(d_p/2, t) = \frac{\alpha \bar{v}}{4} \left( n_g(d_p/2, t) - \frac{n_l(d_p/2, t)}{H^{cc}} \right), \quad (35)$$

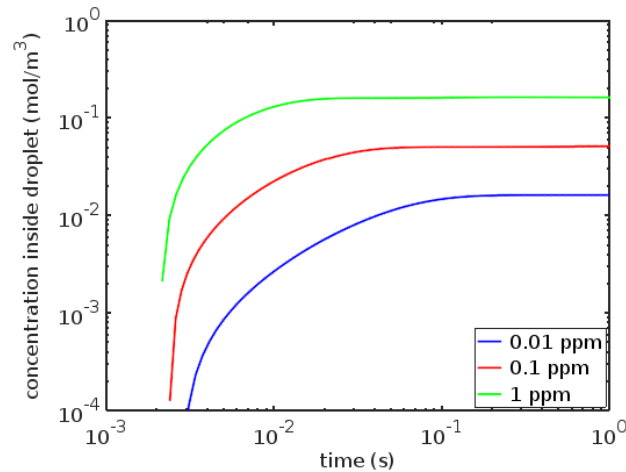
where the gas-liquid interface is located at  $x=d_p/2$  and  $\bar{v}=300$  m/s is mean thermal velocity of SO<sub>2</sub> as given by [20].

### 3 Numerical model and results

Fairly comprehensive attitude to the modeling of soluble gas transport into large droplets during evaporation and condensation is presented in [16] and [21]. Rather complex numerical solution of the presented partial differential equations with moving boundaries led to the idea of using the tool based on finite element (FEM) or finite volume method (FVM).

Regarding to the geometry and physical configuration of the supposed dryer, the estimate of maximal concentrations can be done separately using model of diffusion into the growing droplet and into the condensing film.

In case of droplet, simultaneous simulation of diffusion and droplet growth need to be done. Resulting concentrations inside droplet computed using FEM tool are presented in Figures 6 in dependence on time and initial  $\text{SO}_2$  concentrations. The model is based on equations (27-28) with initial conditions for gas and liquid (32-33) and boundary conditions (34-35) prescribed on gas-liquid interface. The mass accommodation coefficient is set to 0.054, the effective Henry's law constant is computed using (21) where dissociation constants  $K_H$  and  $K_1$  are taken from [12], see Table 1 and formula (24). The droplet growth is defined by formula (7) where droplet temperature is prescribed by formula (6), vapor partial pressure is given by equation (3) and saturation pressure is given by (2).



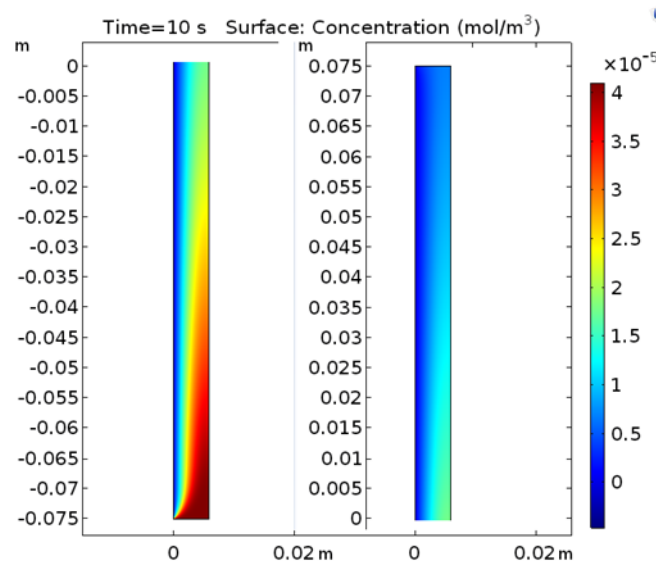
**Figure 6:** Concentration of  $\text{SO}_2$  in growing droplet during time for different values of  $\text{SO}_2$  concentration in air with relative humidity equal to 101% and temperature  $20^\circ\text{C}$ .

The Figure 6 shows the concentration of  $\text{SO}_2$  in growing droplet in dependence on time. Depending on the initial concentration of  $\text{SO}_2$  in gas, the time needed to reach the maximal concentration, which is given by formula (25), ranges between 0.02 and 0.2s. Regarding the dimensions of condensation tubes in P-AMSS (diameter of the tube from 6mm, height from 125 mm) and maximal flow velocity of sampled air (approx. 0.4 m/s), we can suppose, that the maximal concentration can be reached before the droplet is attached to the liquid film on the wall. On the other hand, the droplets are established only in case the relative humidity is bigger than 100%. In other cases, the condensation on the cooled wall with temperature  $5^\circ\text{C}$  is predominant for the drying processes in the dryer. Hence, the modeling of liquid film formation due to condensation and corresponding  $\text{SO}_2$  dissolution is important for total amount of dissolved  $\text{SO}_2$ .

In contrast to the growing droplet, in case of the condensing liquid film, the corresponding analysis can be done for a given film thickness, if we can suppose that in process of measurement a stagnant liquid film is established. However, although the stationary liquid film thickness can be taken into account, note that the mathematical models suffer from other complications such as the modeling of interfacial friction or liquid velocity on the film surface in case of moving vapor in dryer when the formula (8) cannot be used.

In case of a typical P-AMS, the maximum gas volume flow rate is 150 NI/h. For a tube with diameter 6mm, this corresponds to gas velocity of about 0.4m/s and laminar flow regime. Supposing limited influence of this laminar flow on the liquid film, the classical Nusselt formula (8) can serve as a basic estimate of the liquid film thickness along the dryer and the liquid volume in which the gas can diffuse. Supposing the length of the dryer tube 0.15m with diameter 0.006m and standard physical properties of water, the film thickness using (8) equals about 0.1 mm for saturation temperature of about 290K. As in the case of water droplet, several FEM simulations have been done using the theoretical basis outlined in the previous chapter. The main difference was the stagnant interface between liquid and gas phase and resolution of convection using the assumptions of laminar flow regimes in both phases. The constant film thickness of 0.1mm has been defined along the 0.15m long dryer. Zero flux of SO<sub>2</sub> has been prescribed in the upper boundary of liquid film and on the wall. The boundary conditions (34-35) have been prescribed on the gas-liquid interface.

The resulting concentration field of SO<sub>2</sub> in stack gas inside the dryer is depicted in Figure 7.



**Figure 7:** Concentration of SO<sub>2</sub> in gas inside dryer in time 10s for initial SO<sub>2</sub> concentration 1ppm and gas velocity 0.1 m/s.

The resulting ratios of outlet to inlet SO<sub>2</sub> concentrations in gas for different average velocity of gas flow and different inlet concentrations can be read in Table 2. The resulting dependencies show that SO<sub>2</sub> losses increase with decreasing gas velocity and decreasing SO<sub>2</sub> concentration in gas phase. The losses increase only slightly for lower inlet concentrations, while dependence on the gas velocity is much more significant.

**Table 2:** SO<sub>2</sub> concentration after passing through the cooler with constant water film thickness. Resulting values are normalized to the initial concentration before entrance to the cooler.

	$n_g^{SO_2} = 1$ ppm	$n_g^{SO_2} = 10$ ppm	$n_g^{SO_2} = 100$ ppm
$U_g = 0.4$ m/s	46.9 %	48.4 %	52.4 %
$U_g = 0.2$ m/s	28.6 %	29.9 %	33.8 %
$U_g = 0.1$ m/s	12.6 %	13.4 %	16.0 %

## 4 Conclusions

The presented mathematical models of droplet growing and liquid film formation during condensation have been outlined on basis of theoretical as well as experimental findings published in literature. The molecular transport of condensable gas is modeled with regards to the molecular processes on gas-liquid interface as well as with regards to chemical processes inside liquid. The governing equations of diffusion processes are solved using boundary condition on gas-liquid interface which is designed using physical parameters mentioned in theoretical sections.

The concentrations of SO<sub>2</sub> dissolved in one droplet needs to be computed using presented mathematical model of diffusion simultaneously with modeling of droplet growth. From the comparison of droplet growth time and droplet lifetime (time needed for leaving of the droplet from the dryer) it follows that ratio of final and initial droplet diameter is less than two. Thus, the dissolution of SO<sub>2</sub> into droplets should be less significant than dissolution into water film and therefore in following computations, summarized in Table 2, only liquid film condensation has been taken into account.

The average concentration of SO<sub>2</sub> in gas flow leaving the dryer has been simulated for stagnant film thickness, but as well as for droplet using aforementioned diffusion model based on Henry's law theory and chemical reactions connected with dissolution of SO<sub>2</sub> in water. From the comparison of resulting SO<sub>2</sub> concentrations at the inlet to the dryer and the final concentrations at the outlet from the dryer it follows that the SO<sub>2</sub> losses can be significant for the measurements using P-AMSS.

### 4.1 Summary of simplifications and challenges for future

In addition to basic assumptions outlined in section 2.1, following simplifications and notes need to be taken into account for assessing of model relevance in qualitative and quantitative sense:

- Only 2D simplified geometry, see Figure 1, with given dimensions (diameter D=6mm, height L=15cm) has been used for computations of SO<sub>2</sub> diffusion into water film condensed in dryer.
- The water film thickness has been set to constant value  $t=0.1$ mm according to formula (8).
- Constant temperature of stack gas,  $T_\infty$ , and water film,  $T_w$ , has been set ( $T_\infty=20^\circ\text{C}$ ,  $T_w=5^\circ\text{C}$ ).
- Relative humidity has been set to 101%.

- Better modeling of the film thickness for countercurrent vapor flow is needed.
- Initial gas properties (temperature, humidity) should be set according to real values.
- Adjustment of parameters and validation using real process or experiments is needed.

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