

Ion-atom physics

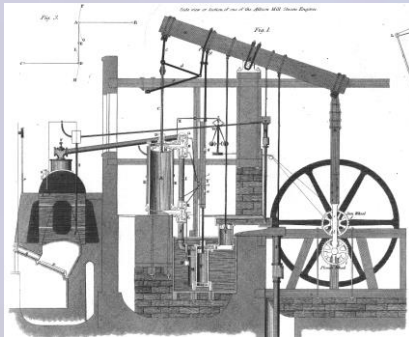
Carlo Sias

Istituto Nazionale di Ricerca Metrologica - INRIM
European Laboratory for Nonlinear Spectroscopy - LENS

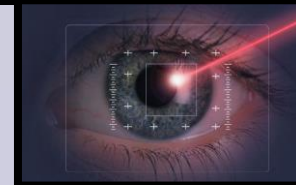
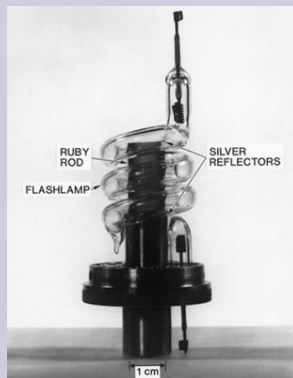
Gressoney, 13 September 2018

Understanding physics enables new technology

Thermodynamics



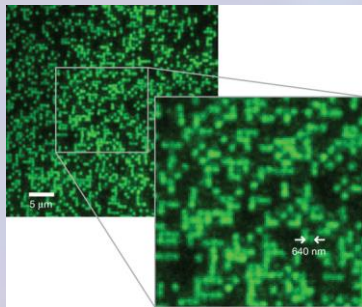
Optics, spectroscopy



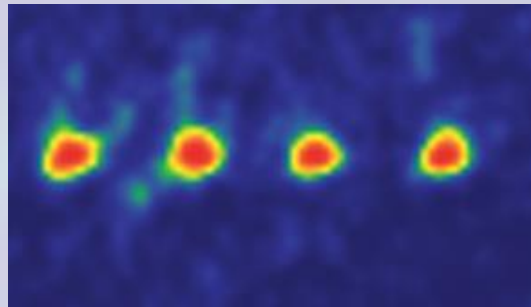
..and future

Where are we with quantum mechanics?

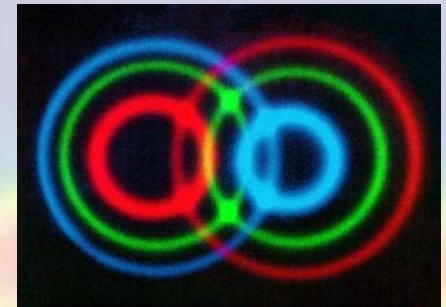
Neutral atoms



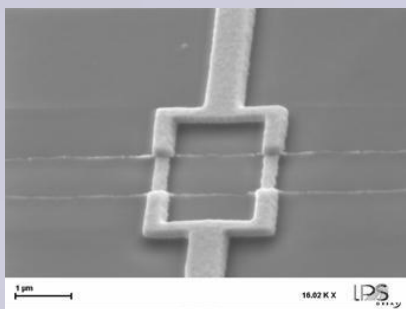
Ions



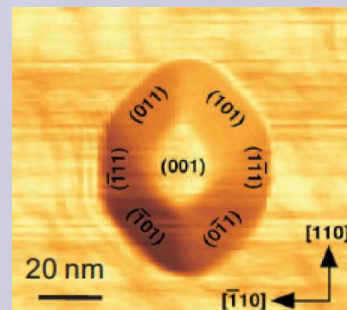
Photons



Squids



Quantum dots



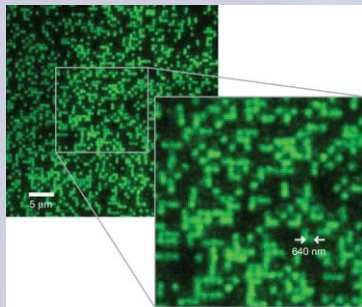
... and many others

..at a good point!

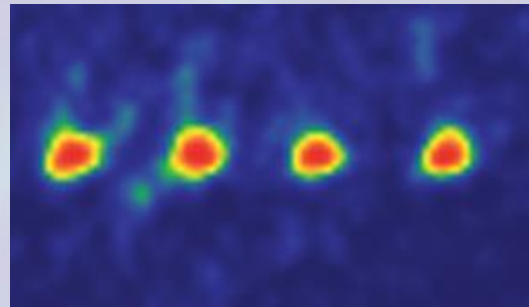
..and future

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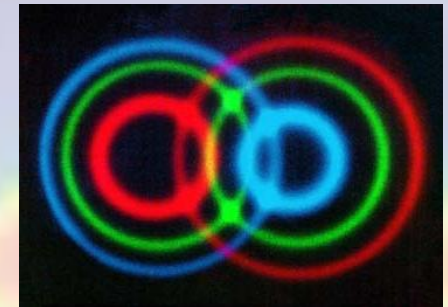
Neutral atoms



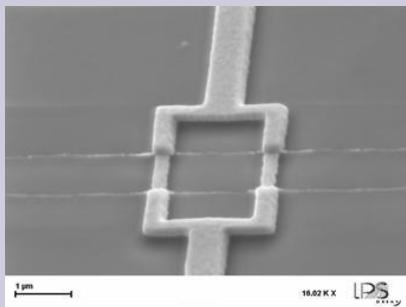
Ions



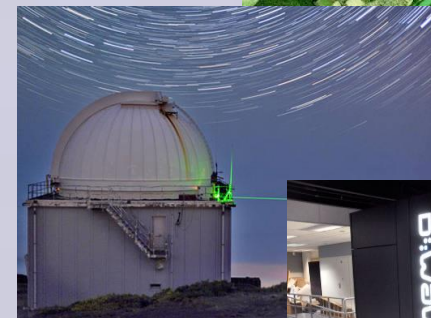
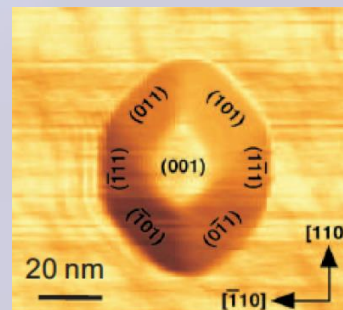
Photons



Squids



Quantum dots

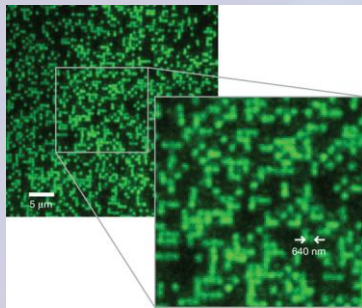


..at a good point!

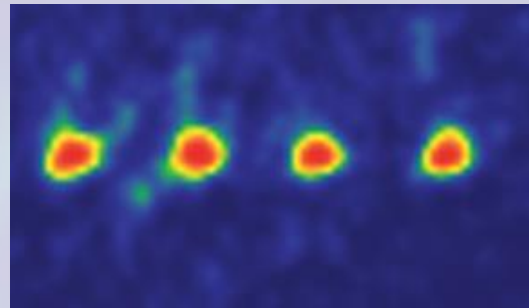
..and future

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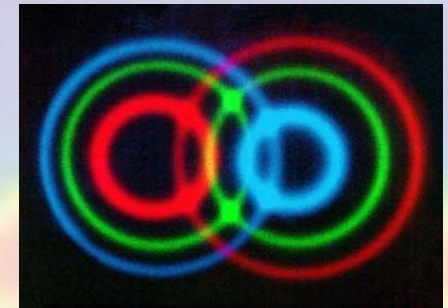
Neutral atoms



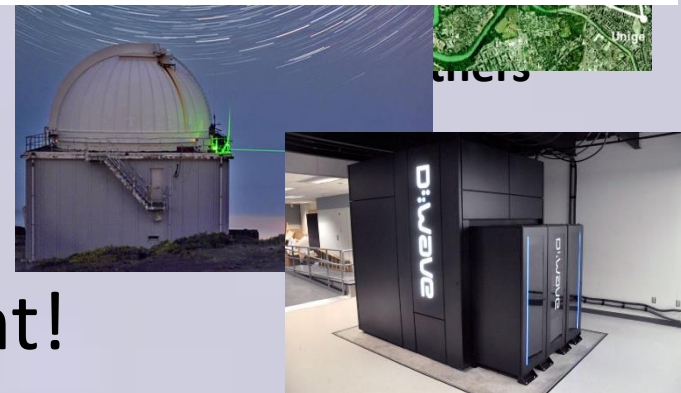
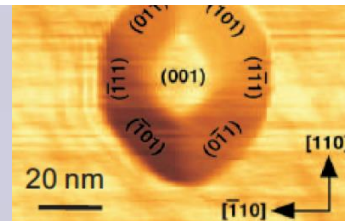
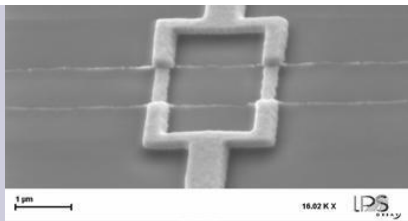
Ions



Photons



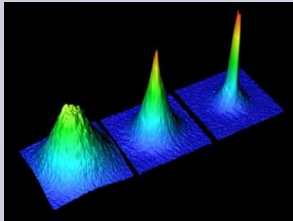
Can we think of a composite system formed by two **coupled** quantum systems?



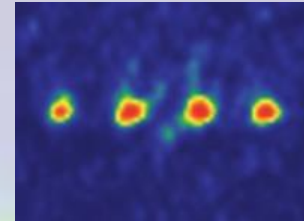
..at a good point!

Ultracold atoms - overview

Same fields, different approaches



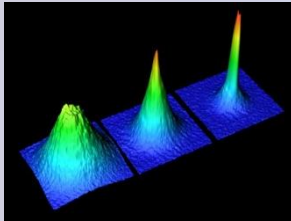
Coherent matter



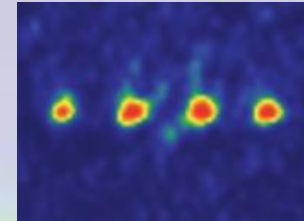
Single particles

Ultracold atoms - overview

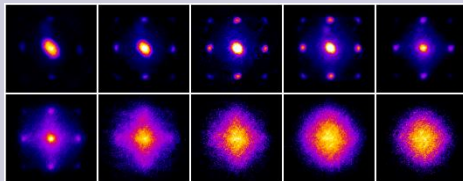
Same fields, different approaches



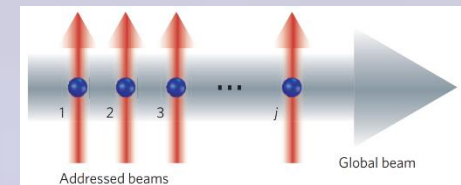
Coherent matter



Single particles



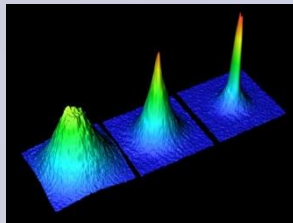
Many-body physics



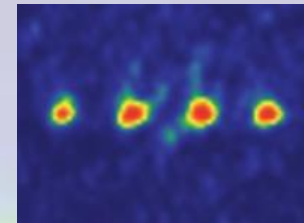
Blatt, Roos Nat. Phys. 2012

Ultracold atoms - overview

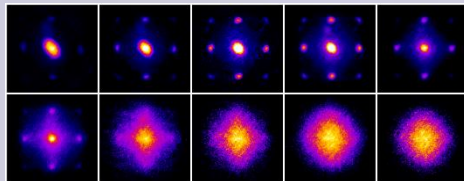
Same fields, different approaches



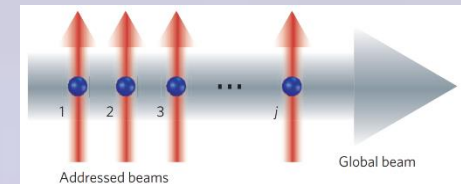
Coherent matter



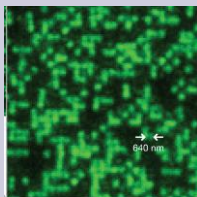
Single particles



Many-body physics

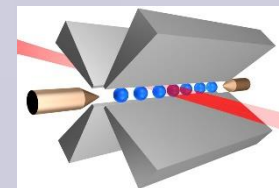


Blatt, Roos Nat. Phys. 2012



Greiner group

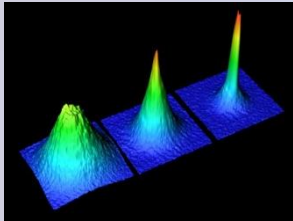
Single particle manipulation



Blatt group

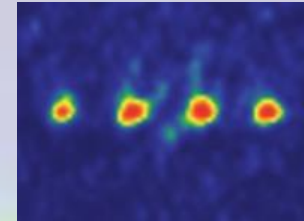
Ultracold atoms - overview

Same fields, different approaches

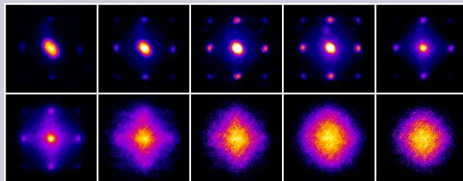


Coherent matter

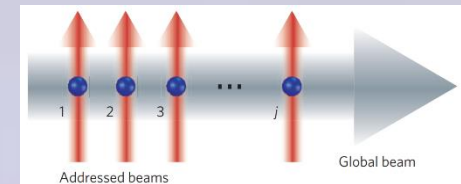
CLOCKS!



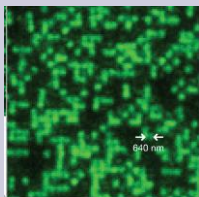
Single particles



Many-body physics

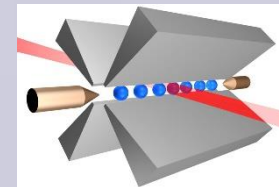


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Greiner group

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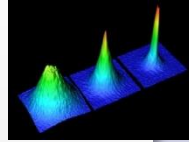


Blatt group

Hybrid quantum systems

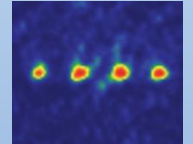
Quantum degenerate atoms

- Macroscopic quantum system at nano Kelvin temperature
- Collective quantum states of 10^6 particles
- Very long coherence times



Ultracold trapped ions

- Single particle detection and manipulation
- Pristine source for quantum information processing & precision spectroscopy



+ Interactions

Hybrid quantum system

A new approach to

- **Ultracold collisions & quantum chemistry**
- **Quantum information processing and decoherence**
- **Quantum many-body physics with impurities**
- **Metrology**

Ion-neutral interactions

An old problem in physics..

Ion-neutral collisions
(1905)

**UNE FORMULE FONDAMENTALE DE THÉORIE CINÉTIQUE;
PAR M. P. LANGEVIN.**

Ion-neutral interactions

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Ion-neutral collisions
(1905)

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Scattering of Ions by Polarization Forces

ERICH VOGT* AND GREGORY H. WANNIER
Bell Telephone Laboratories, Murray Hill, New Jersey
(Received May 19, 1954)

Scattering from neutral-
ion potential

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Scattering from neutral-
ion potential

First atom-ion
experiment (He^+ -Cs)

Exchange-Collision Technique for the rf Spectroscopy of
Stored Ions*†

F. G. MAJOR‡ AND H. G. DEHMELT
Department of Physics, University of Washington, Seattle, Washington 98105
(Received 6 September 1967)

Buffer gas sympathetic cooling

First application: sympathetic cooling of charged particles

Increase of confinement
lifetime of Hg^+ by buffer
gas cooling (Ne)

The Three-Dimensional Quadrupole Ion Trap

P. H. DAWSON and N. R. WHETTEN

Received December 16, 1968

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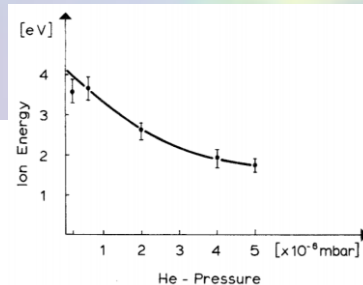
P. H. DAWSON and N. R. WHETTEN

Received December 16, 1968

Appl. Phys. 25, 249–251 (1981)

Trapped Ion Density Distribution in the Presence of He-Buffer Gas

H. Schaaf, U. Schmeling, and G. Werth



Cooling of Ba^+ by the presence of He buffer gas

+ theoretical analysis

Computed energy and spatial statistical properties of stored ions cooled by a buffer gas

F. Vedel, J. André, M. Vedel, and G. Brincourt

Brownian Motion of a Parametric Oscillator: A Model for Ion Confinement in Radio Frequency Traps

R. Blatt¹, P. Zoller², G. Holzmueller², and I. Siemers¹

Ion-neutral quantum mixtures

Can we exploit our knowledge in trapping and cooling particles to study collisions at lower temperature?



Ion-neutral quantum mixtures

Can we exploit our knowledge in trapping and cooling particles to study collisions at lower temperature?

Yes!

PRL 102, 223201 (2009)

PHYSICAL REVIEW LETTERS

Observation of Cold Collisions between Trapped Ions and Trapped Atoms

Andrew T. Grier, Marko Cetina, Fedja Oručević, and Vladan Vuletić

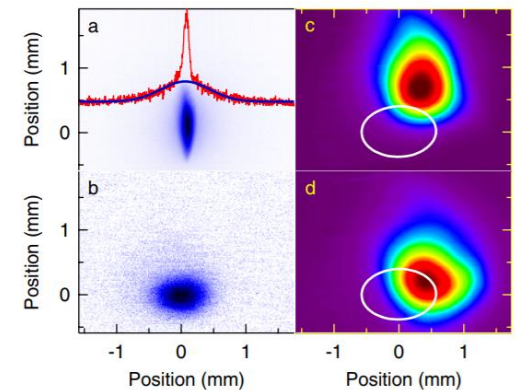


FIG. 3 (color online). (a) (1,0,1) camera image of the ion crystal (blue or dark gray) and cross-section showing highly non-Gaussian shape of crystal (red or gray). (b) (0,1,1) camera image of the ion crystal. (c) Typical low-overlap setting between MOT (colored or shaded contours) and $1/e^2$ contour of ions (white). (d) Same as (c) but for a higher overlap setting.

Ion-neutral quantum mixtures

Can we exploit our knowledge in trapping and cooling particles to study collisions at lower temperature?

Yes!

A fast growing community

Yb-Yb ⁺	Rb-Rb ⁺	Rb-N ₂ ⁺
Rb-Yb ⁺	Rb-Ca ⁺	Li-Ca ⁺
Li-Yb ⁺	Ca-Yb ⁺	Ca-BaCl ⁺
Rb-Ba ⁺	Rb-Sr ⁺	Rb-K ⁺
Yb-Yb ⁺	Na-Na ⁺

PRL 102, 223201 (2009)

PHYSICAL REVIEW LETTERS

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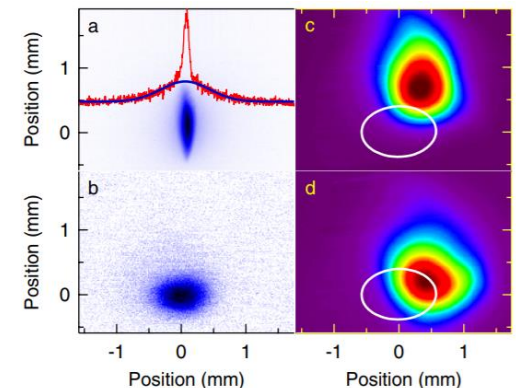


FIG. 3 (color online). (a) (1,0,1) camera image of the ion crystal (blue or dark gray) and cross-section showing highly non-Gaussian shape of crystal (red or gray). (b) (0,1,1) camera image of the ion crystal. (c) Typical low-overlap setting between MOT (colored or shaded contours) and $1/e^2$ contour of ions (white). (d) Same as (c) but for a higher overlap setting.

Outline

A brief intro of atom-ion collisions

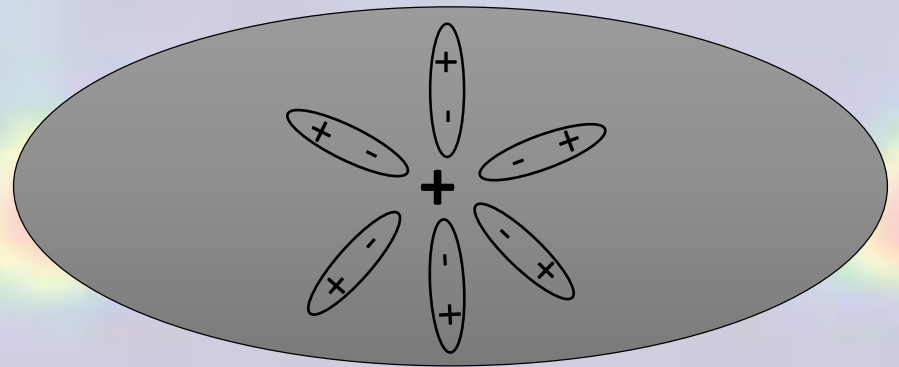
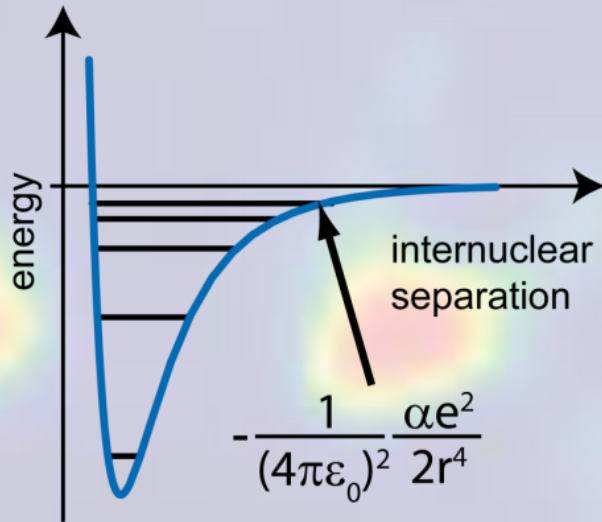
Hybrid atom-ion experiments – state-of-the-art

Hybrid atom-ion experiments: perspectives

Atom-ion systems for improving ion clocks

Atom-ion interactions

Charge-induced dipole interaction: R^{-4} scaling



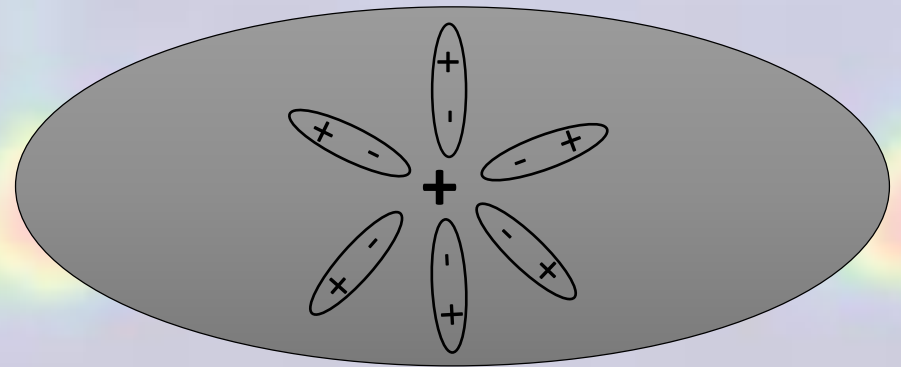
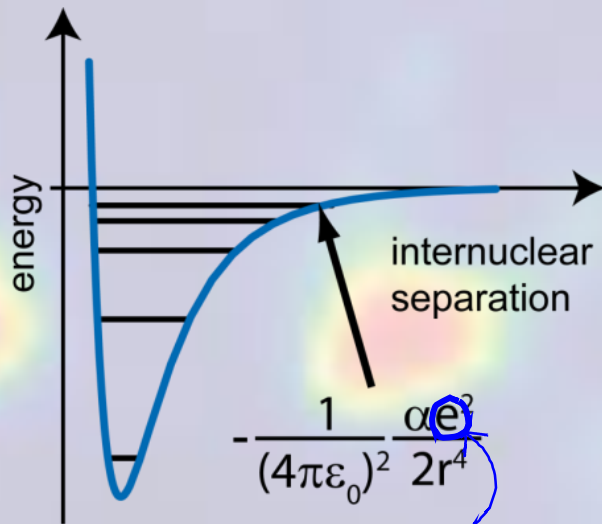
Longer ranged than atom-atom interactions!

$$\underline{R^* \gtrsim 100\text{nm}}$$

For atom-atom interactions
typically $R^* \approx 1\text{ nm}$

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Charge-induced dipole interaction: R^{-4} scaling



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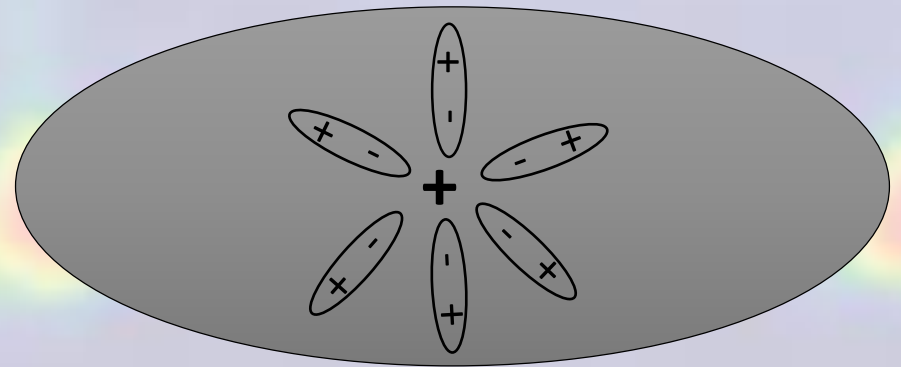
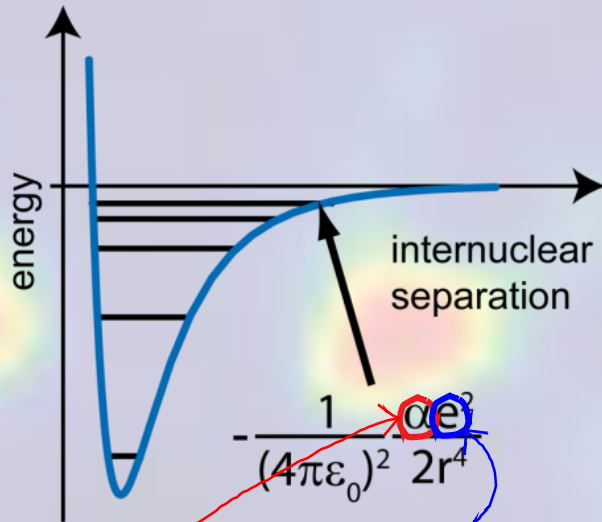
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Depends only from:

- The ion's charge

Atom-ion interactions

Charge-induced dipole interaction: R^{-4} scaling



Longer ranged than atom-atom interactions!

$$\underline{R^* \gtrsim 100\text{nm}}$$

For atom-atom interactions
typically $R^* \approx 1\text{nm}$

Depends only from:

- The ion's charge
- The DC polarizability of the neutral particle

→ independent from the ion's element!

Atom-ion collisions

First theory by Langevin (classical)

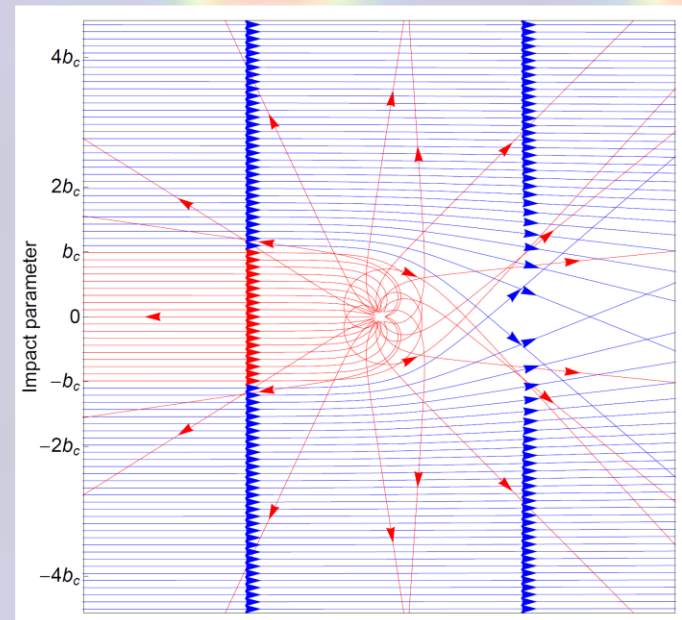
$$V(r) = -\frac{C_4}{2r^4} \quad C_4 = \frac{\alpha e^2}{(4\pi\epsilon_0)^2}$$

We can define a *critical impact parameter*

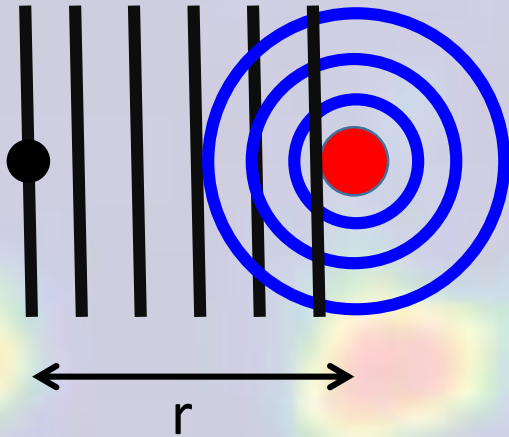
$$b_c = \left(\frac{2C_4}{E_c} \right)^{1/4}$$

Large deflection $b < b_c$

Small deflection $b > b_c$



..a bit of theory



Colliding particles – Schroedinger eq.

$$\left(-\frac{\hbar^2}{2\mu} \nabla^2 + V(\mathbf{r}) \right) \Psi(\mathbf{r}) = \frac{\hbar^2 k^2}{2\mu} \Psi(\mathbf{r})$$

Solution (at large r):
In: plane wave
Out: spherical wave

$$\Psi(\mathbf{r}) \sim e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$

..a bit of theory

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Hypothesis: $V(r)$ is central \rightarrow cylindrical simmetry
Ok in atom-atom, atom-ion NOT in dipole-dipole!

\rightarrow Noether theorem, a quantity is conserved: angular momentum

..a bit of theory

Expansion in partial waves

$$\Psi(\mathbf{r}) \sim e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$



..a bit of theory

Expansion in partial waves

$$\Psi(\mathbf{r}) \sim e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$

(from expansion of spherical
Bessel function)

$$e^{ikz} \sim \frac{1}{2ikr} \sum_{l=0}^{\infty} (2l+1) (e^{ikr} - (-1)^l e^{-ikr}) P_l(\cos \theta)$$

..a bit of theory

Expansion in partial waves

$$\Psi(\mathbf{r}) \sim e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$
$$f(\theta) = \sum_{l=0}^{\infty} (2l + 1) f_l P_l(\cos \theta)$$

$$e^{ikz} \sim \frac{1}{2ikr} \sum_{l=0}^{\infty} (2l + 1) (e^{ikr} - (-1)^l e^{-ikr}) P_l(\cos \theta)$$

..a bit of theory

Expansion in partial waves

$$\Psi(\mathbf{r}) \sim e^{ikz} + f(\theta) \frac{e^{ikr}}{r}$$

Here is the physics

$$f(\theta) = \sum_{l=0}^{\infty} (2l+1) f_l P_l(\cos \theta)$$

Legendre polynomials

$$e^{ikz} \sim \frac{1}{2ikr} \sum_{l=0}^{\infty} (2l+1) (e^{ikr} - (-1)^l e^{-ikr}) P_l(\cos \theta)$$

Let's re-write these terms $f_l = \frac{1}{k} e^{i\delta_l} \sin \delta_l$ and solve the S.E.

(crunch crunch..)

Phase acquired in the collision

Incoming wave

$$R_l(r) \sim \frac{1}{2ikr} \frac{1}{\sqrt{4\pi(2l+1)}} (e^{2i\delta_l} e^{ikr} - (-1)^l e^{-ikr})$$

Radial part of $\Psi(\mathbf{r})$

Scattered wave

..a bit of theory

Cross section:

$$\sigma = \frac{4\pi}{k^2} \sum_l (2l + 1) (\sin \delta_l)^2$$



..a bit of theory

Cross section:

$$\sigma = \frac{4\pi}{k^2} \sum_l (2l + 1) (\sin \delta_l)^2$$

Summarizing: a plane wave – written in terms of sum of spherical waves – is changed after the collision since each spherical wave of the expansion gets a **different** phase shift

..a bit of theory

Cross section:

$$\sigma = \frac{4\pi}{k^2} \sum_l (2l + 1) (\sin \delta_l)^2$$

Summarizing: a plane wave – written in terms of sum of spherical waves – is changed after the collision since each spherical wave of the expansion gets a **different** phase shift

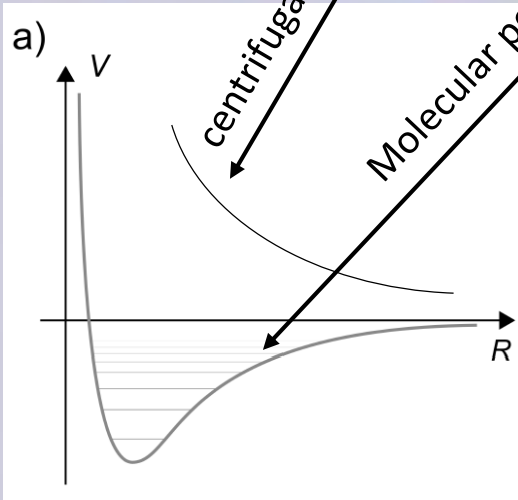
Important: energy conservation! The only effect of collisions are phase shifts that modify the «envelope» of the wave but not its amplitude!

..a bit of theory

How do we calculate the phase shifts?

We solve the Schroedinger equation! (radial part)

$$\left(-\frac{1}{r} \frac{d^2}{dr^2} r + \frac{l(l+1)}{r^2} + \frac{2\mu}{\hbar^2} V(r) \right) R_l(r) = k^2 R_l(r)$$

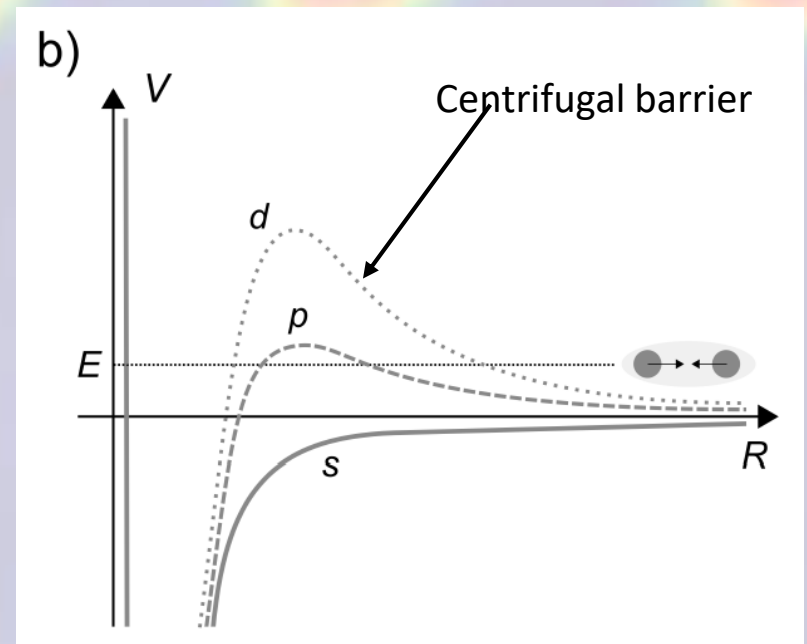
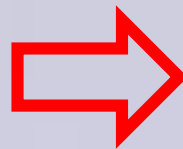
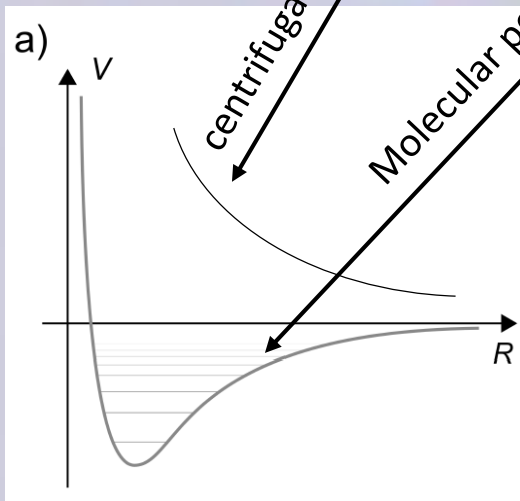


..a bit of theory

How do we calculate the phase shifts?

We solve the Schroedinger equation! (radial part)

$$\left(-\frac{1}{r} \frac{d^2}{dr^2} r + \frac{l(l+1)}{r^2} + \frac{2\mu}{\hbar^2} V(r) \right) R_l(r) = k^2 R_l(r)$$

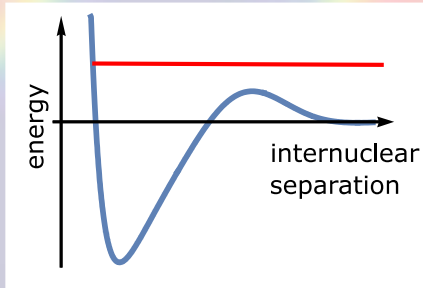
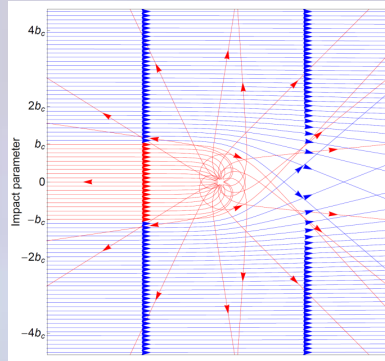


Atom-ion collisions

Large deflection

$$b < b_c$$

Collisions from the
hard wall

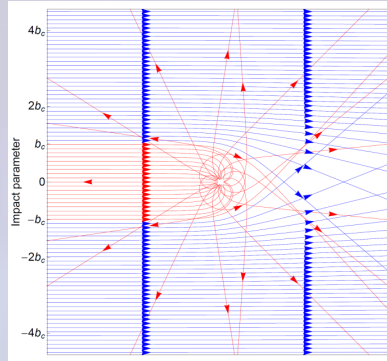
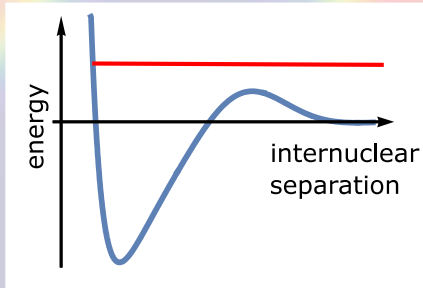


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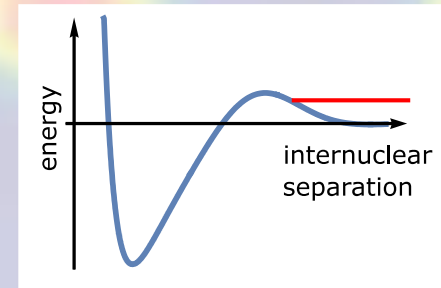
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Small deflection

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Collisions from the
centrifugal barrier

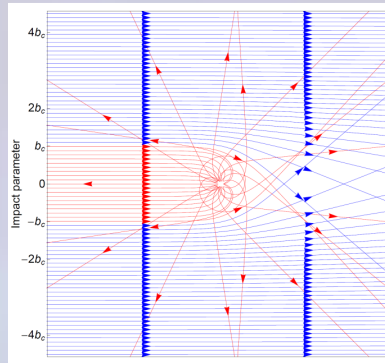


Atom-ion collisions

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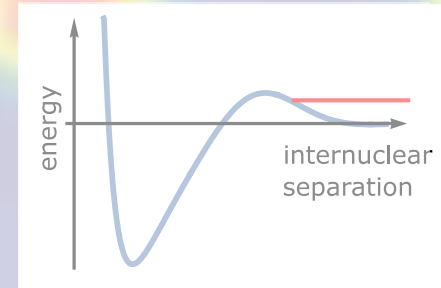
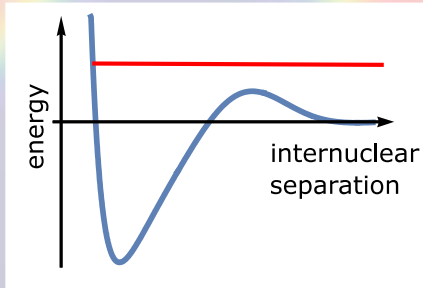
Collisions from the hard wall



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Collisions from the centrifugal barrier



↓

$$\sigma_L = \pi b_c^2 = \pi \sqrt{\frac{2C_4}{E_c}}$$

(classical)

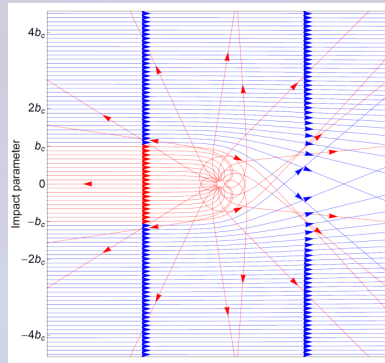


Atom-ion collisions

Large deflection

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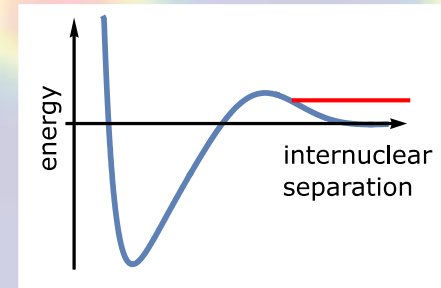
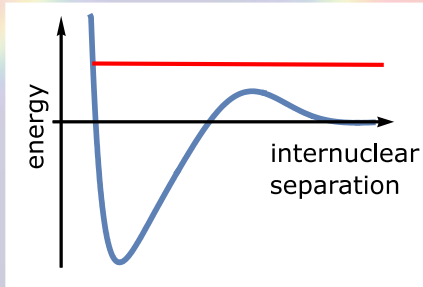
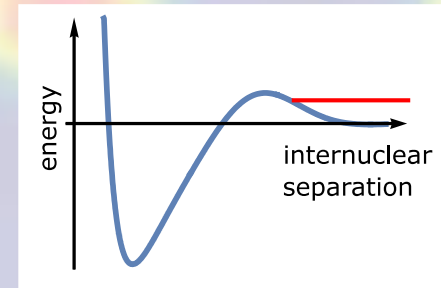
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$$\sigma_L = \pi b_c^2 = \pi \sqrt{\frac{2C_4}{E_c}}$$

(classical)

$$\sigma_{tot} = \pi \left(1 + \frac{\pi}{16}\right) \left(\frac{\mu C_4^2}{\hbar^2}\right)^2 E_c^{-1/3}$$

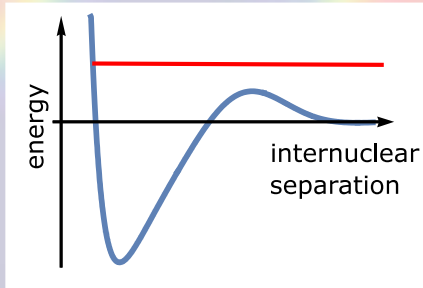
(semiclassical)

Atom-ion collisions

Large deflection

$$b < b_c$$

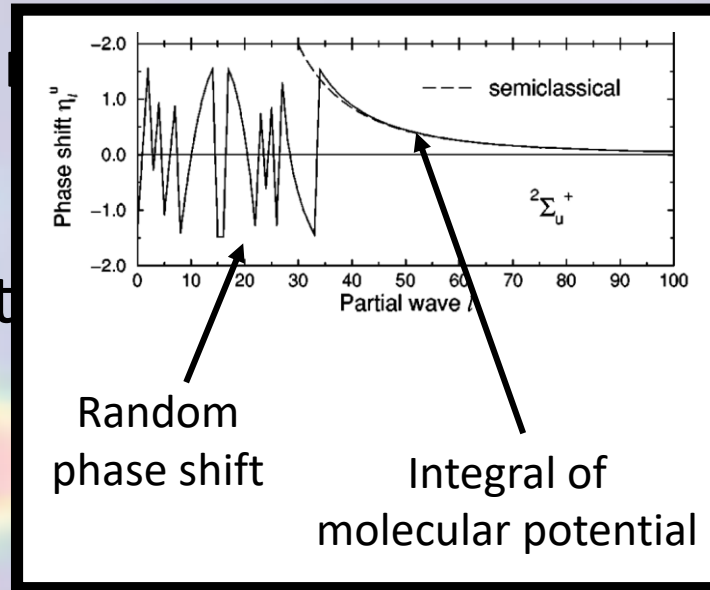
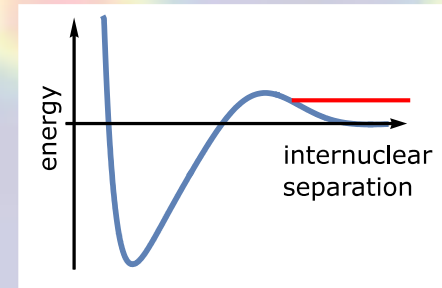
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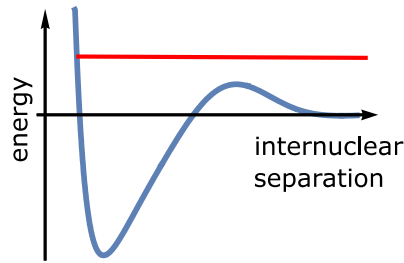
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Atom-ion collisions

Large deflection

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Langevin collisions



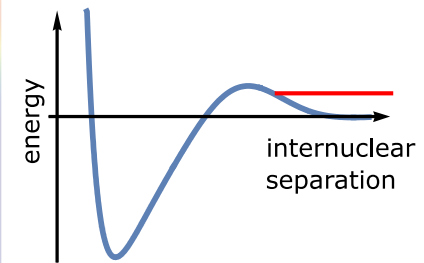
$$\sigma_L = \pi b_c^2 = \pi \sqrt{\frac{2C_4}{E_c}}$$

(classical)

Small deflection

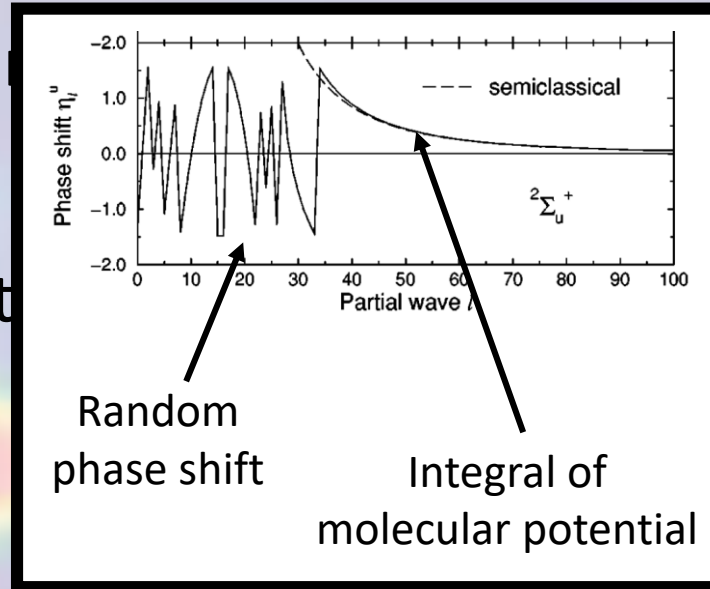
$$b > b_c$$

soft collisions



$$\sigma_{tot} = \pi \left(1 + \frac{\pi}{16}\right) \left(\frac{\mu C_4^2}{\hbar^2}\right)^2 E_c^{-1/3}$$

(semiclassical)



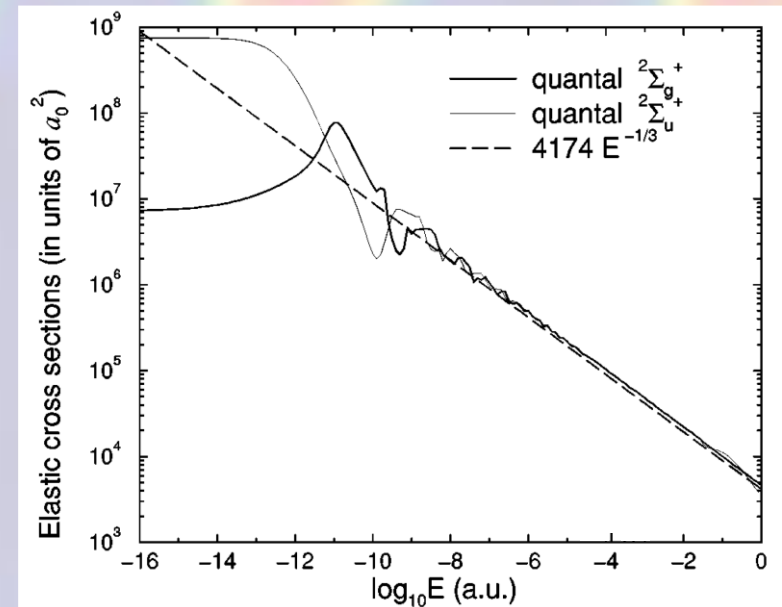
Atom-ion collisions

Cross section

$$\sigma_{tot} = \pi \left(1 + \frac{\pi}{16} \right) \left(\frac{\mu C_4^2}{\hbar^2} \right)^2 E_c^{-1/3}$$

Langevin cross section

$$\sigma_L = \pi \sqrt{2C_4} E_c^{-1/2}$$



Atom-ion collisions

Cross section

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Langevin cross section

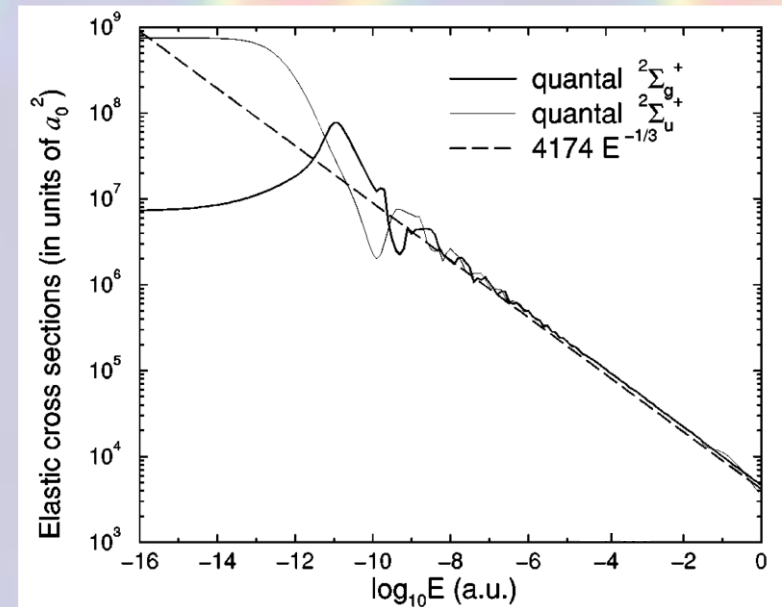
$$\sigma_L = \pi \sqrt{2C_4} E_c^{-1/2}$$

Scattering rate:

$$\gamma = n \sigma_L v$$

$$\gamma_{tot} = n \sigma_{tot} v \propto E_c^{1/6}$$

$$\gamma_L = n \sigma_L v = \text{const.}$$



Atom-ion collisions

Cross section

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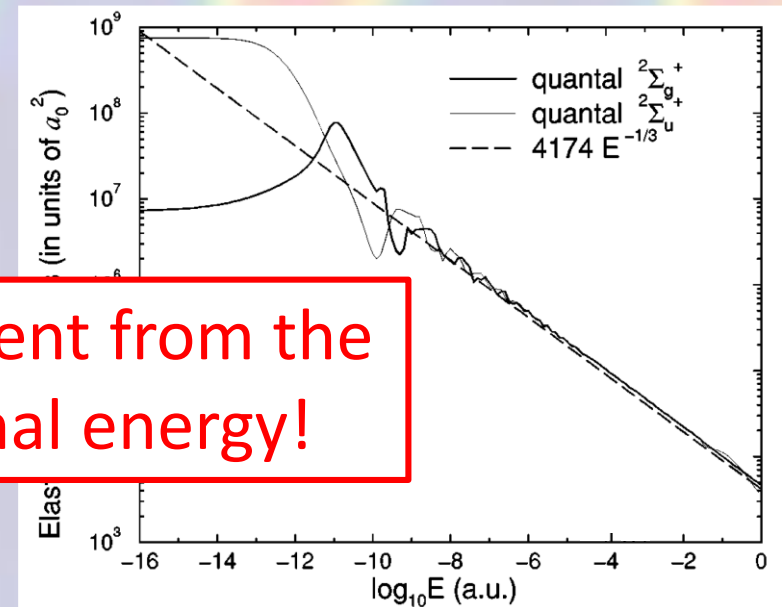
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Independent from the collisional energy!



Atom-ion collisions

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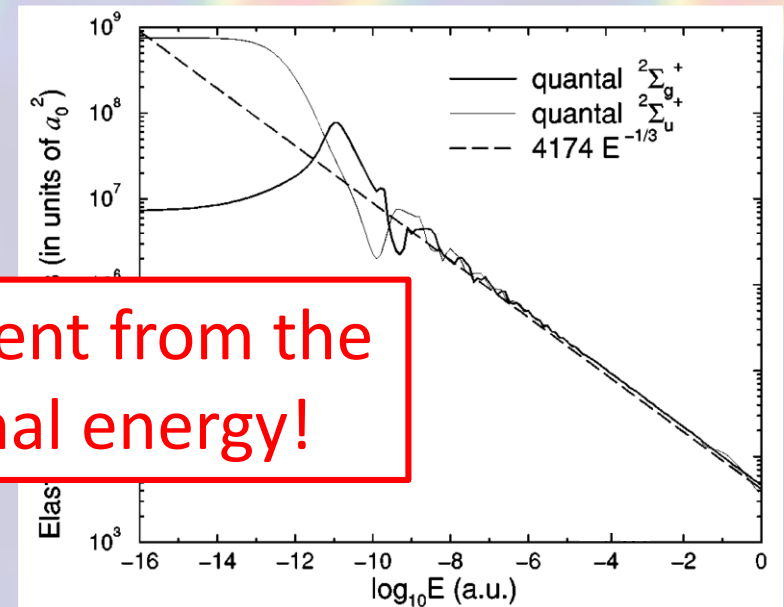
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Independent from the collisional energy!

At large collisional energies soft collisions are the most frequent



Outline

A brief intro of atom-ion collisions

Hybrid atom-ion experiments – state-of-the-art

Hybrid atom-ion experiments: perspectives

Atom-ion systems for improving ion clocks

The experiments

Paul trap + Atom trap

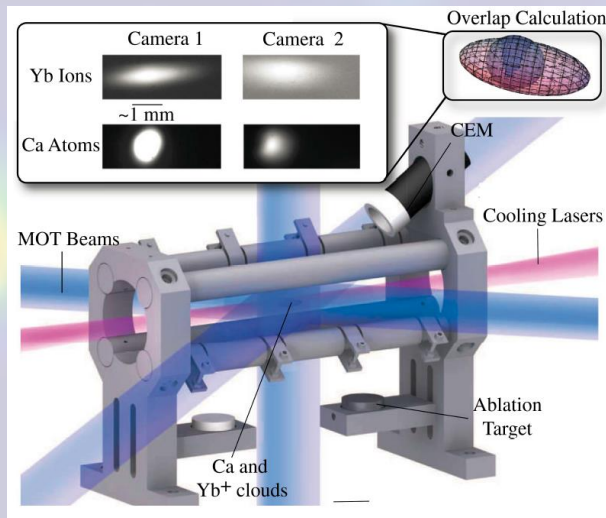
(Magneto-optical trap,
magnetic trap, optical trap..)



The experiments

Paul trap + Atom trap

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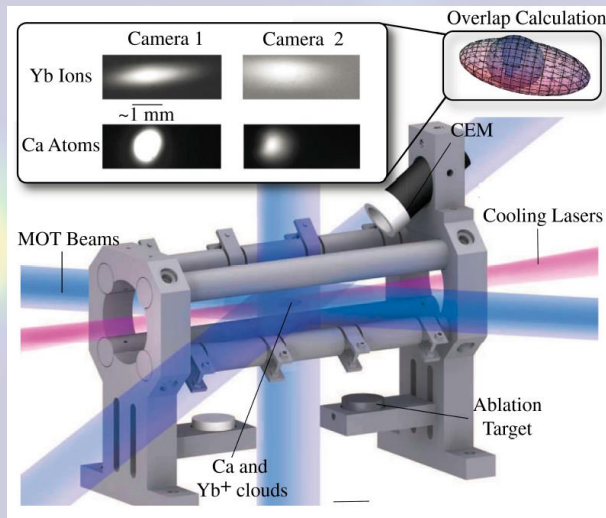


Hudson group, UCLA

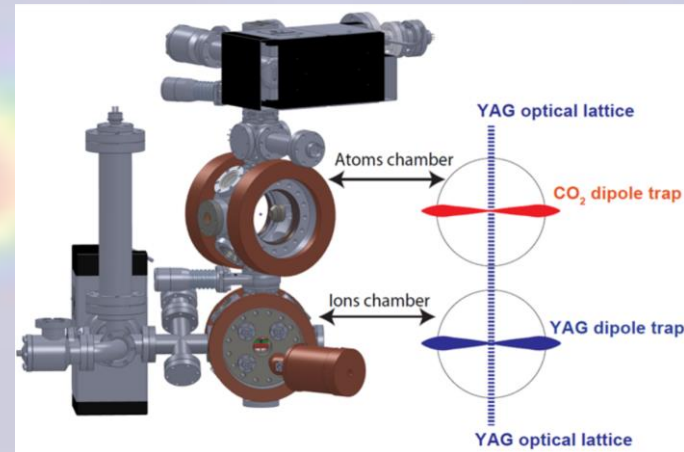
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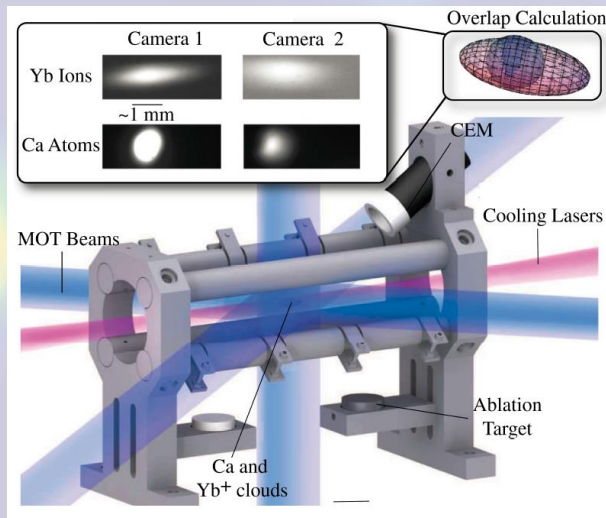


Ozeri group, Weizmann

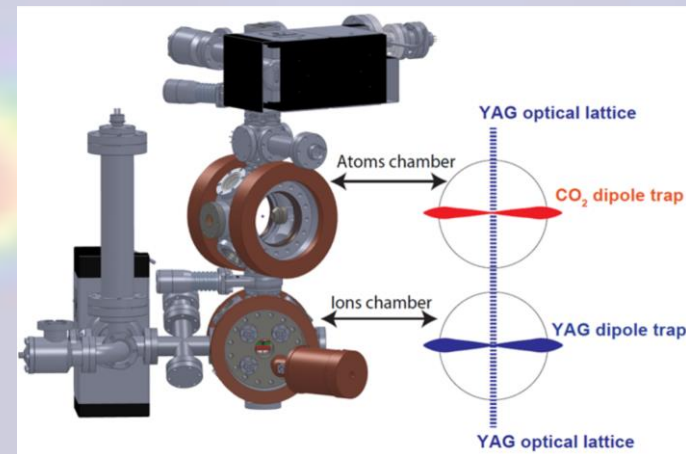
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Paul trap + Atom trap

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Hudson group, UCLA



Ozeri group, Weizmann

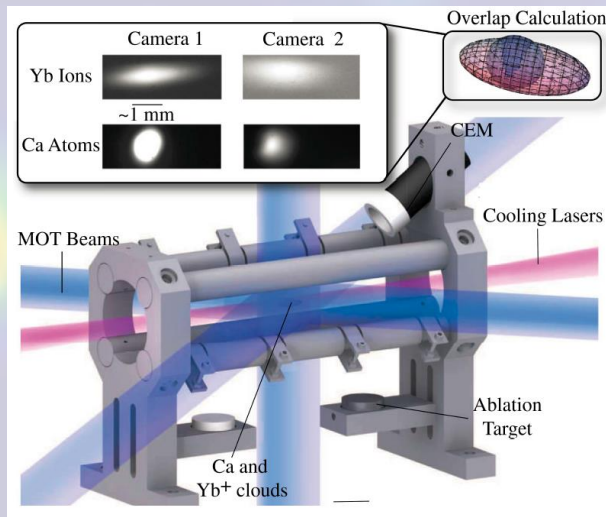
PROs: high repetition rate
large atom number

CONs: relatively high temperature,
poor control of internal state

The experiments

Paul trap + Atom trap

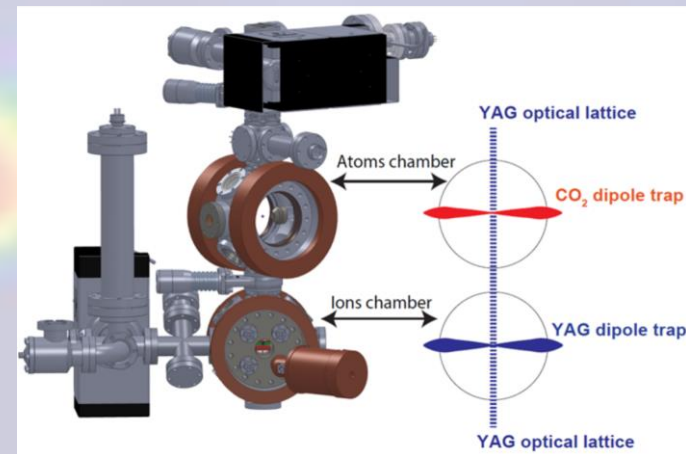
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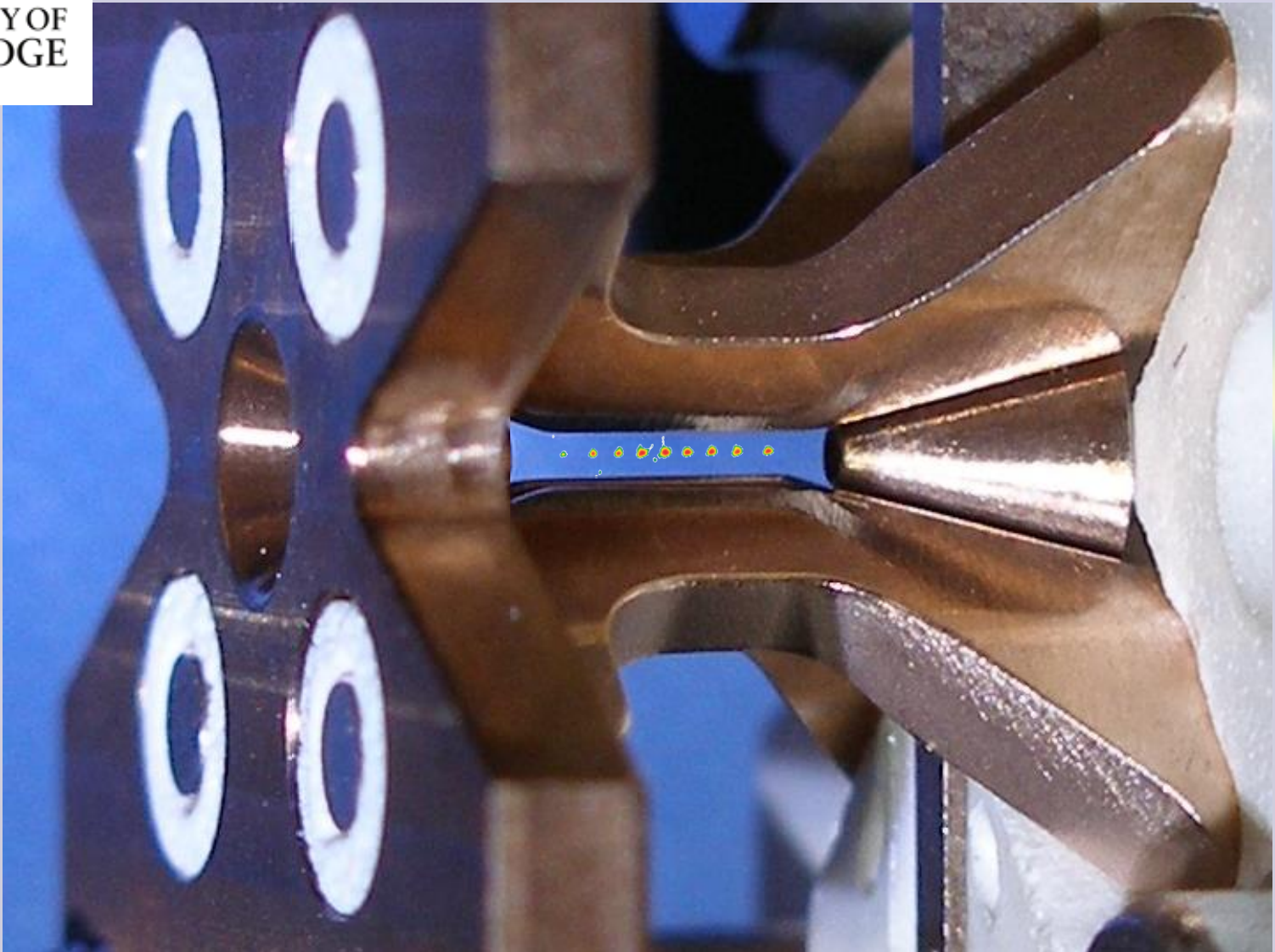
PROs: lower temperature,
coherence

CONs: slower repetition rate,
high collisional rate

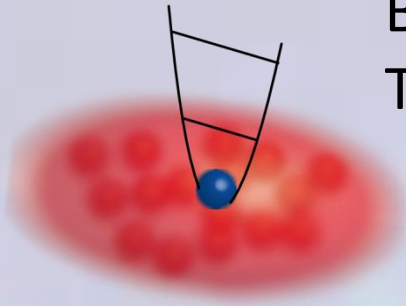


UNIVERSITY OF
CAMBRIDGE

A hybrid apparatus



Elastic collisions: sympathetic cooling



Bose-Einstein condensate

$T \approx 100$ nK

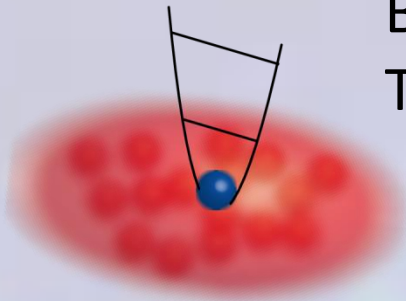
Goal: continuous cooling
of an ion-based quantum
hardware

Ion trap

(level spacing ≈ 5 mK)



Elastic collisions: sympathetic cooling



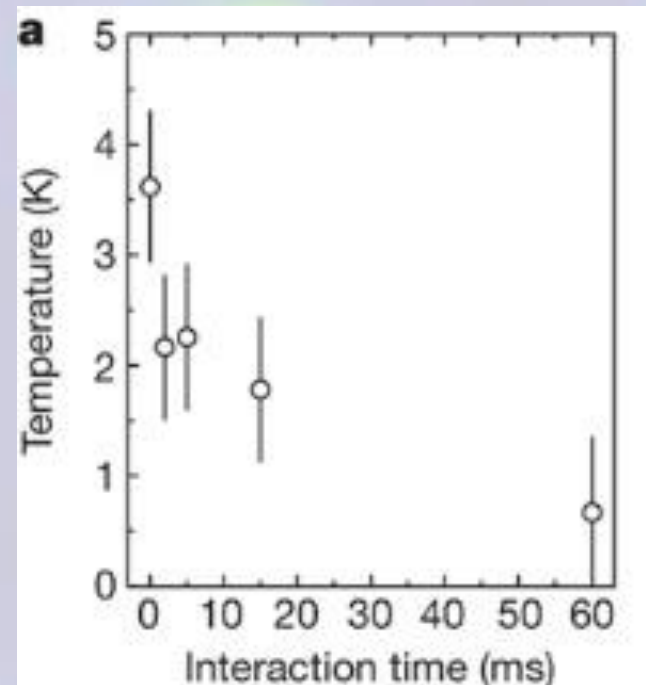
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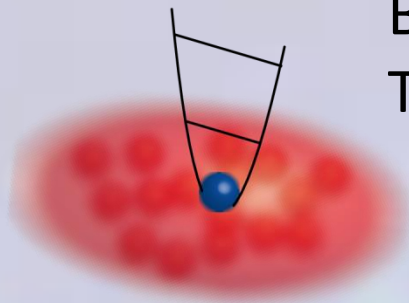
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Elastic collisions: sympathetic cooling

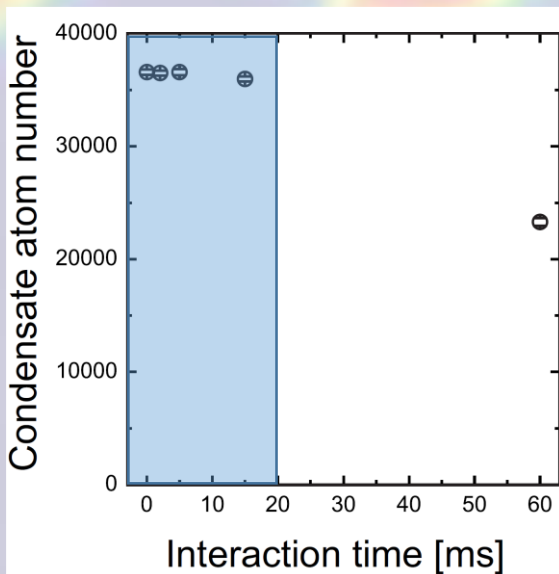


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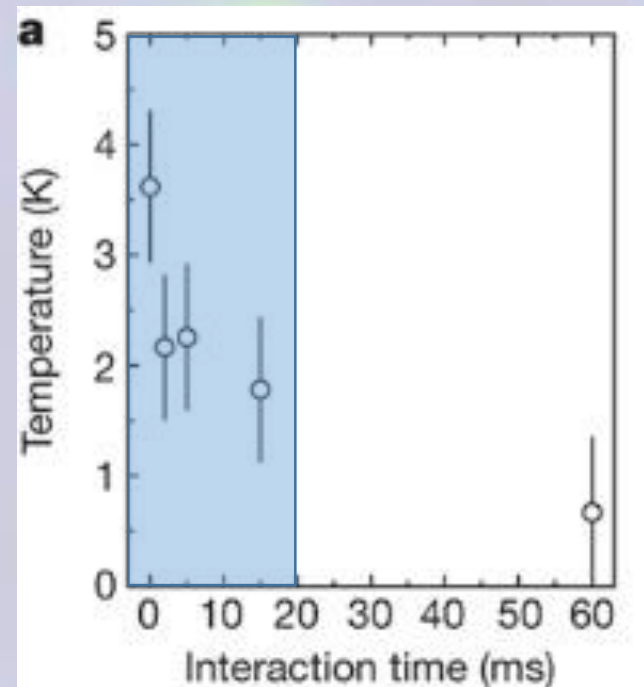
$T \approx 100$ nK

Goal: continuous cooling of an ion-based quantum hardware

Ion trap
(level spacing ≈ 5 mK)



**Very efficient!
(at the beginning)**



Inelastic collisions: chemical reactions

Possible effects of inelastic collisions: (example: Rb-Yb⁺ mixture)

- $\text{Yb}^+ + \text{Rb} \rightarrow \text{Yb} + \text{Rb}^+$ Charge Exchange
- $\text{Yb}^+ + \text{Rb} \rightarrow (\text{YbRb})^+$ Molecular association
- $(\text{Yb}^+)^* + \text{Rb} \rightarrow \text{Yb}^+ \text{Rb}$ Collisional quenching

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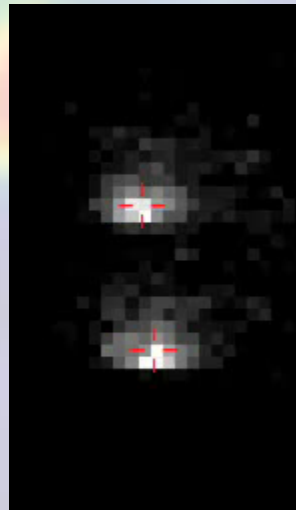
Problems:

1. New ion invisible to spectroscopy
2. Non-radiative decays cause trap losses
3. Heating mechanism

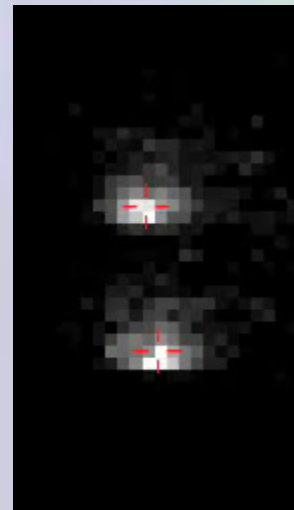
Inelastic collisions: chemical reactions

Detection of inelastic processes

Creation of a
dark ion



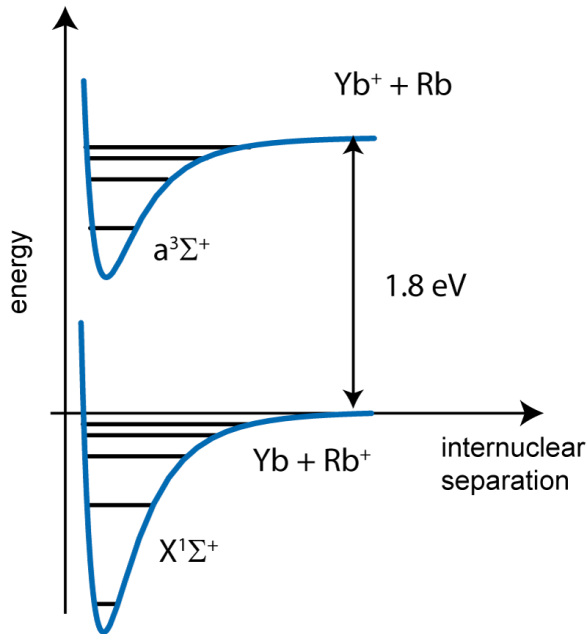
Ion loss



Not enough to ascertain what happened

Inelastic collisions: chemical reactions

Detection of inelastic processes

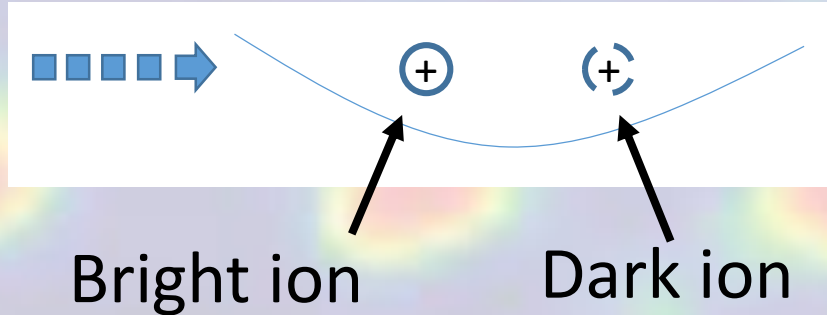


Only purely radiative processes make the ion survive

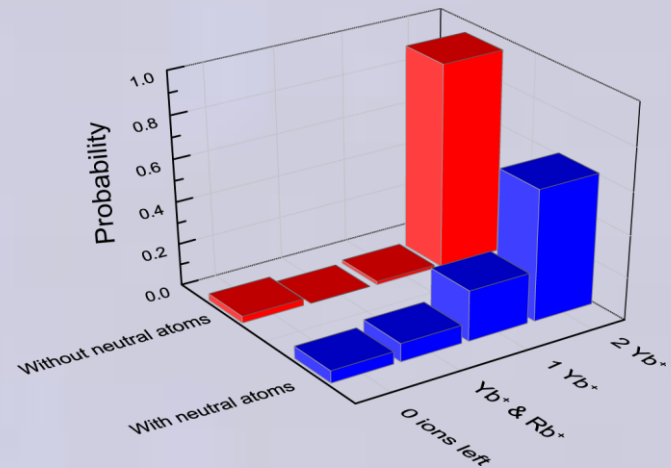
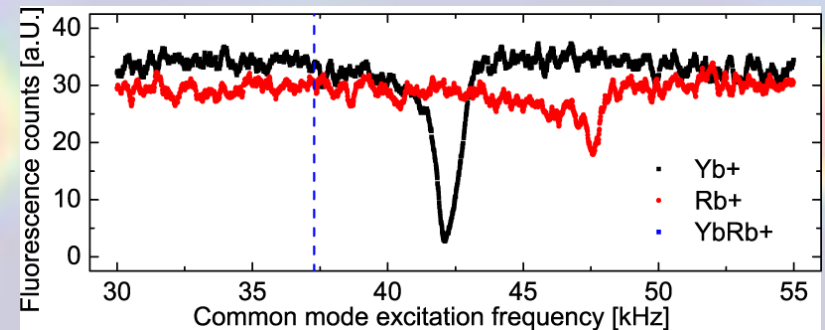
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Inelastic collisions: chemical reactions

We can perform mass spectrometry of a dark ion by modulating e.g. the cooling laser of an ancillary ion

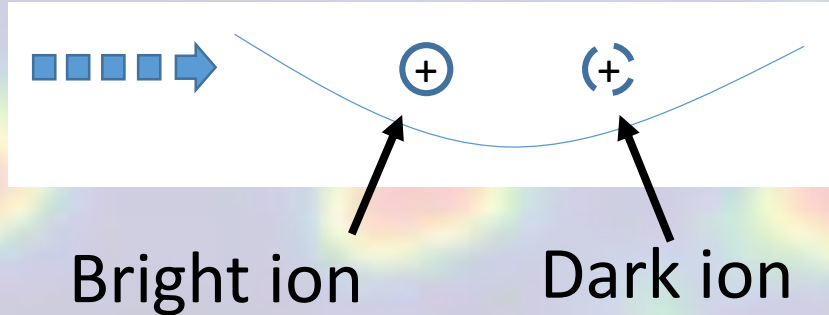


Example: Rb-Yb⁺

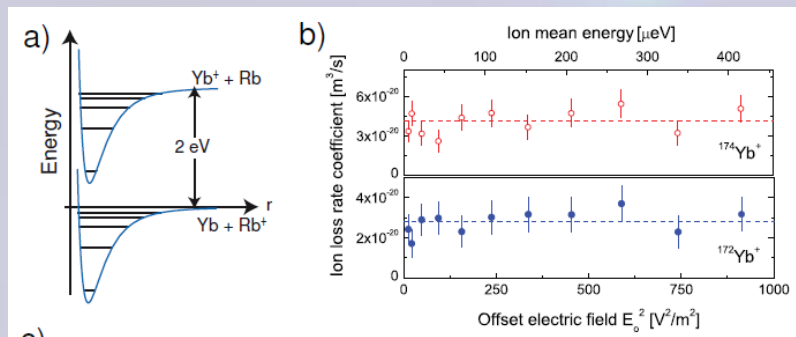
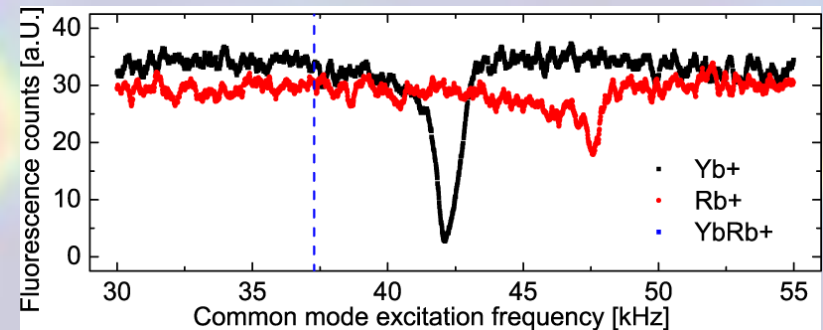


Inelastic collisions: chemical reactions

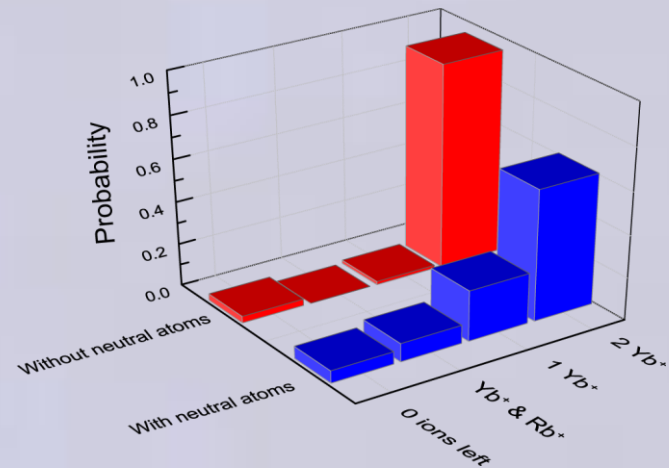
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Example: Rb-Yb⁺






Reaction rate independent from collisional energy: Langevin!



Controlling chemical reactions

everything depends on the ion-neutral pair

	Rb-Yb⁺ (Cambridge)	Rb-Sr⁺ (Weizmann)	Rb-Ca⁺ (Basel)
Charge exchange rate	$\sim 10^{-5} \gamma_L$	$\sim 2 \times 10^{-5} \gamma_L$	$< 10^{-3} \gamma_L$
Creation of molecules			

Controlling chemical reactions

everything depends on the ion-neutral pair
... and on the state

Table 1 | Measured proportionality constant ϵ and branching ratios.

	$^2S_{1/2}$	$^2D_{3/2}$	$^2F_{7/2}$	$^2P_{1/2}$
ϵ	$10^{-5 \pm 0.3}$	1.0 ± 0.2	0.018 ± 0.004	0.1 ± 0.2
Charged particle lost	65%	87%	84%	
Rb ⁺ identified	35%	12%	15%	
Dark Yb ⁺ identified		< 1%		
Hot ion (unidentified)			1%	
Number of events	283	754	225	

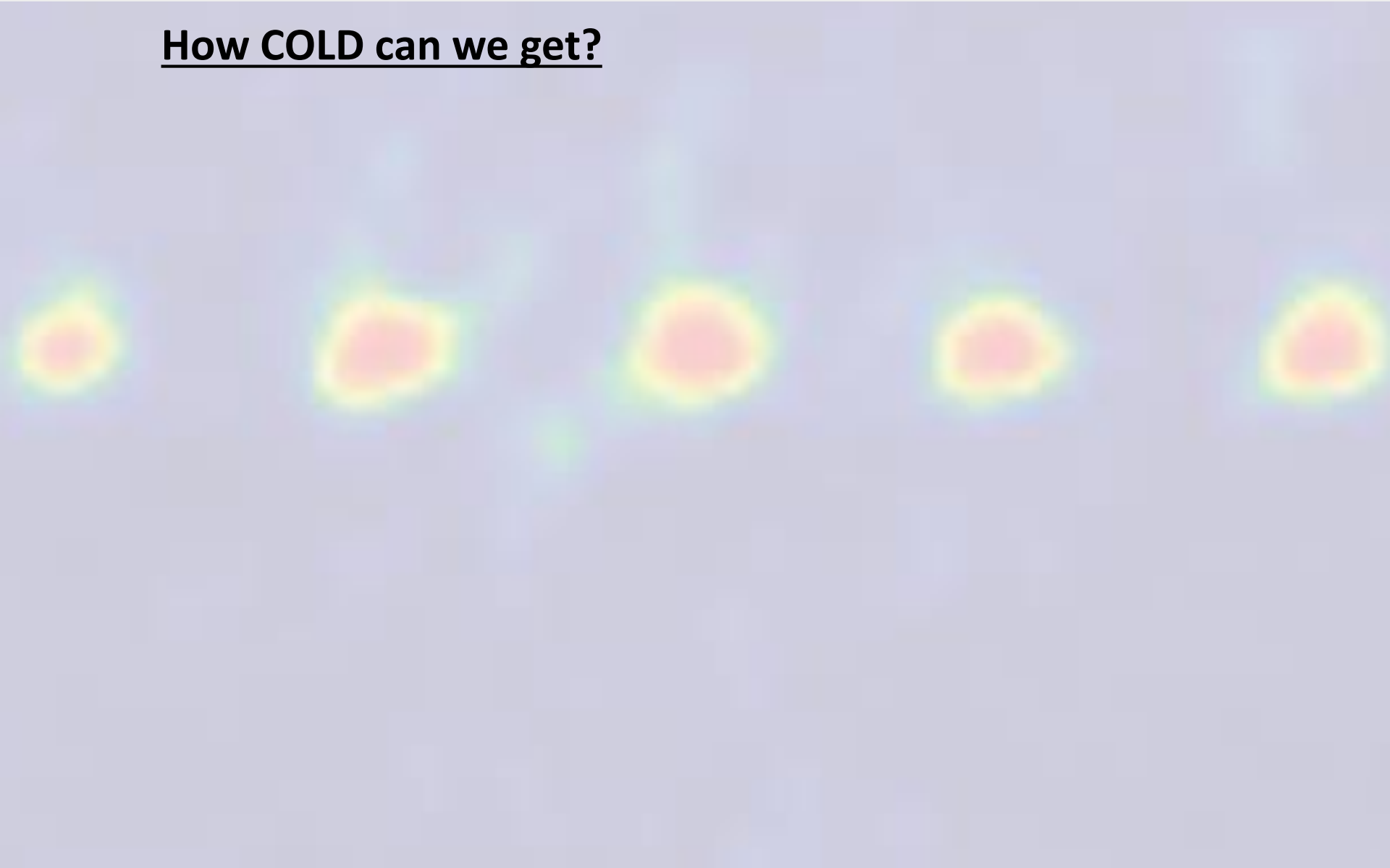
Rb-Yb⁺
(Cambridge)

$$\gamma_{react} = \gamma_L \epsilon$$

State-dependent reactivity

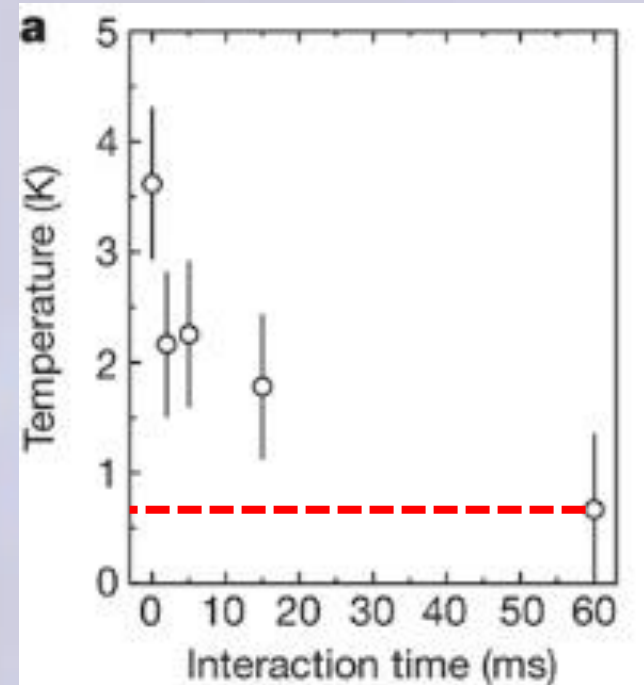
.. a problematic story

How COLD can we get?



.. a problematic story

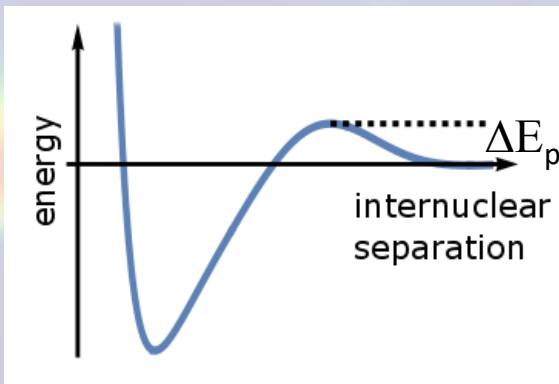
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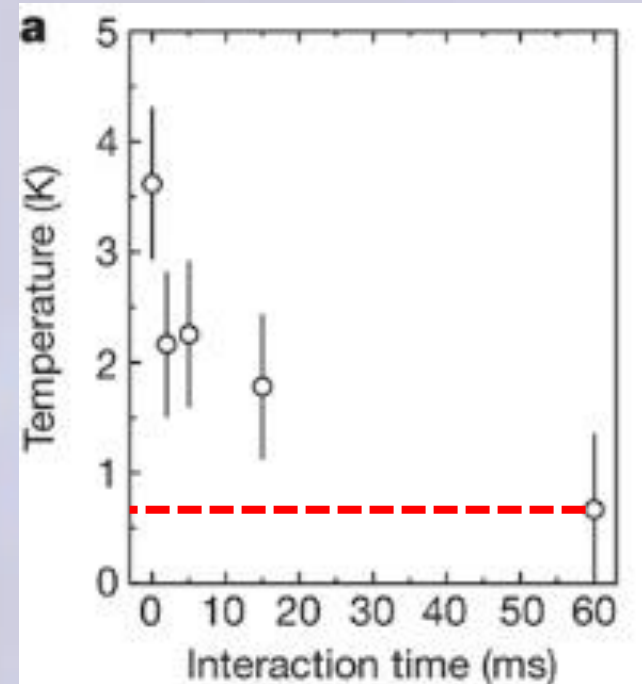
.. a problematic story

How COLD can we get?

At these temperature atom-ion collisions are **CLASSICAL**



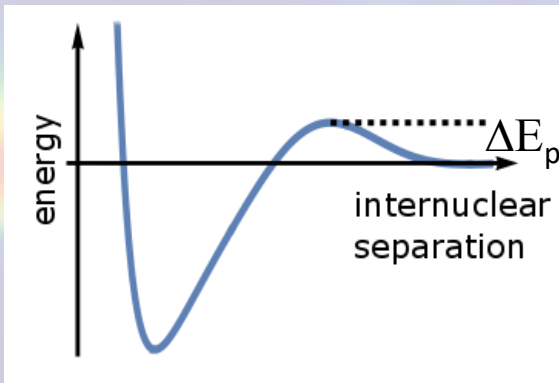
$$\Delta E_p / k_B \approx 1-10 \mu\text{K}$$



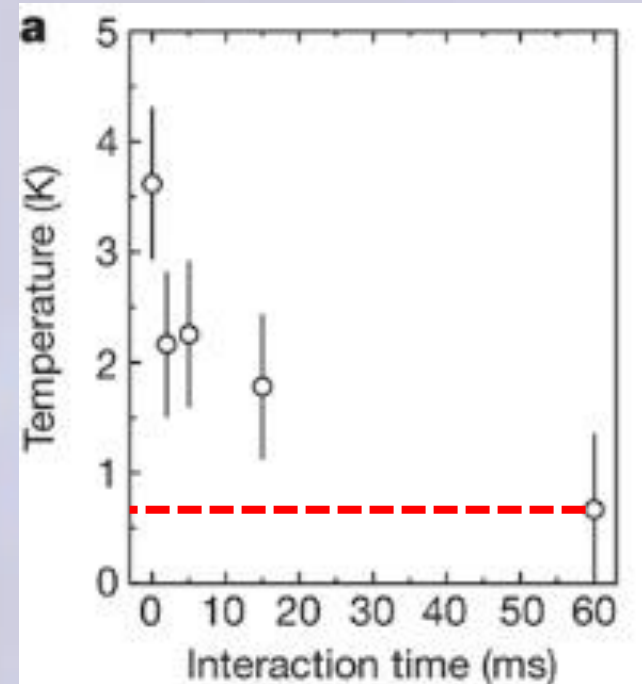
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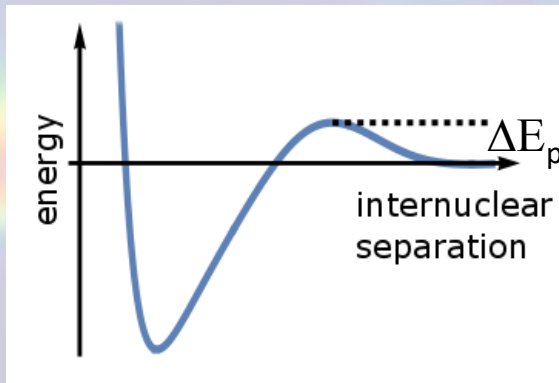
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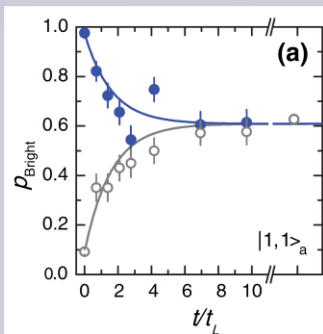
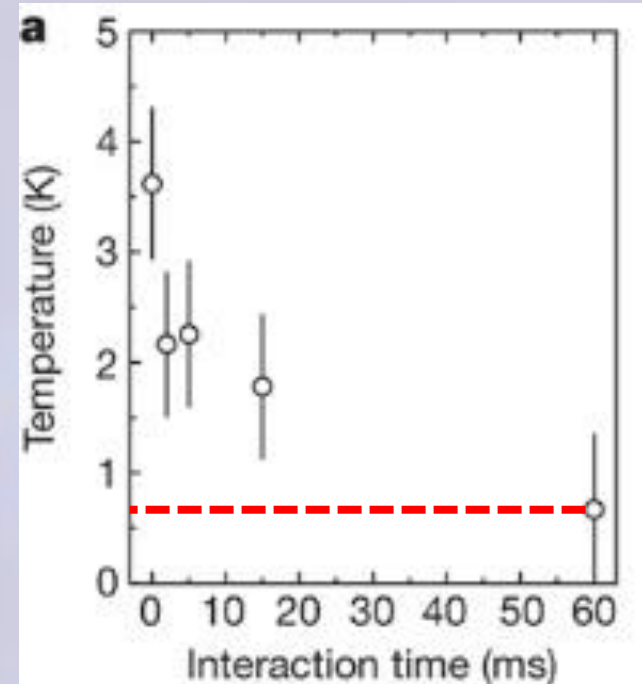
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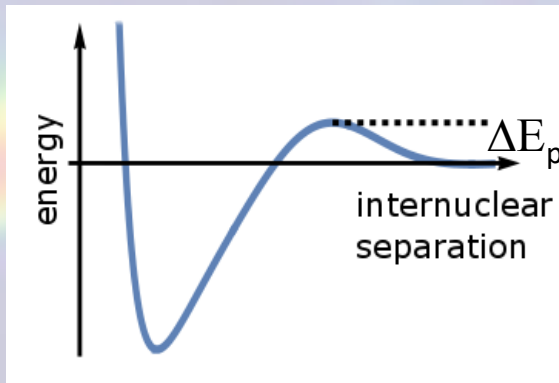


→ **Problem 1:** coherence does not survive

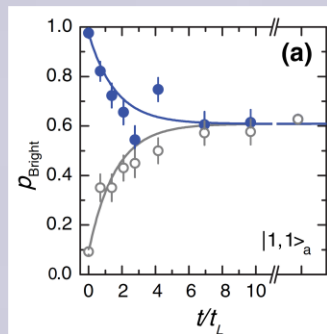
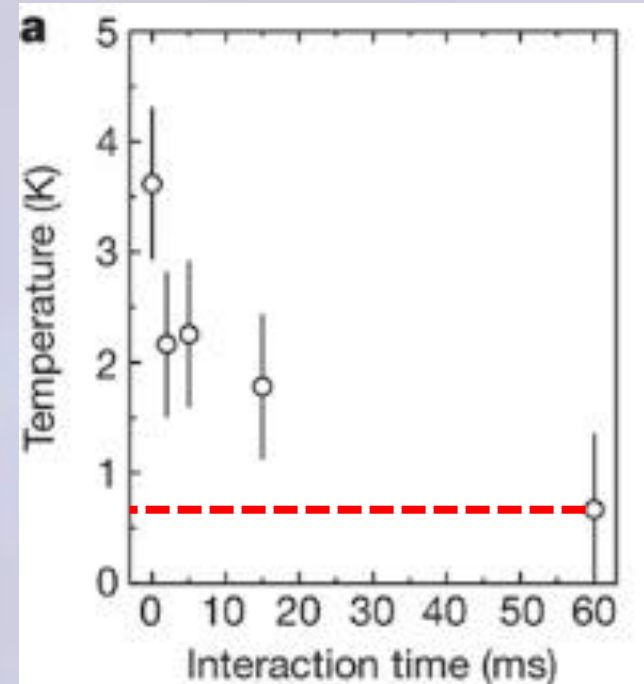
.. a problematic story

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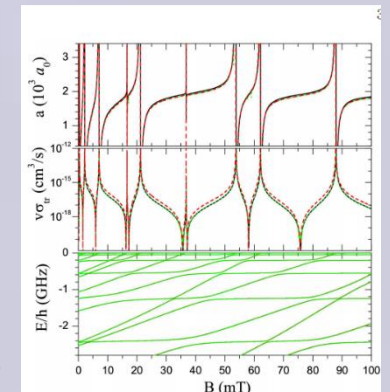


$$\Delta E_p / k_B \approx 1-10 \mu\text{K}$$



→ **Problem 1:** coherence does not survive

Problem 2 ← We can't see/use Feshbach resonances



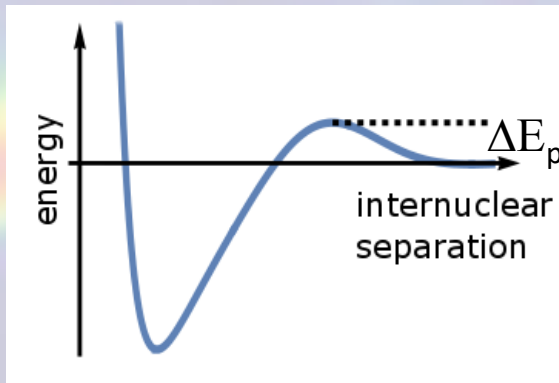
Z. Idziaszek et al. PRA 2009

L. Ratschbacher et al. PRL 2013

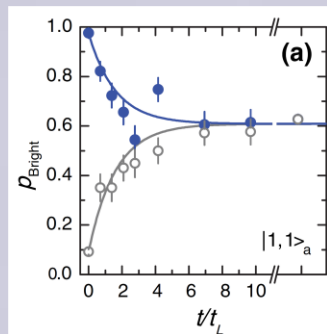
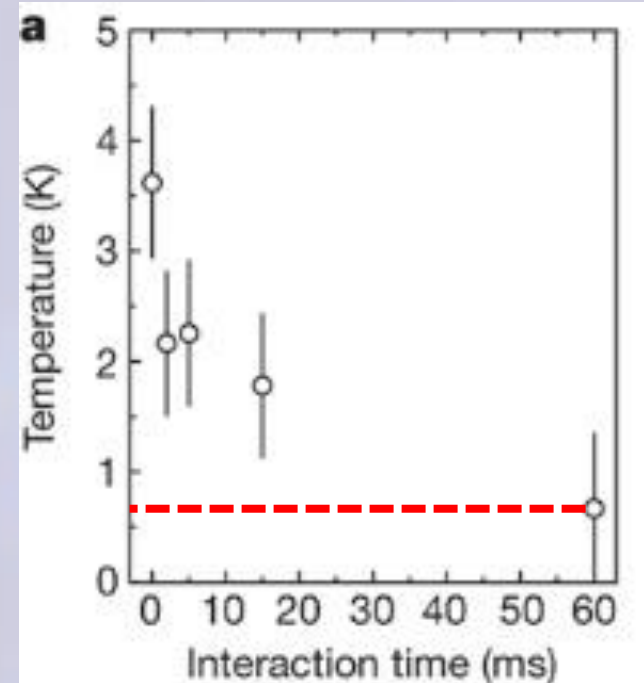
.. a problematic story

How COLD can we get?

At these temperature atom-ion collisions are **CLASSICAL**



$$\Delta E_p / k_B \approx 1-10 \mu\text{K}$$

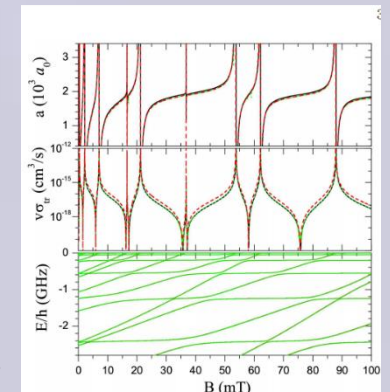


→ Problem 1: coherence

Physical problem:

We cannot bring atom-ion physics to a **quantum level**

Feshbach resonances



Z. Idziaszek et al. PRA 2009

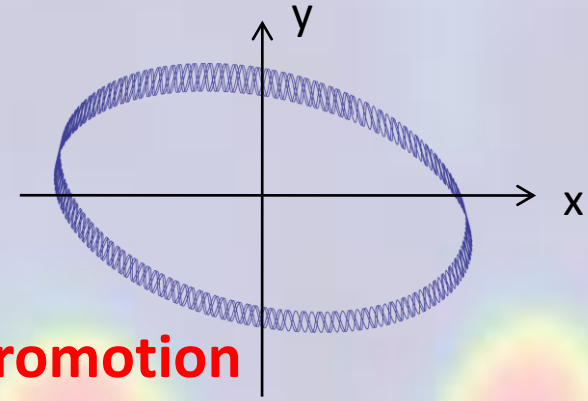
L. Ratschbacher et al. PRL 2013

Origin of problems: micromotion

$$r_{ion,x} = A_x \sin(\omega_x t + \varphi_x) \left[1 + \frac{q}{2} \sin(\Omega_T t) \right]$$
$$r_{ion,y} = A_y \sin(\omega_y t + \varphi_y) \left[1 - \frac{q}{2} \sin(\Omega_T t) \right]$$

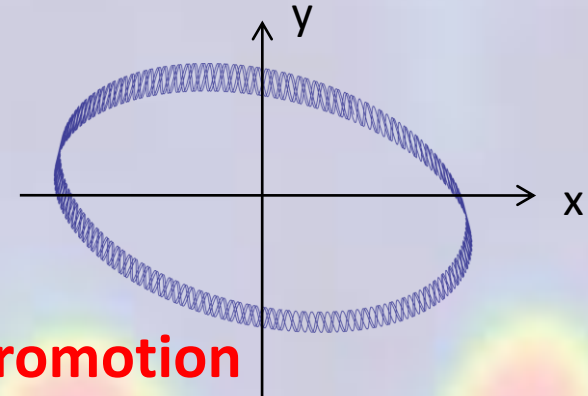
Secular motion

Micromotion

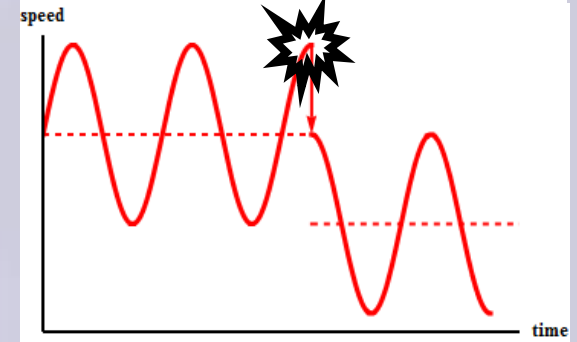
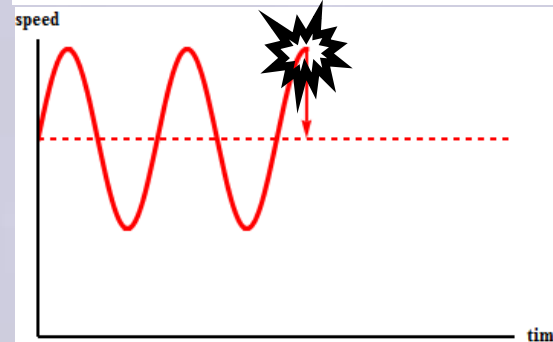
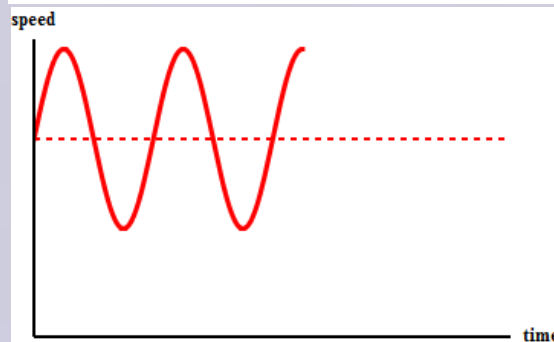
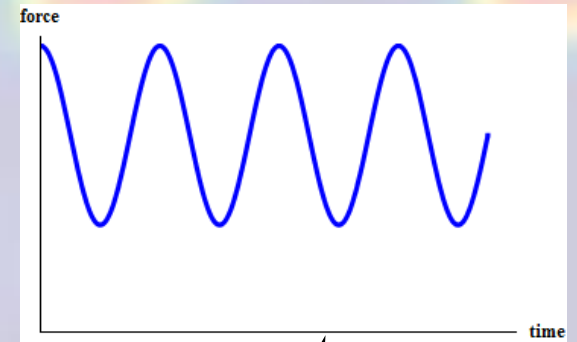
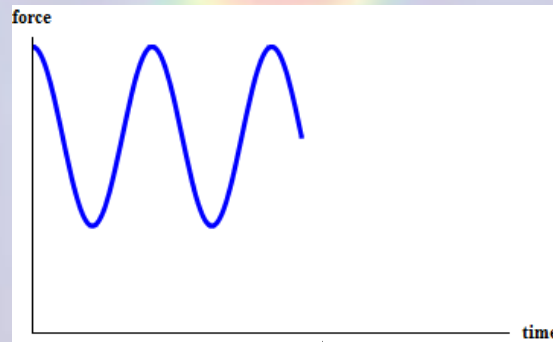
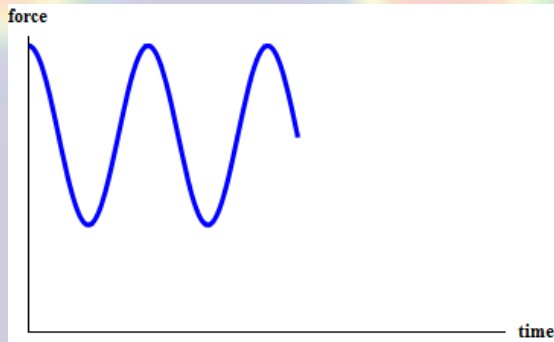


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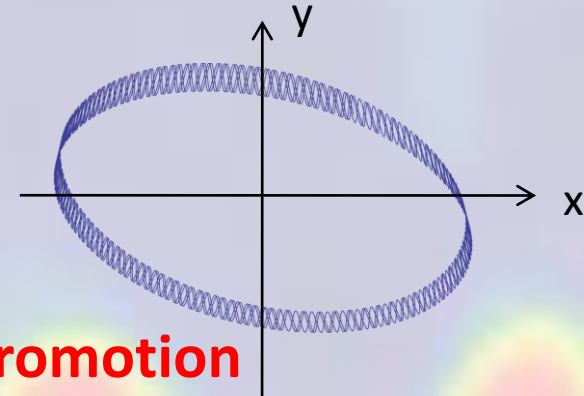
Secular motion **Micromotion**



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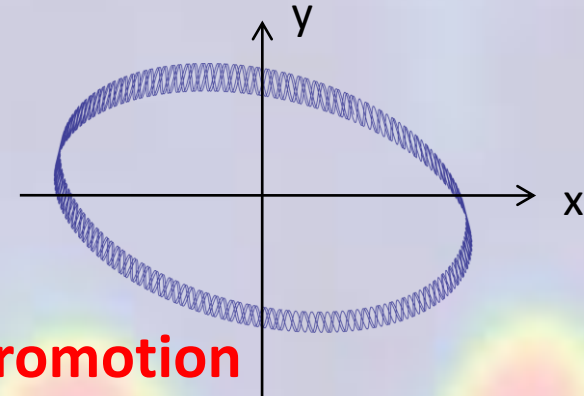
Collisions in the presence of micromotion can couple energy from the radiofrequency field to the colliding particles

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→ **Secular motion**

→ **Micromotion**

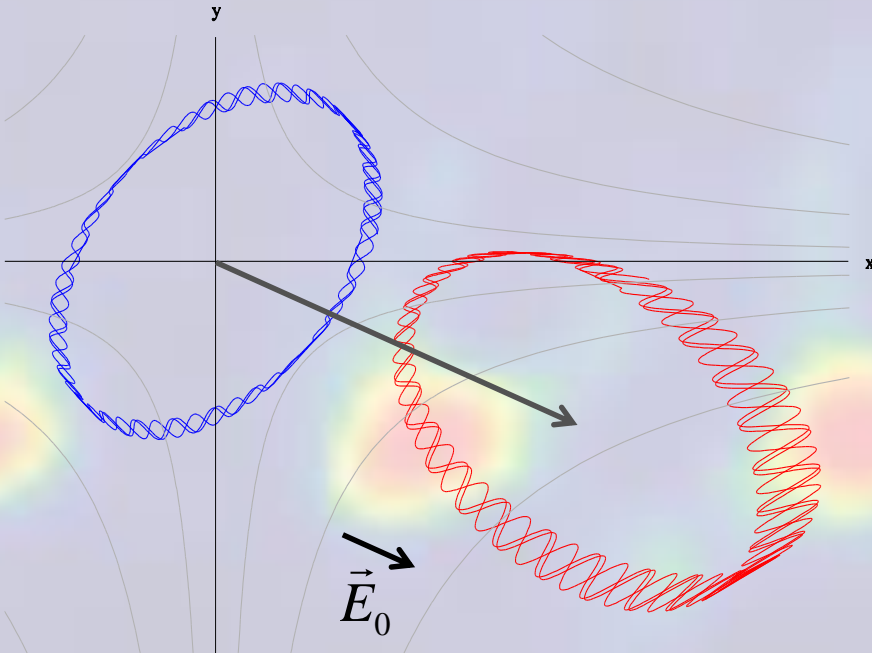


Collisions in the presence of micromotion can couple energy from the radiofrequency field to the colliding particles

Problem = Opportunity!

We can use micromotion to tune the collisional energy

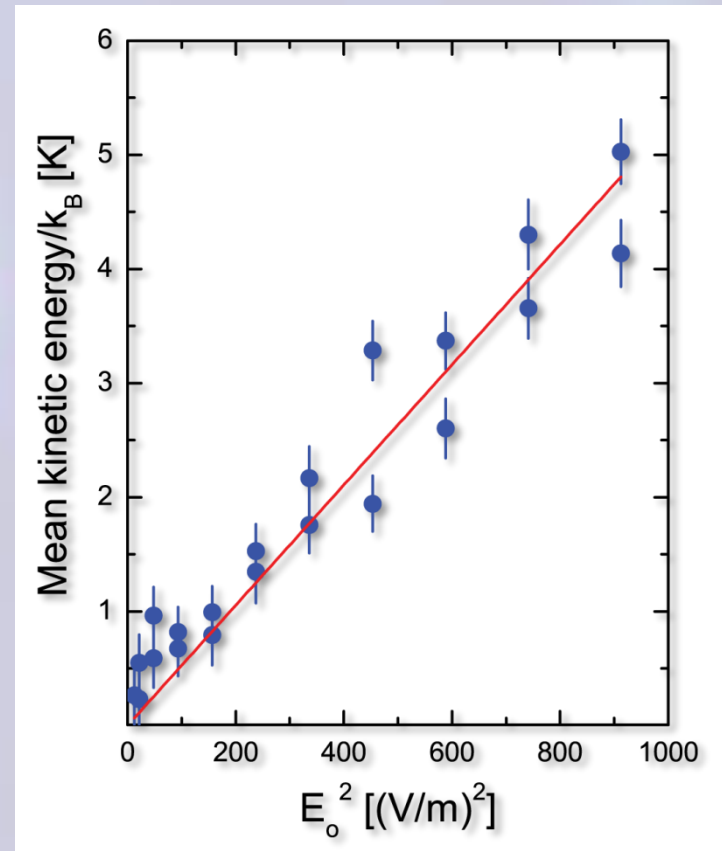
Tuning the collisional energy



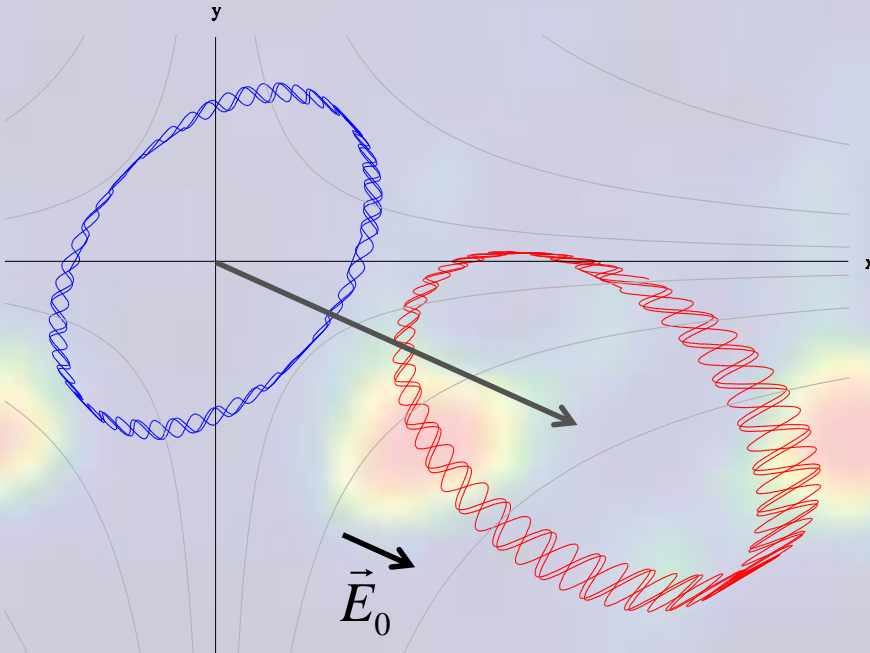
By applying an offset electric field \vec{E}_0 the mean kinetic energy changes:

$$\Delta E \approx \frac{1}{2} m \Omega^2 a_{mm}^2 \propto E_0^2$$

Where a_{mm} is the micromotion amplitude



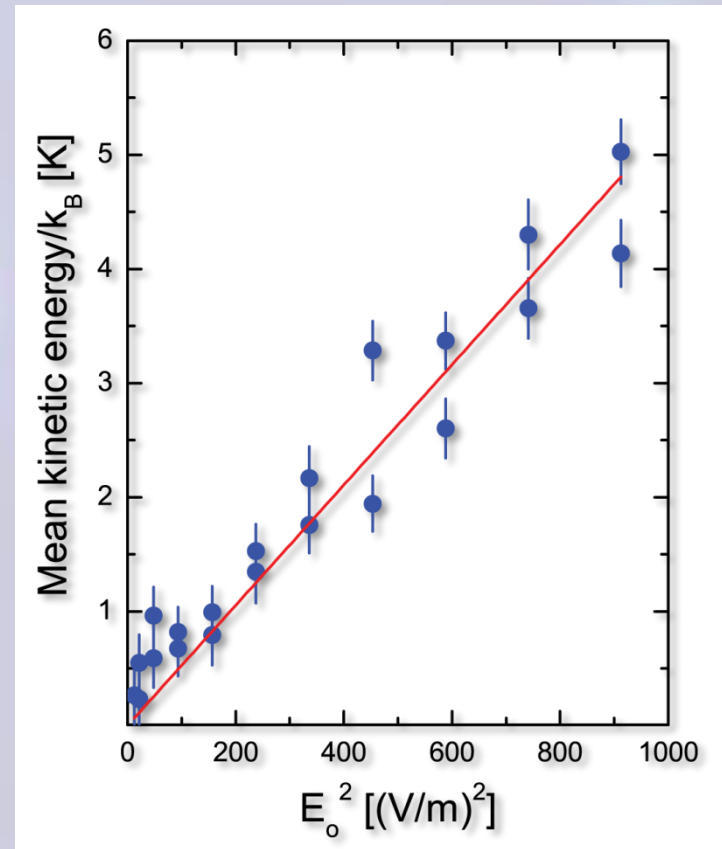
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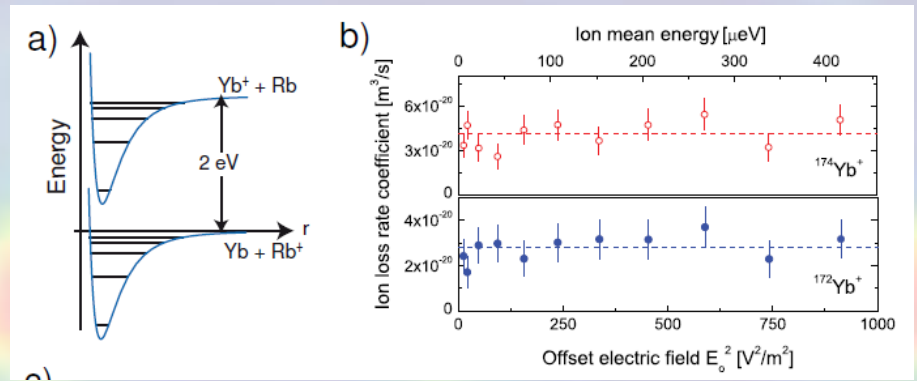
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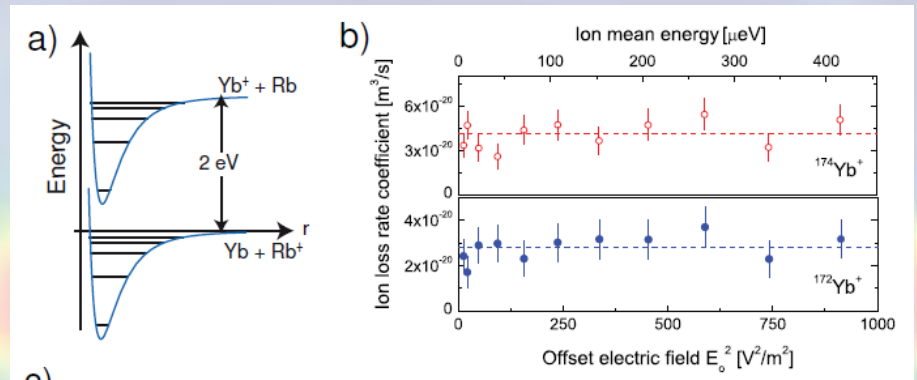
Studying ion-neutral collisions

Langevin collisions (no energy dependence of the scattering rate constant)

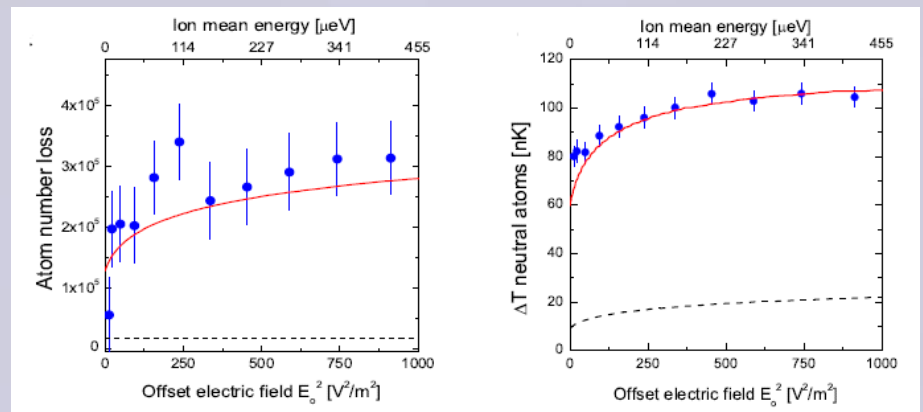


Studying ion-neutral collisions

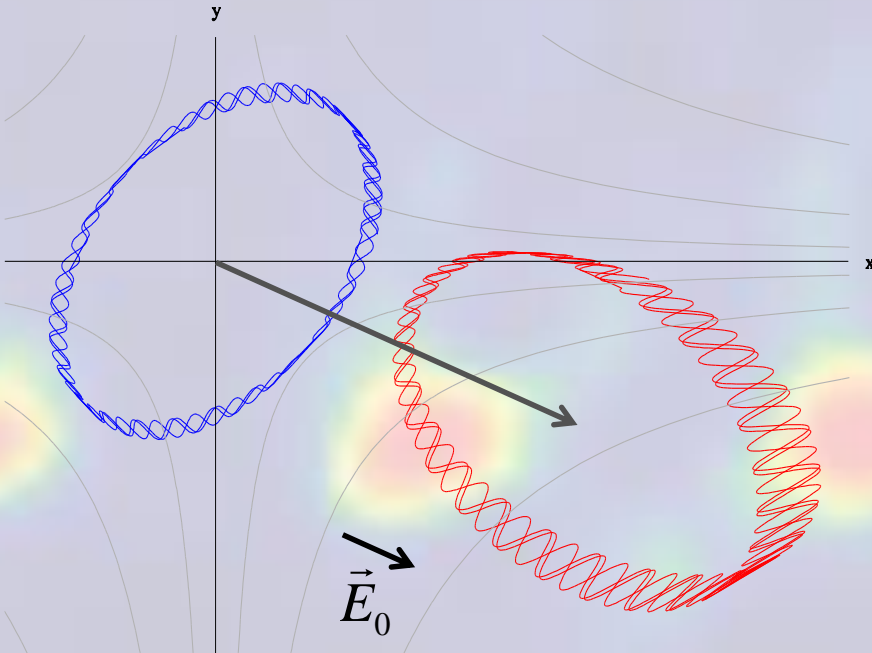
Langevin collisions (no energy dependence of the scattering rate constant)



«Soft» collisions (energy dependent scattering rate constant)



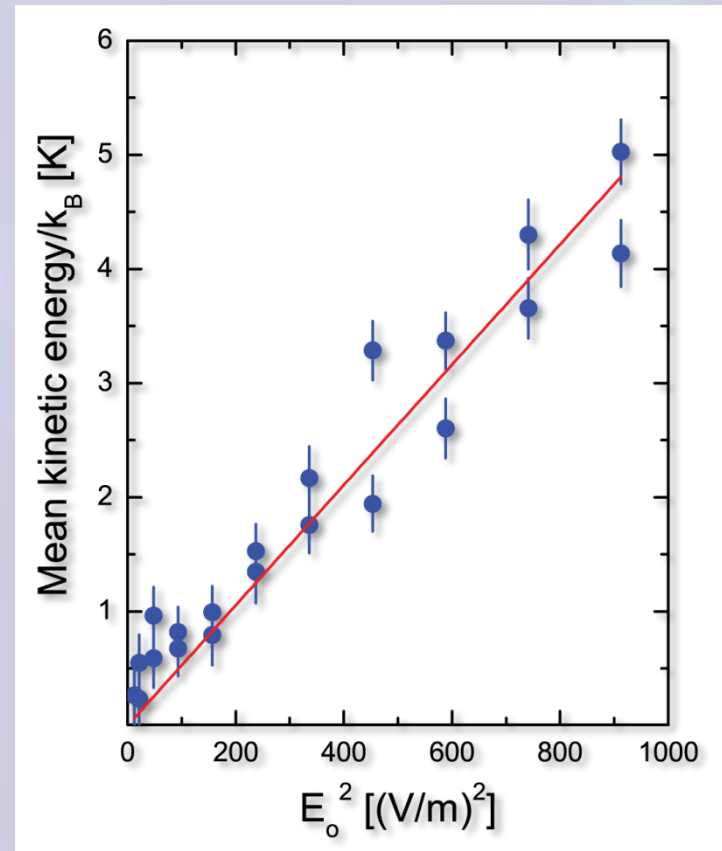
Tuning the collisional energy



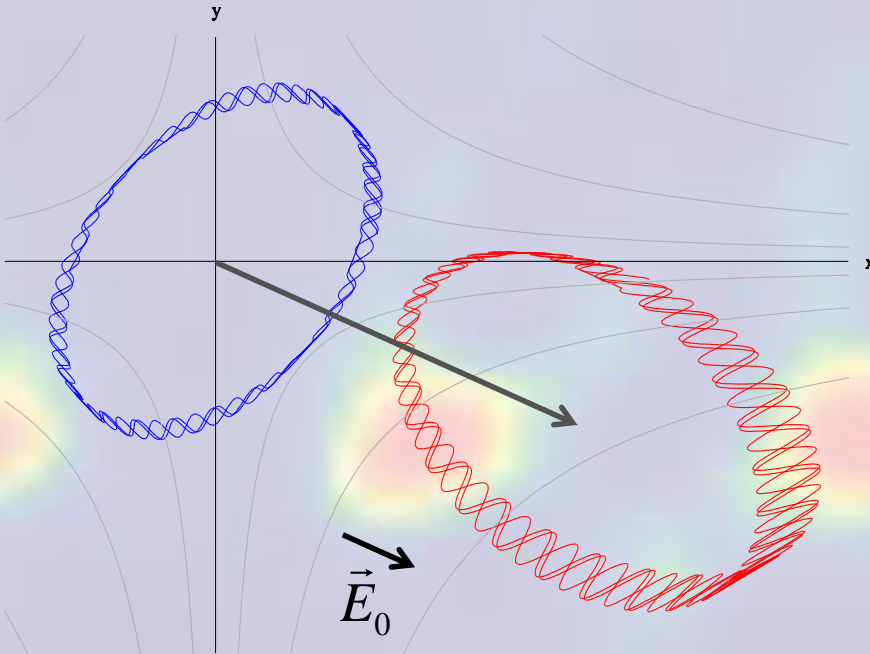
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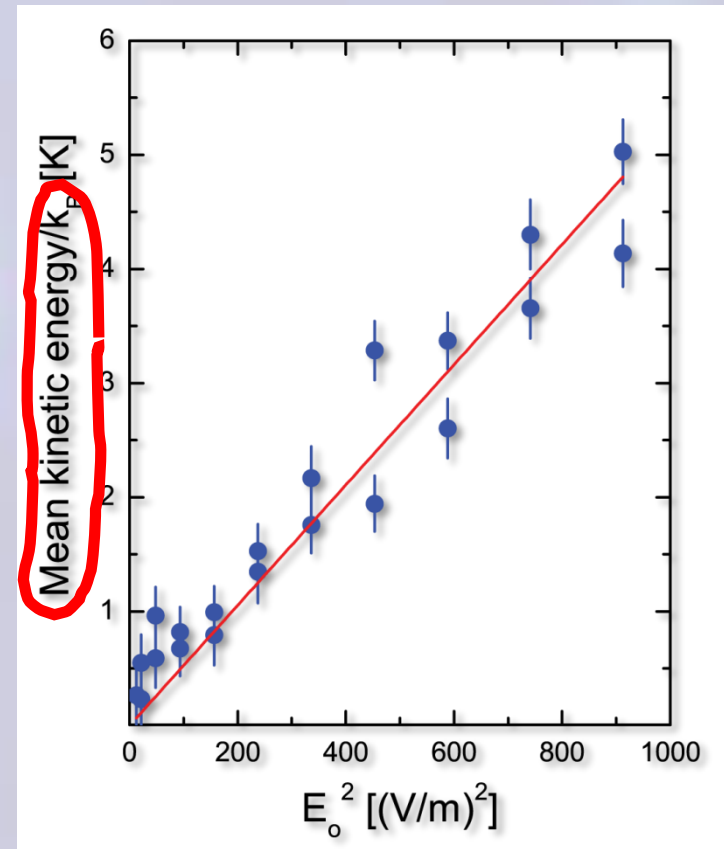
Tuning the collisional energy



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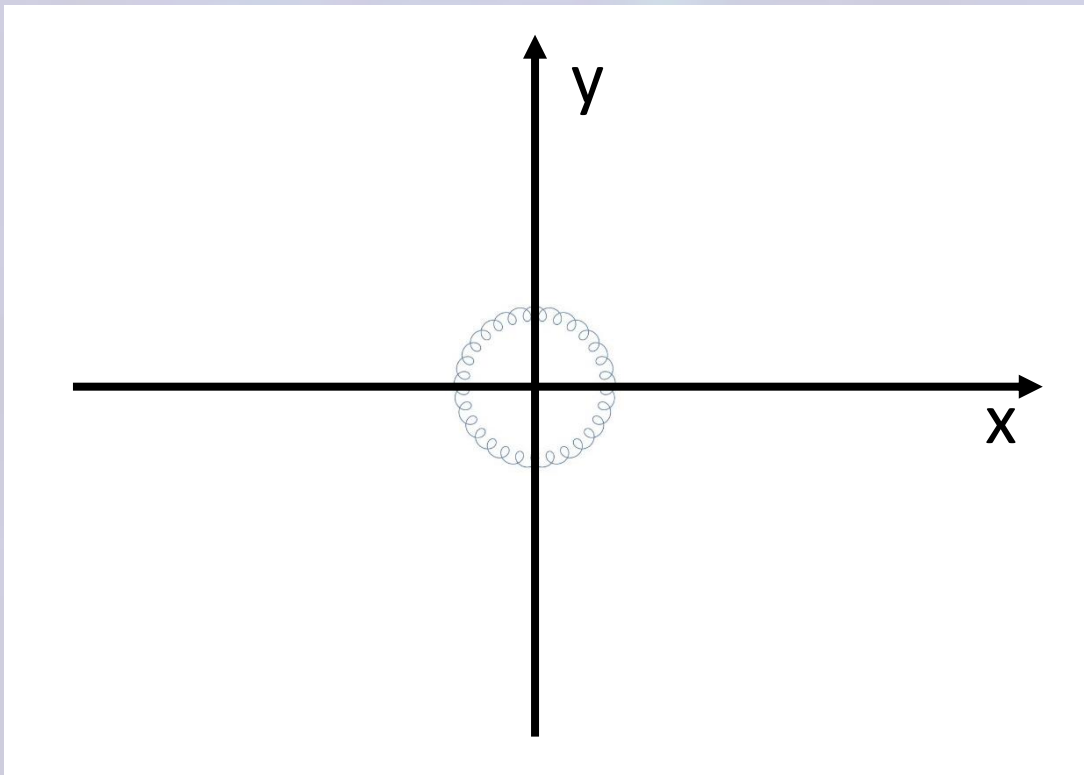
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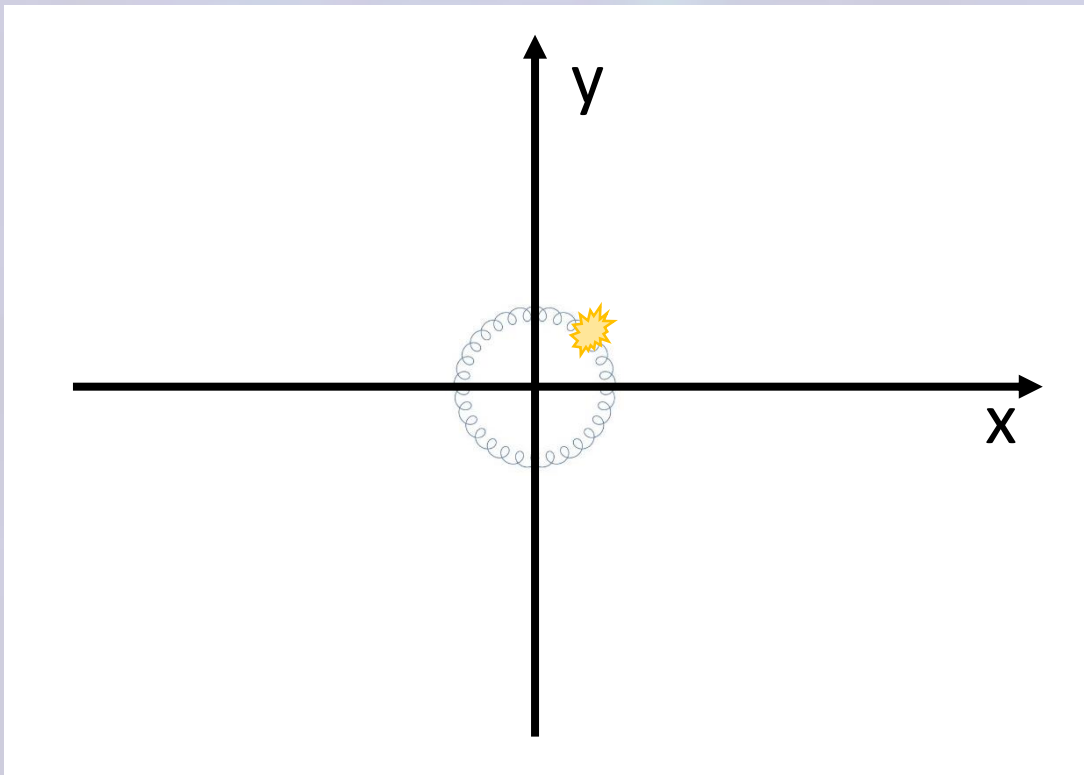
Distribution energy

The energy distribution of an ion in a buffer gas is *not* a thermal distribution



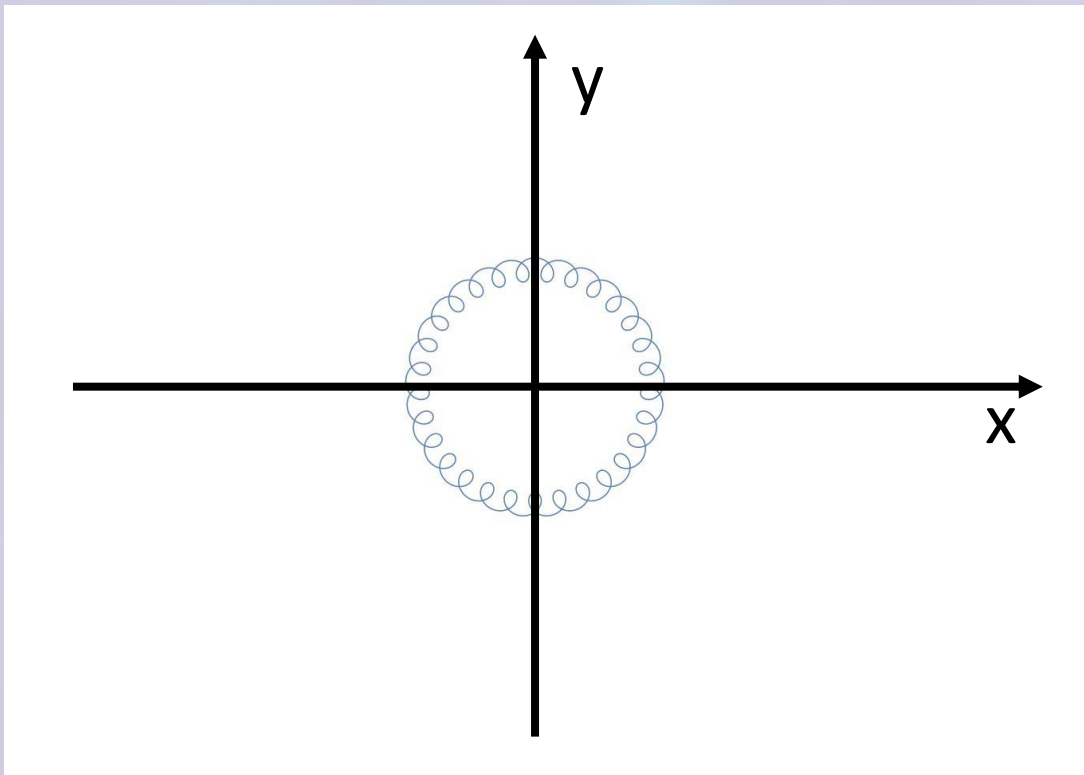
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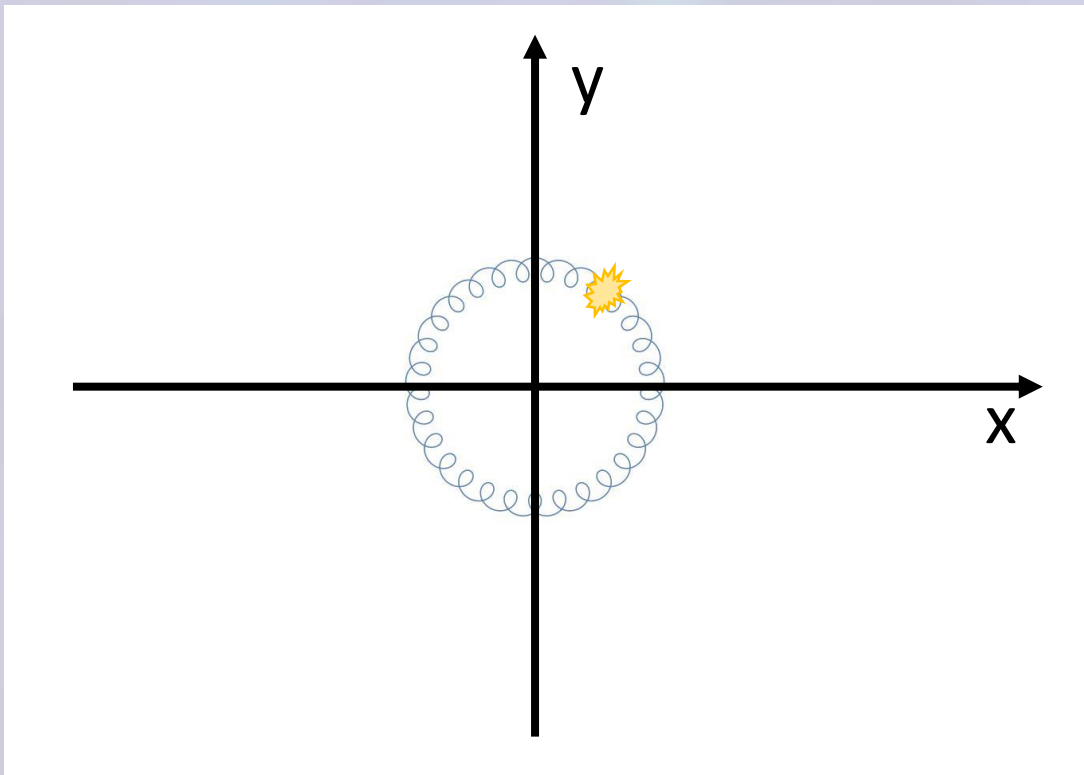
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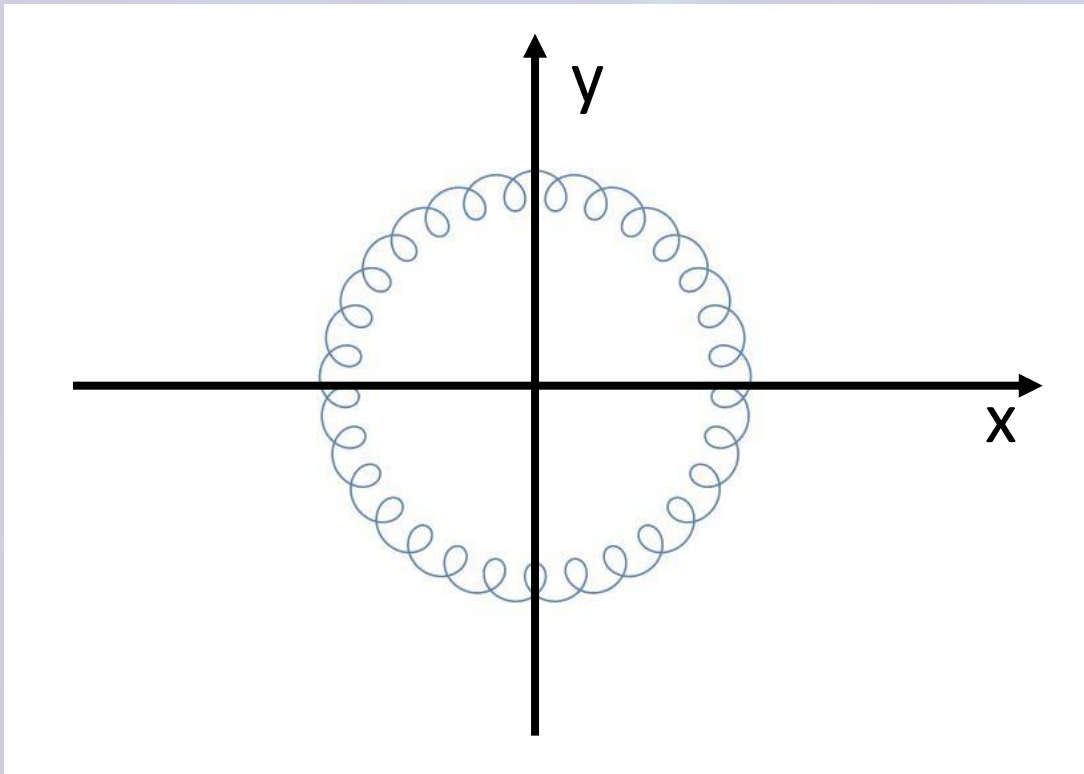
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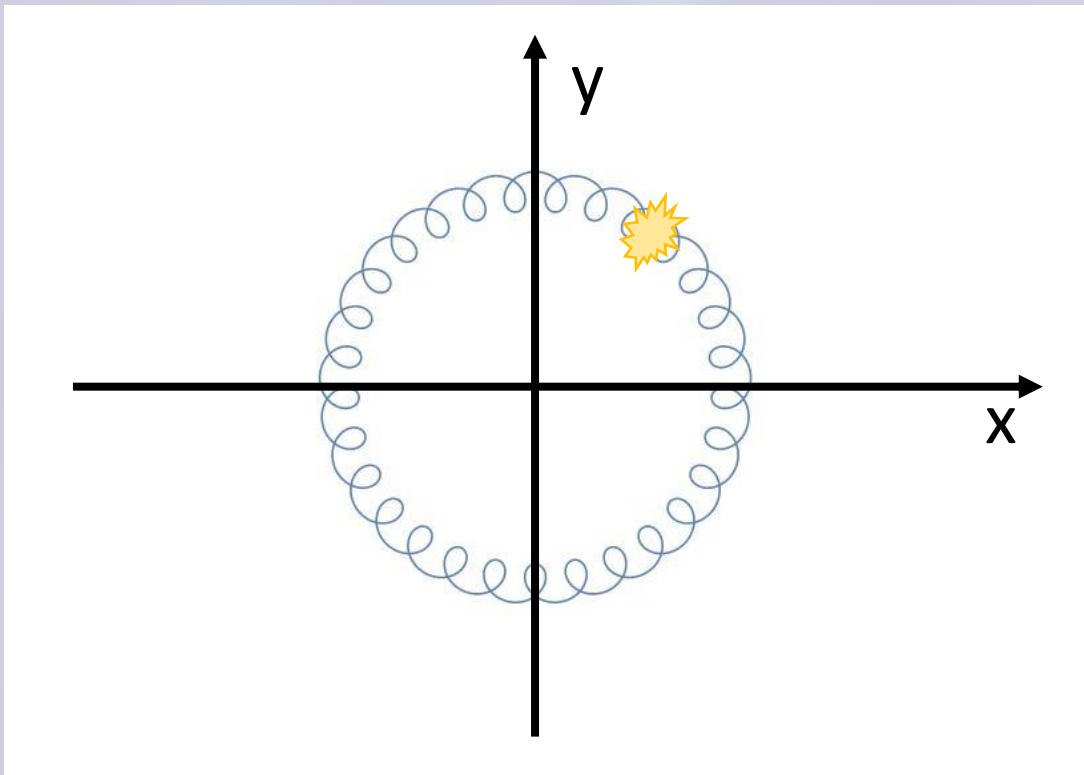
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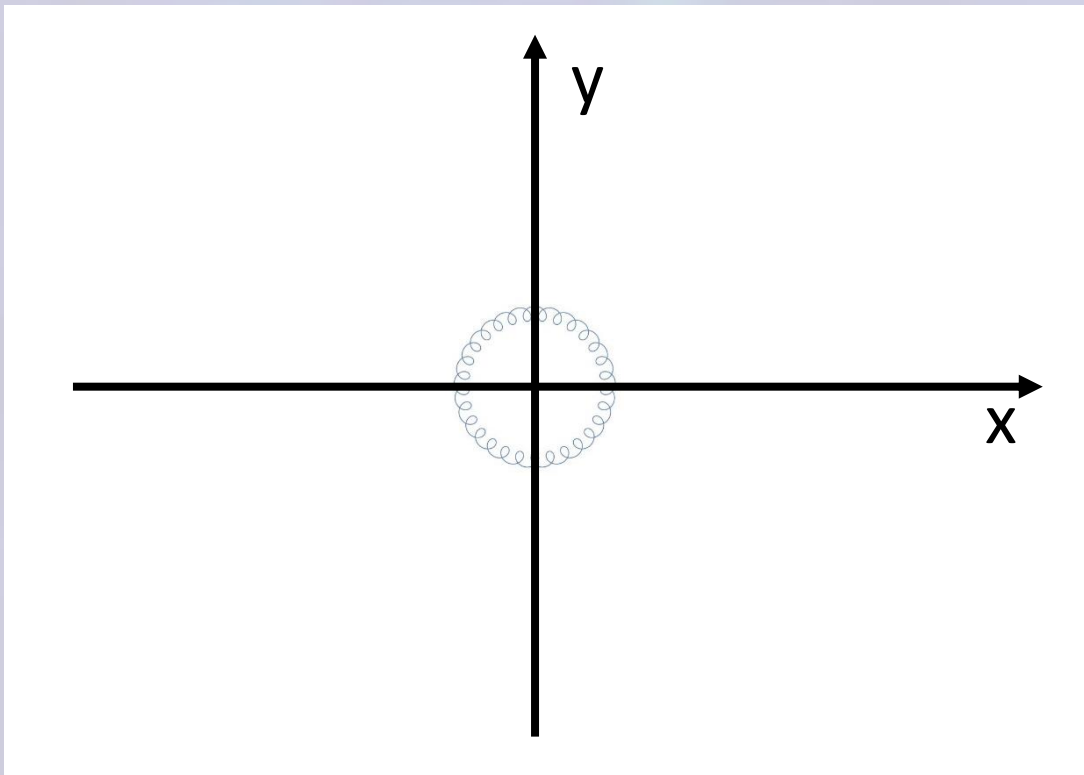
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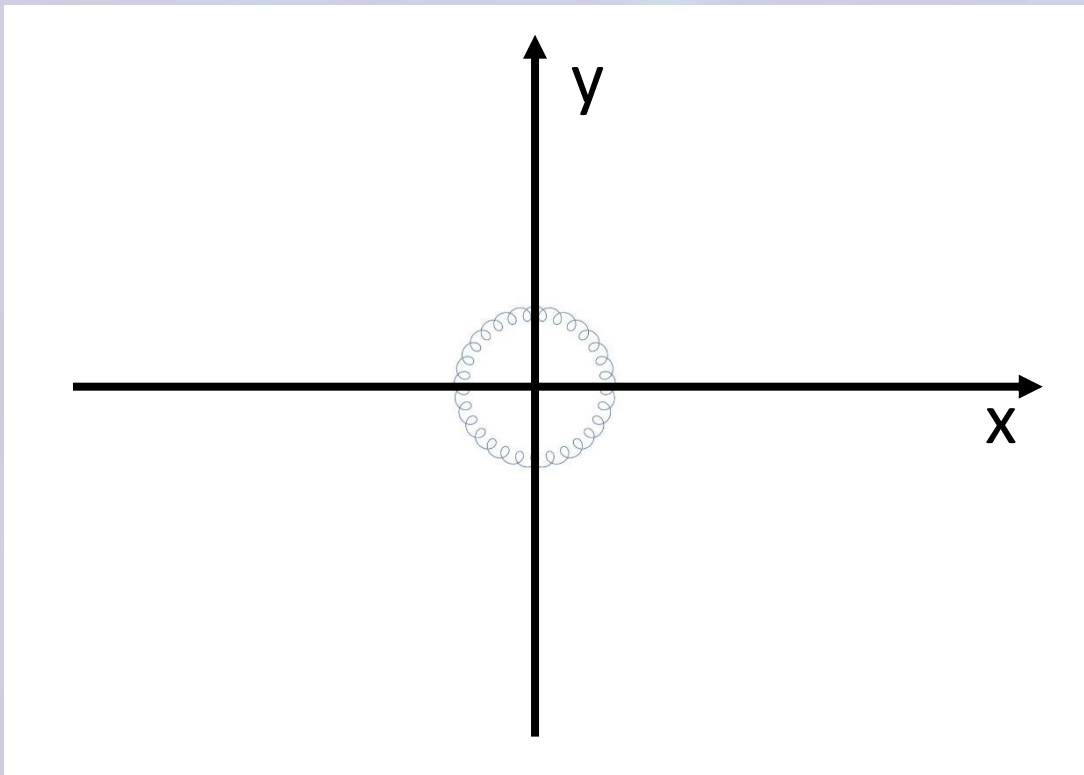
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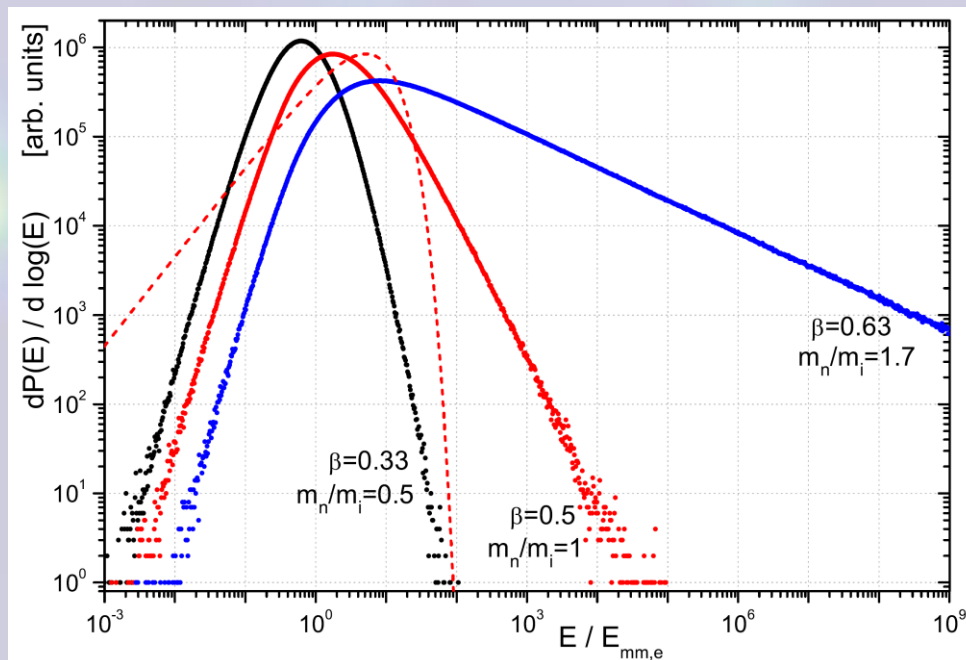
The energy distribution of an ion in a buffer gas is *not* a thermal distribution



It's like playing Texas hold'em only with all-ins! ©Roe Ozeri

Distribution energy

The energy distribution of an ion in a buffer gas is *not* a thermal distribution

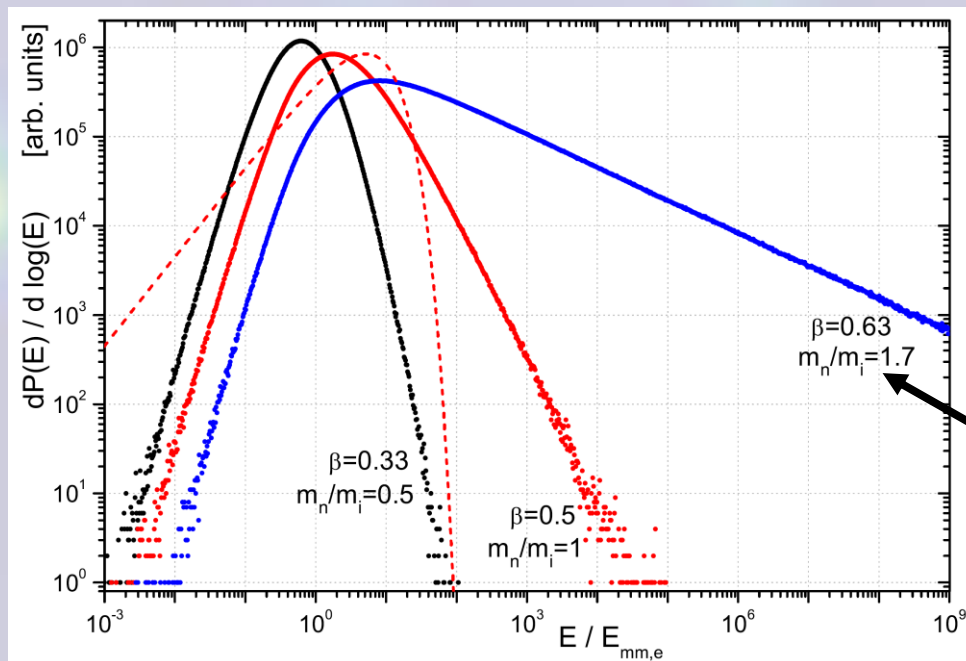


C. Zipkes et al., New J. Phys. 2011

It's a **power law distribution**.. This is complex physics!

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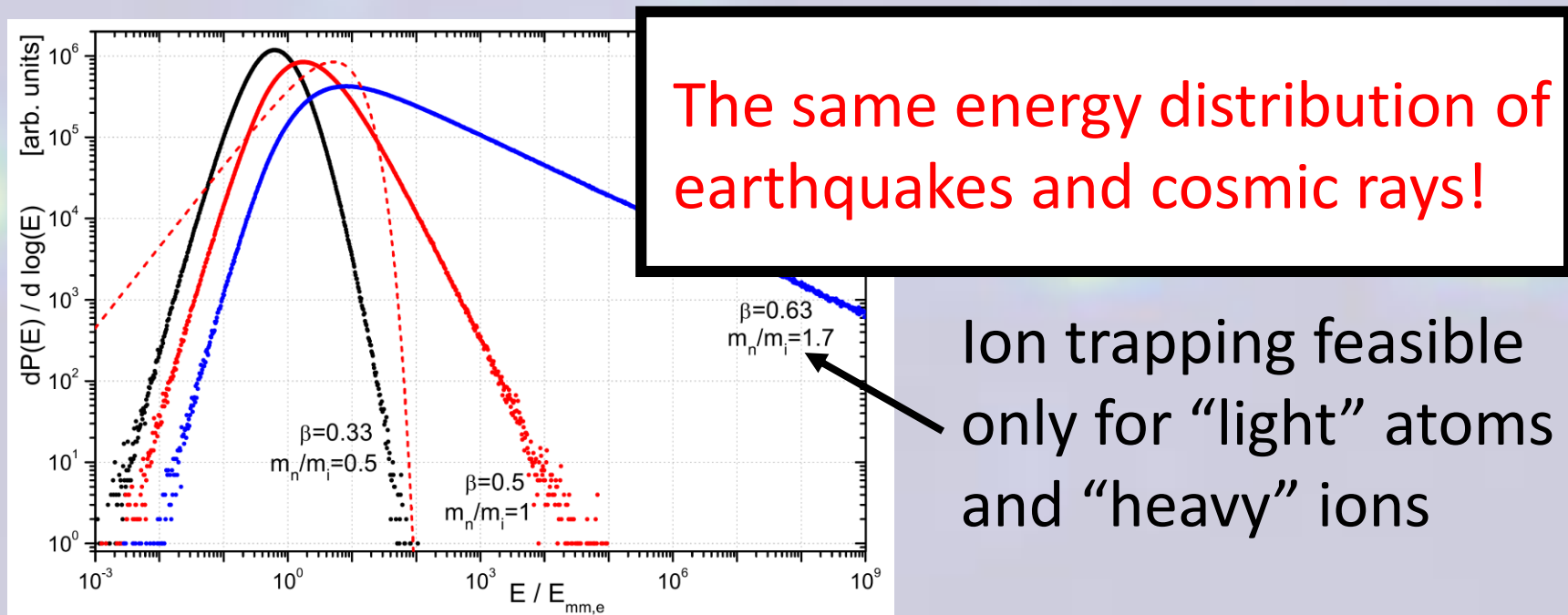
C. Zipkes et al., New J. Phys. 2011

Ion trapping feasible only for “light” atoms and “heavy” ions

It's a **power law distribution**.. This is complex physics!

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The energy distribution of an ion in a buffer gas is *not* a thermal distribution



C. Zipkes et al., New J. Phys. 2011

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Outline

A brief intro of atom-ion collisions

Hybrid atom-ion experiments – state-of-the-art

Hybrid atom-ion experiments: perspectives

Atom-ion systems for improving ion clocks

Still many issues:

- Can we enter a full quantum regime (few partial waves) in atom-ion interactions?
- Can we reach a coherent evolution of the ion-neutral mixture?
- Can we control (in a more deterministic way) inelastic processes?
-

Reducing micromotion

Atoms are an exceptional tool to measure micromotion!

Method	$\Delta\epsilon$ (V/m)
Photon-correlation spectroscopy	0.9
Micromotional sideband spectroscopy	7 / 1 / 0.4
Ion-cavity emission spectroscopy	1.8
Parametric excitation of secular motion	6 / 0.4
Neutral atom loss	0.02
Monitor displacement	≤ 11.8
Trajectory analysis	0.09

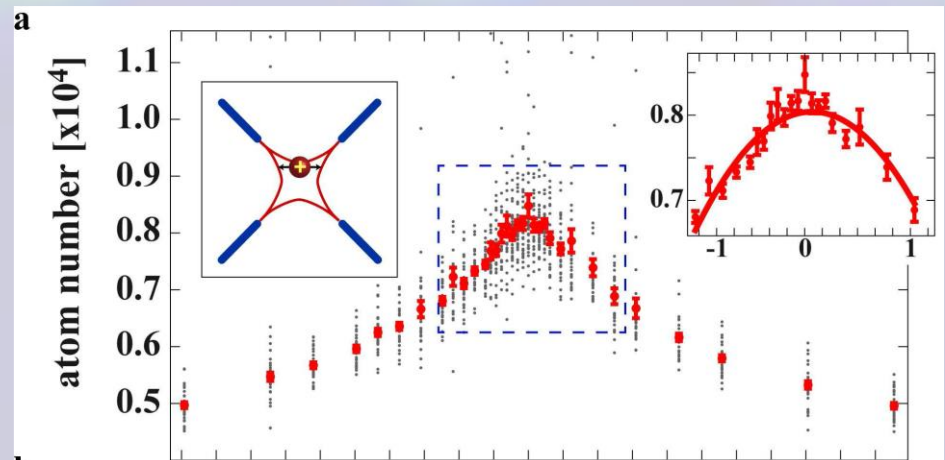
Gloger et al. PRA **92**, 043421 (2015)

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Haerter et al. APL **102**, 221115 (2013)

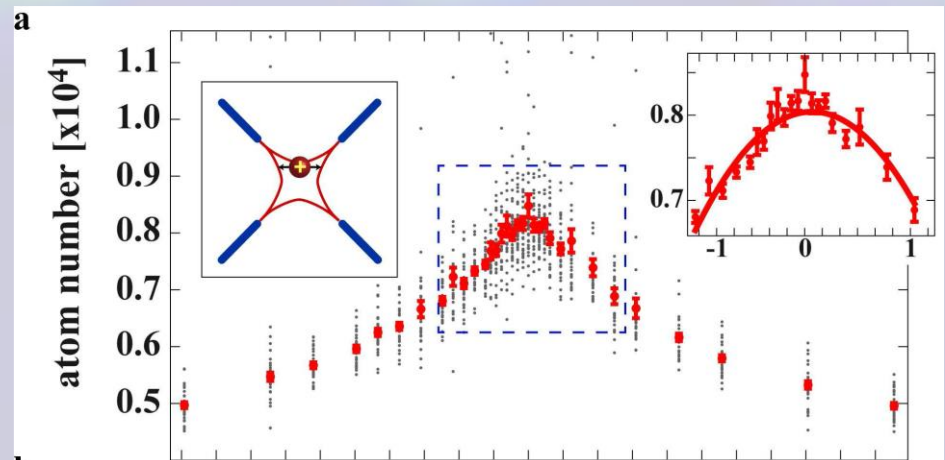
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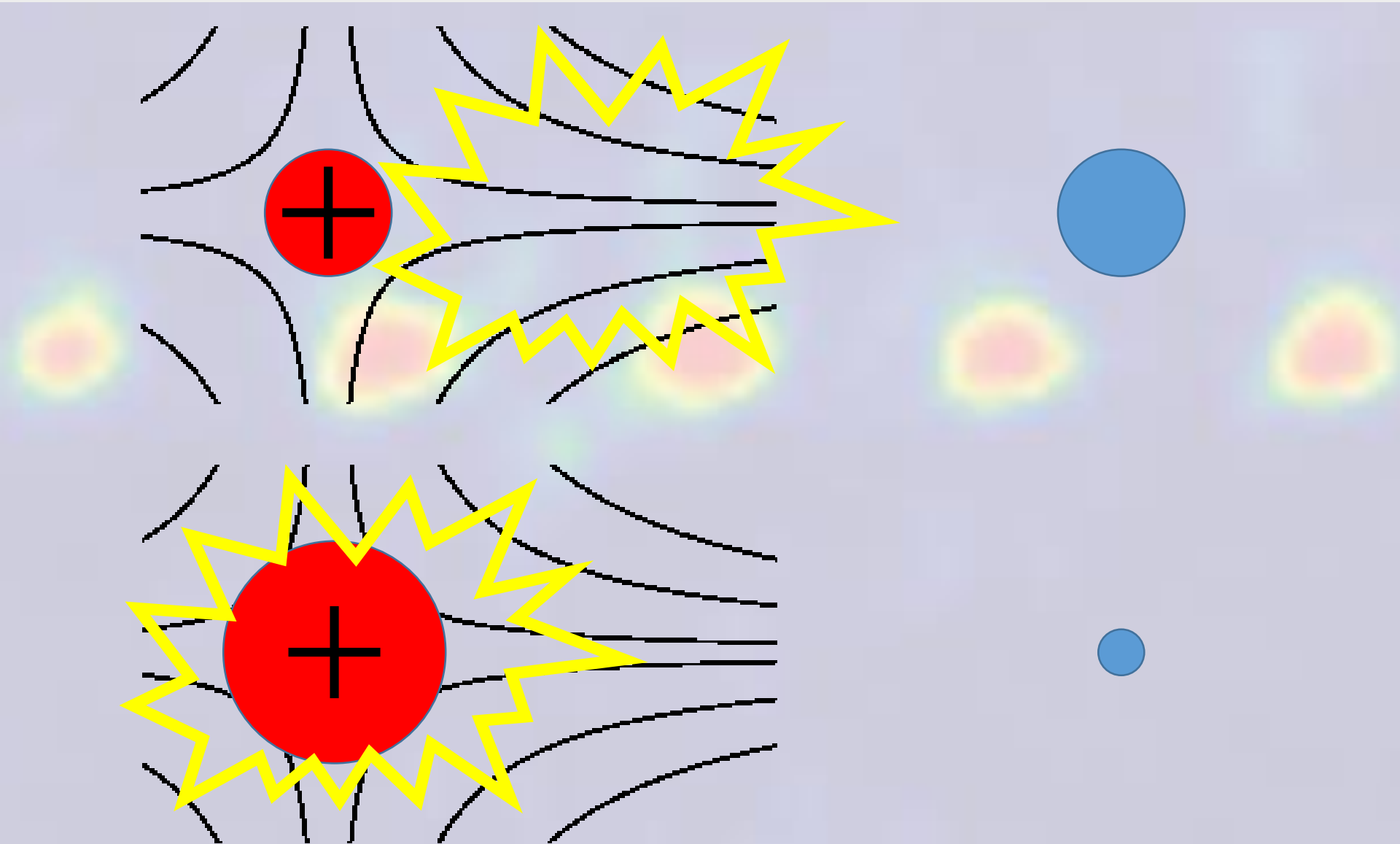
Gloger et al. PRA **92**, 043421 (2015)

Is that sufficient?



Haerter et al. APL **102**, 221115 (2013)

Reducing micromotion



Reducing micromotion

What mass ratio to reach s-wave scattering?

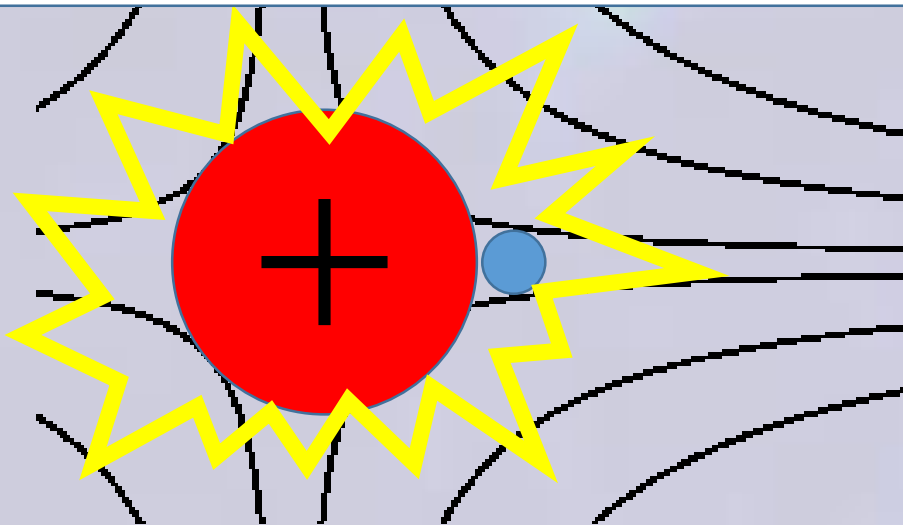
Cetina et al. PRL **109**, 253201 (2012)

→ Yb⁺ - Li

Krych et al. PRA **91**, 023430 (2015)

→ Yb⁺ - Li

→ Ba⁺ - Li



Static ion trapping

Paul trap \rightarrow dipolar trap

Optical trapping of an ion

Ch. Schneider, M. Enderlein, T. Huber and T. Schaetz*

First demonstrated in 2010



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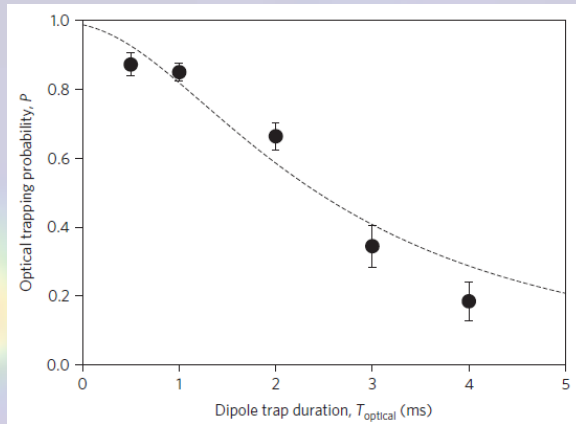
PROS:

Intrinsically without micromotion

No limitations in mass ratio

Static ion trapping

Paul trap \rightarrow dipolar trap



CONS:

SHORT lifetime because of difficulties in laser cooling

BUT atoms can provide cooling!

Optical trapping of an ion

Ch. Schneider, M. Enderlein, T. Huber and T. Schaetz*

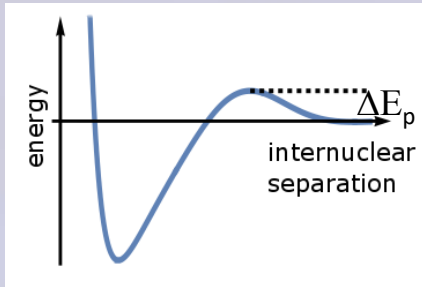
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Control of ion-neutral interactions



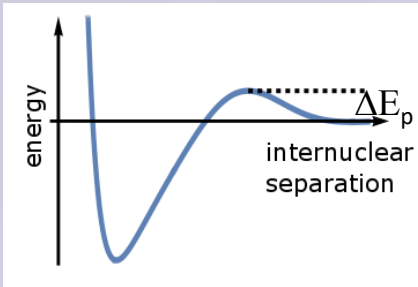
S-wave regime - collisions
described by a single phase shift:
scattering length

Elastic collisions



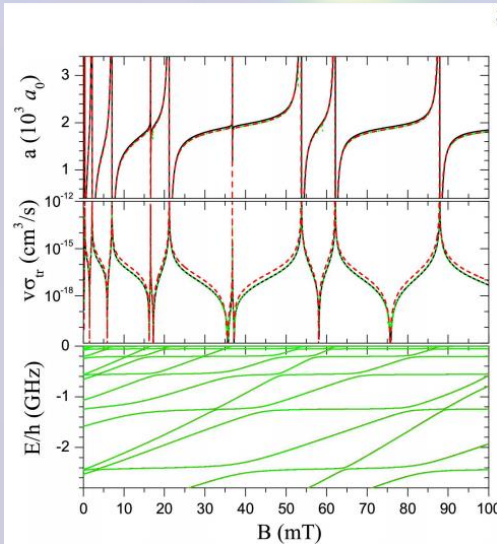
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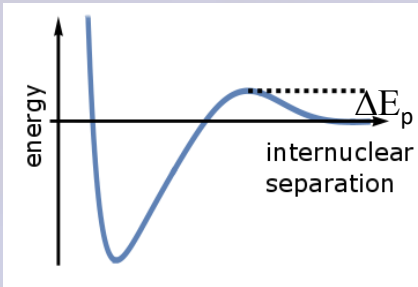
S-wave regime - collisions described by a single phase shift: scattering length

Observe Feshbach resonances to control ion-neutral interactions



Control of ion-neutral interactions

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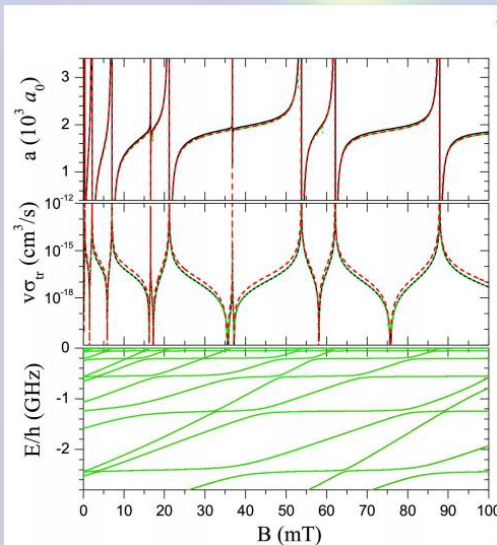


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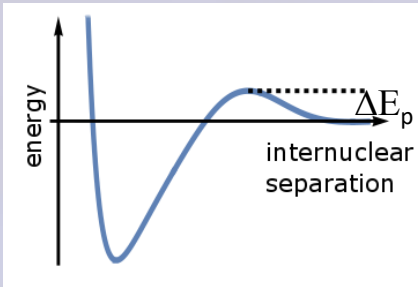
Some advice:

- Use atom-ion pair with “tall” ($>10\mu\text{K}$) p-wave barrier



Control of ion-neutral interactions

Elastic collisions

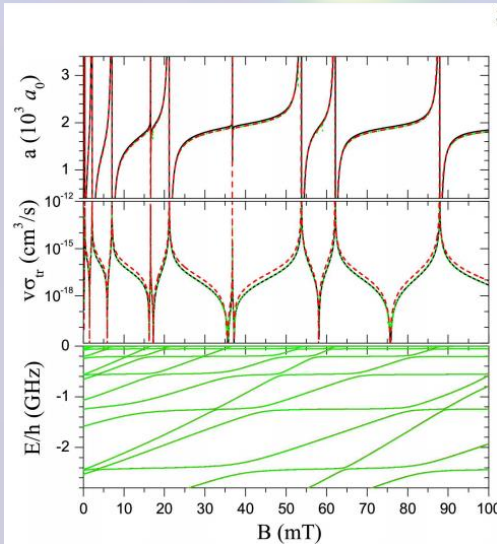


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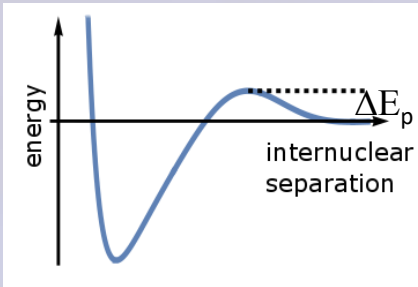
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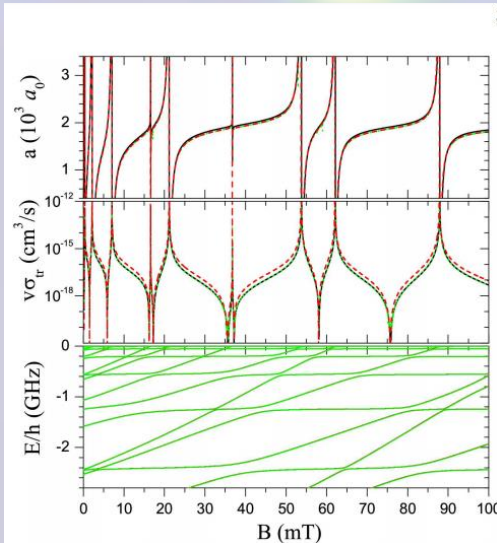
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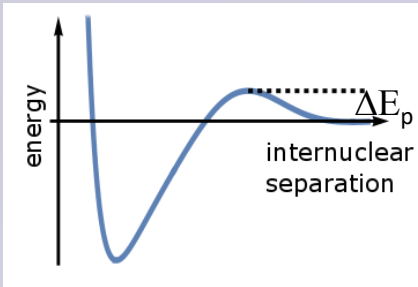


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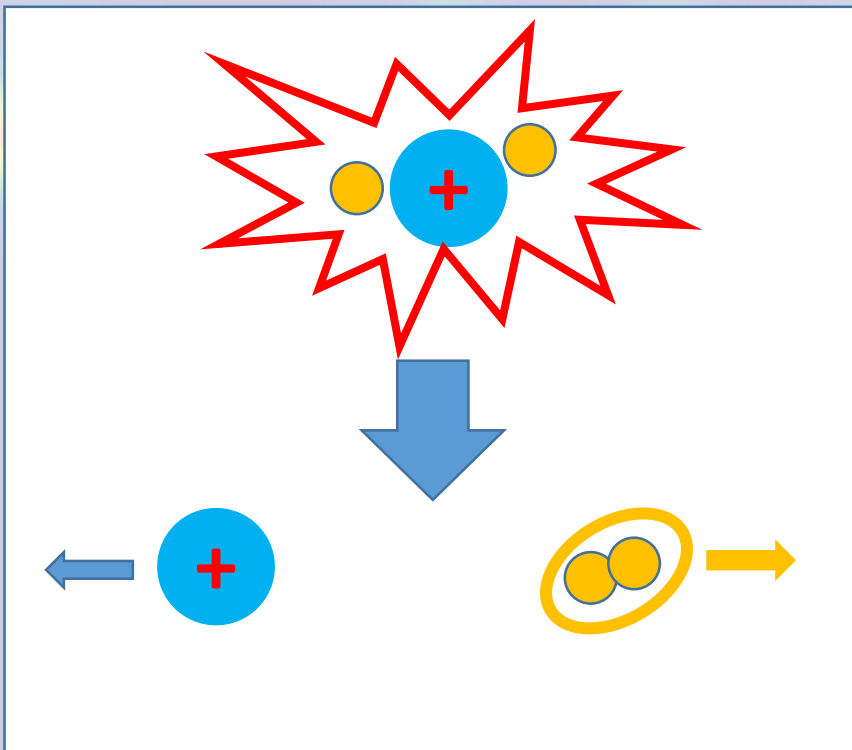
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Control of ion-neutral interactions

Elastic collisions



S-wave regime - collisions described by a single phase shift: scattering length



Some advice:

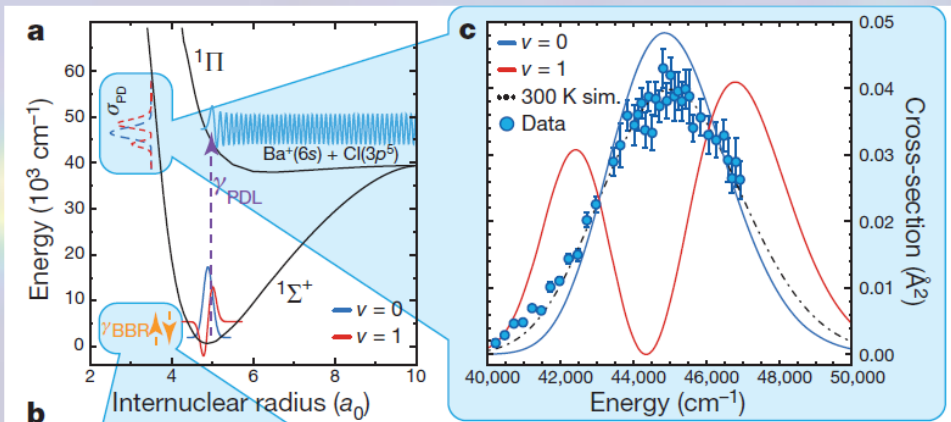
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See works by Denschlag group

Control of ion-neutral interactions

Inelastic collisions

Control the formation of molecules
and achieve sympathetic cooling



W.G. Rellergert et al. Nature 2013

Photodissociation thermometry of a
 BaCl^+ ions immersed in a Ca MOT

Also Heidelberg, Basel, Bangalore, ...

More pieces of advice:

- Use atom-ion pairs for which you know you get a molecule
- Alternatively: begin from molecules!

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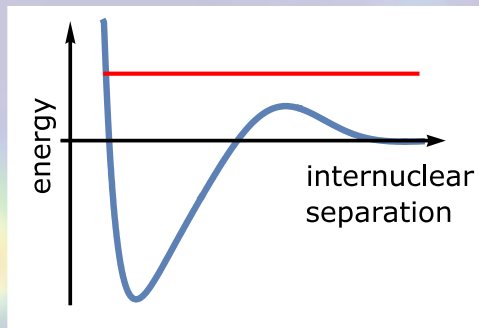
Effects of collisions on a clock

Effects: Decoherence (Langevin, soft collisions)
Quench (Langevin collisions)

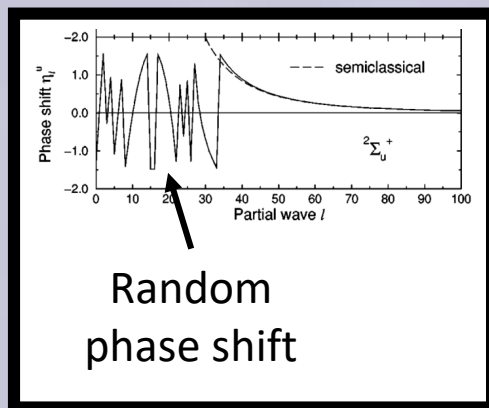


Effects of collisions on a clock

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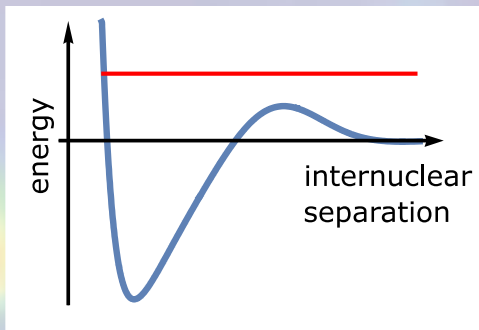


Langevin: phase shift
during clock operation

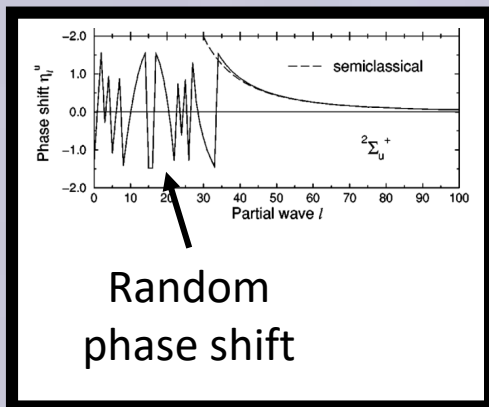


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Langevin: phase shift during clock operation



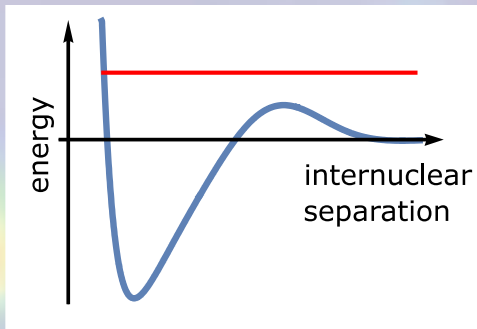
Model (Rosenband et al. Science 2008):

1. assume $\pi/2$ phase shift at each Langevin collision
 2. Integrate over optical Bloch equation
- Average frequency shift:

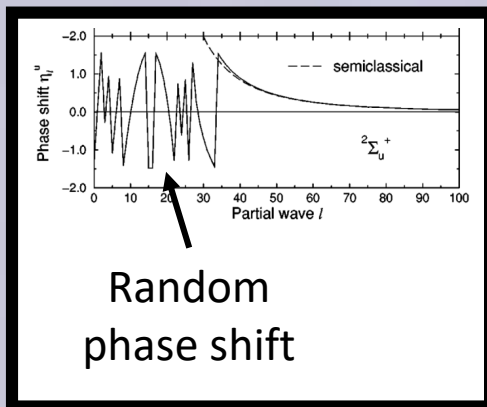
$$\Delta\nu_{\text{coll}} = 0.15 \gamma_L$$

Effects of collisions on a clock

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 Quench (Langevin collisions)



Langevin: phase shift during clock operation



Model (Rosenband et al. Science 2008):

1. assume $\pi/2$ phase shift at each Langevin collision
 2. Integrate over optical Bloch equation
- Average frequency shift:

estimated shift uncertainty

$\text{Al}^+ : 0.5 \times 10^{-18}$

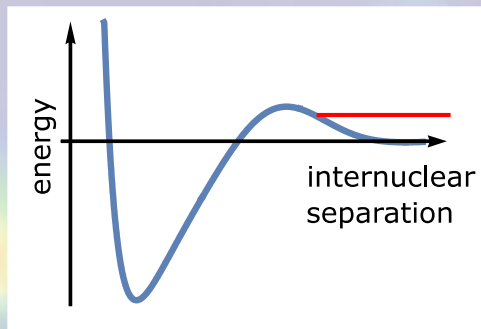
$\text{Sr}^+ : 2 \times 10^{-18}$

$$\Delta\nu_{\text{coll}} = 0.15 \gamma_L$$

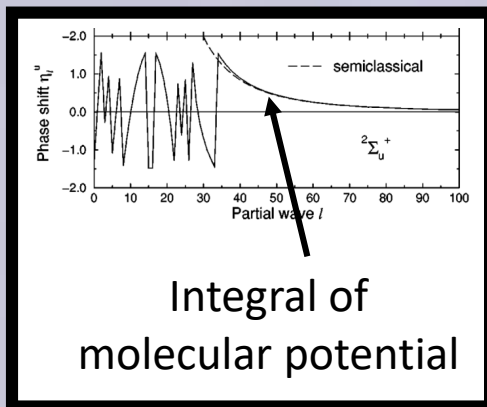
More refined model:
 Vutha et al. PRA 2017

Effects of collisions on a clock

Effects: **Decoherence** (Langevin, soft collisions)
Quench (Langevin collisions)



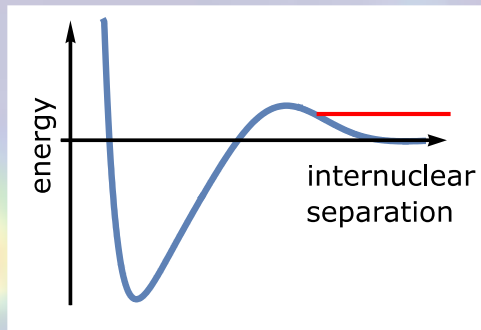
Soft collisions:
differential phase



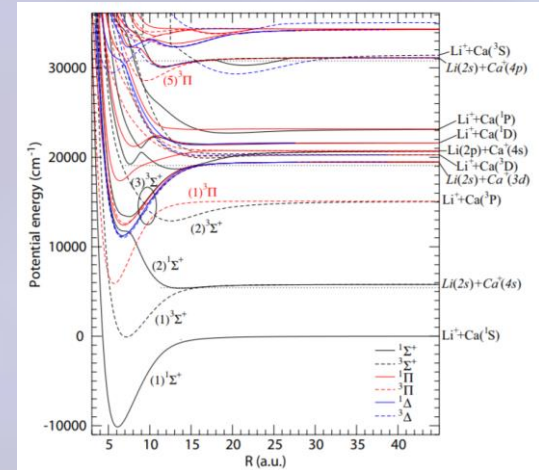
Effects of collisions on a clock

Effects: **Decoherence** (Langevin, soft collisions)

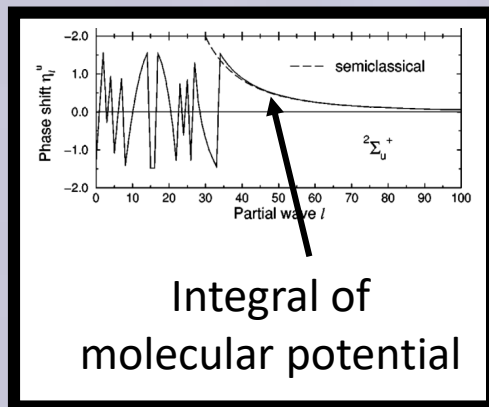
Quench (Langevin collisions)



Complicated structure!



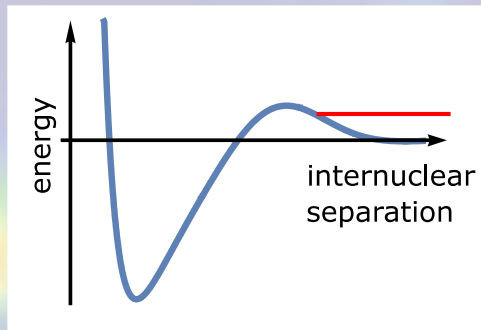
Li-Ca⁺ (Mukaiyama, Dulieu)



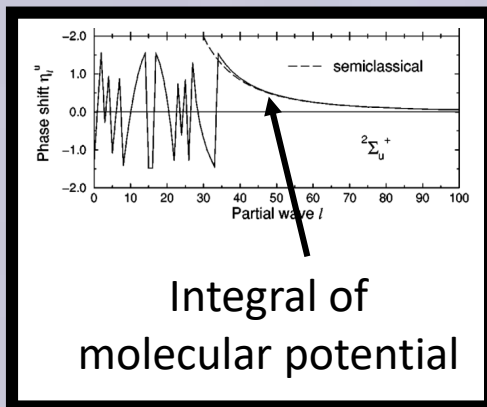
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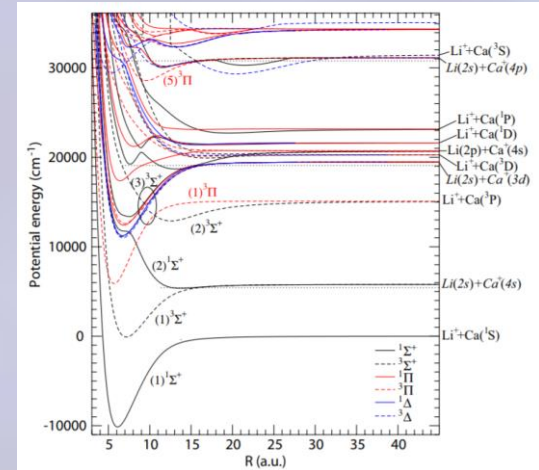
Quench (Langevin collisions)



Soft collisions:
differential phase



Complicated structure!



Li-Ca⁺ (Mukaiyama, Dulieu)

Model (Rosenband et al. Science 2008):

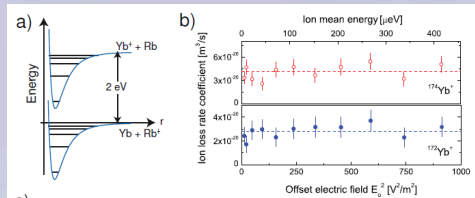
Consider the induced dipole on the atom as a fluctuating E-field

**estimated shift
uncertainty**

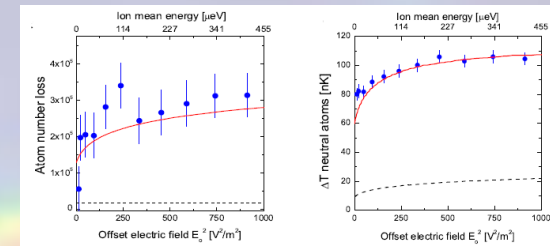
$$A\text{I}^+ < 1 \times 10^{-19}$$

Effects of collisions on a clock

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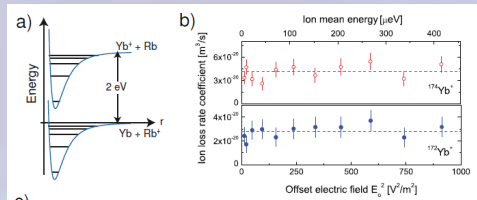


Precise detection of Langevin and soft collisions rates

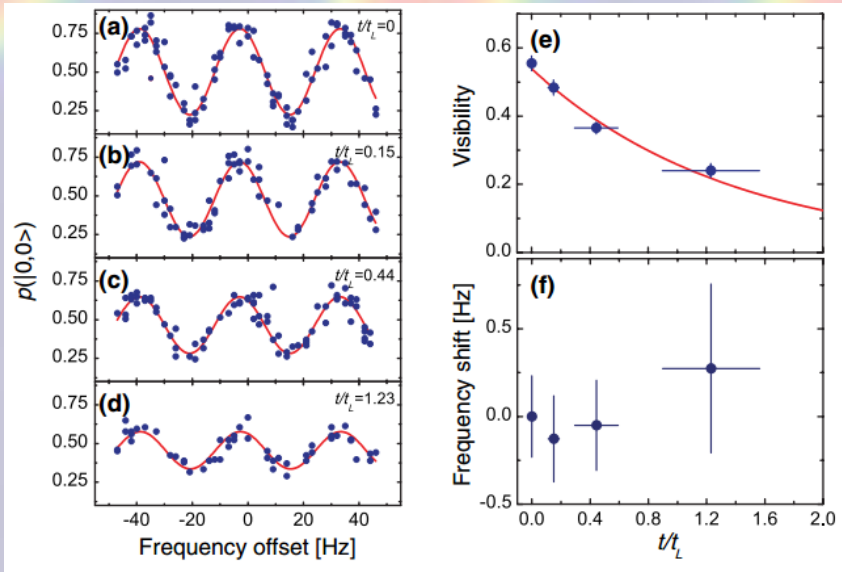
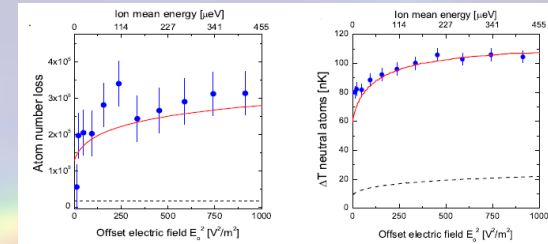


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Precise detection of Langevin and soft collisions rates



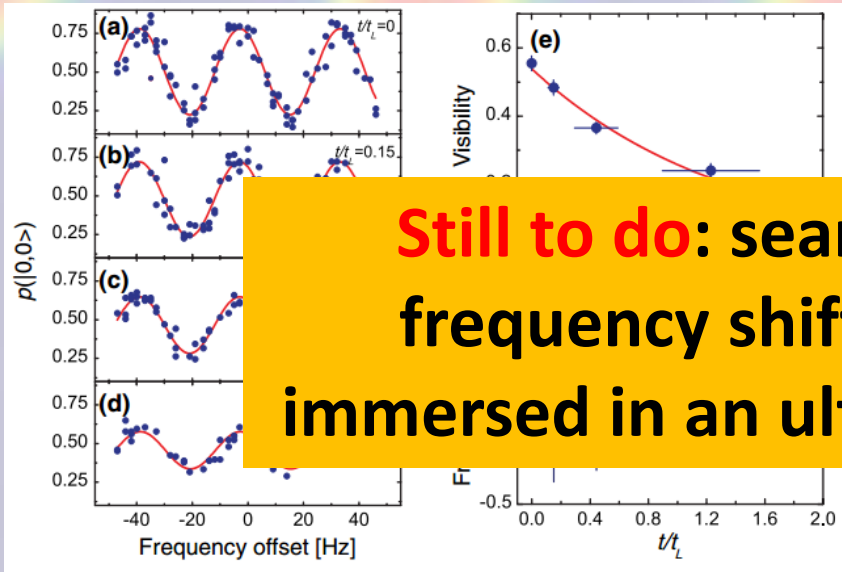
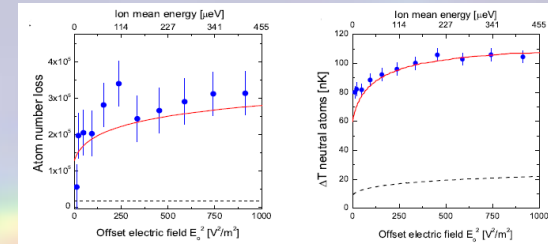
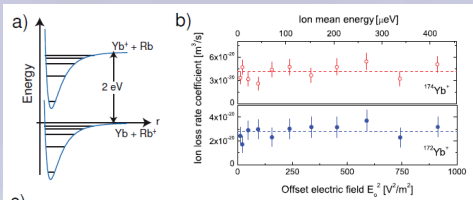
“knobs” that we can use:

- Density of the neutrals
- Collisional energy
- Optical transitions
- Internal states of the ions

Effects of collisions on a clock

Effects: **Decoherence** (Langevin, soft collisions)
 Quench (Langevin collisions)

Precise detection of Langevin
 and soft collisions rates



Still to do: search for optical frequency shifts with an ion immersed in an ultracold buffer gas

“knobs” that we can use:

- Density of the neutrals
- Energy
- Ions
- Internal states of the ions

Effects of collisions on a clock

Effects: Decoherence (Langevin, soft collisions)

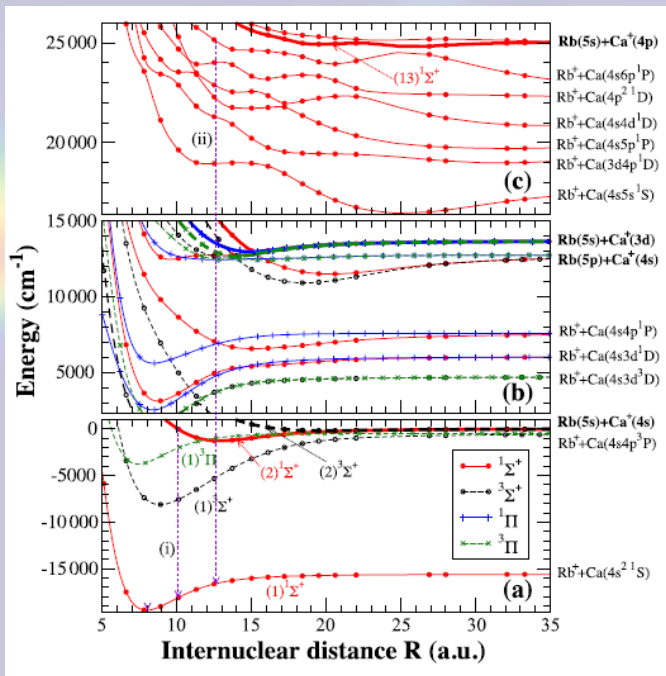
Quench (Langevin collisions)



Effects of collisions on a clock

Effects: Decoherence (Langevin, soft collisions)

Quench (Langevin collisions)



F.H.J. Hall et al. PRL 107, 243202 (2011)

Are collisionally-created molecular ions a common phenomena?

NO: Yb⁺-Rb, Ba⁺-Rb, Sr⁺-Rb, Yb⁺-Li, ...

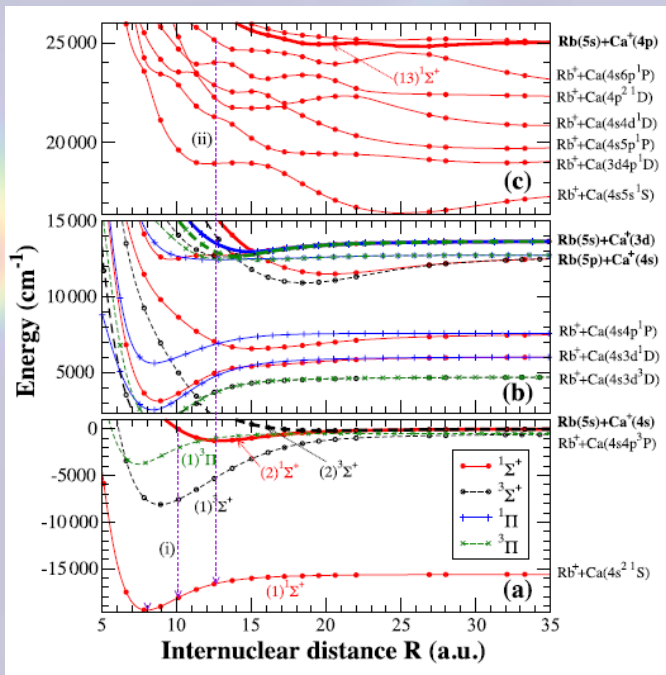
YES: Ca⁺-Rb, many hydrates (like YbH⁺)

→ A problem in HCl, since C_4 depends *quadratically* on the charge

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Questions: How big can Coulomb crystals be before being too much affected by these inelastic processes?

Can we photodissociate these molecules and recycle the ion?

Conclusions

- Neutral-ion interaction's main term is a long-range R^{-4} potential
- Collisions can be of two types: Langevin and soft collisions
- Langevin collisions' scattering rate is independent from the collisional energy, while soft collisions' scattering rate is not
- Hybrid atom-ion experiments have dramatically advanced our understanding of neutral-ion collisions
- In hybrid experiments we can control the density, collisional energy, the internal state of the colliding particles, the occurrence of chemical reactions
- Applications to ion clocks: searching for frequency shifts and for procedures to “recycle” an ion undergoing a chemical reaction

The group and funds



Elia Perego



Amelia Detti



Lucia Duca



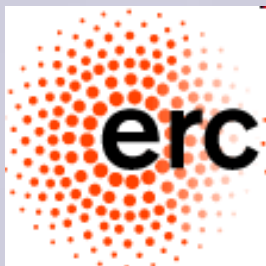
Massimo Inguscio

PhD and postdoc opportunities!

www.ultracoldplus.eu

quantumgases.lens.unifi.it

Funds:



PlusOne



Ultracoldplus



CC4C



UltraCrystals

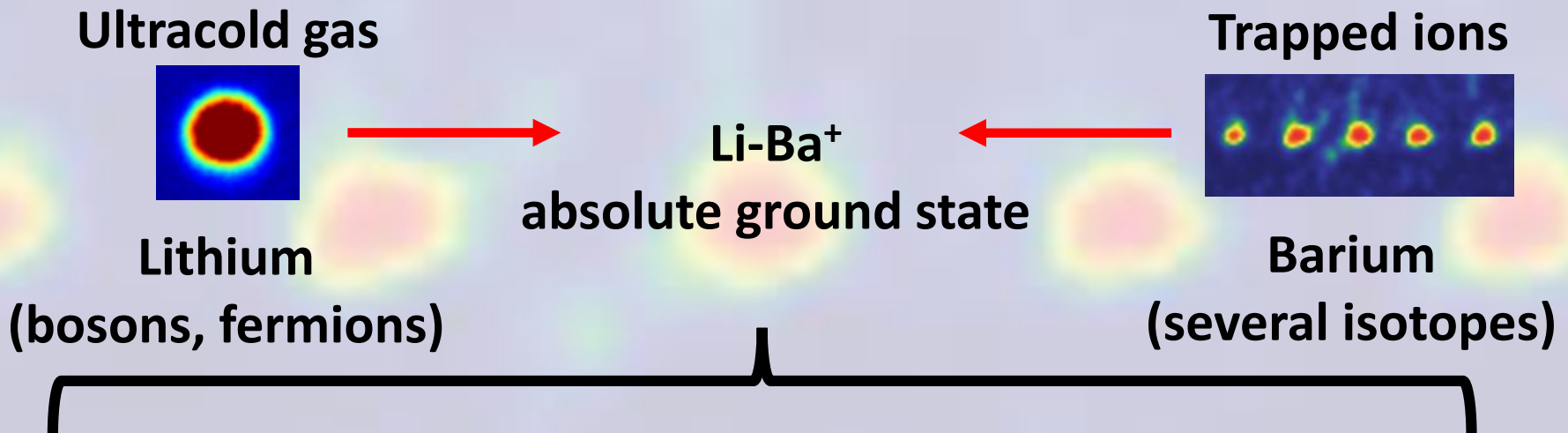
Our experiment

Our experimental setup: a novel atom-ion mixture



Our experiment

Our experimental setup: a novel atom-ion mixture



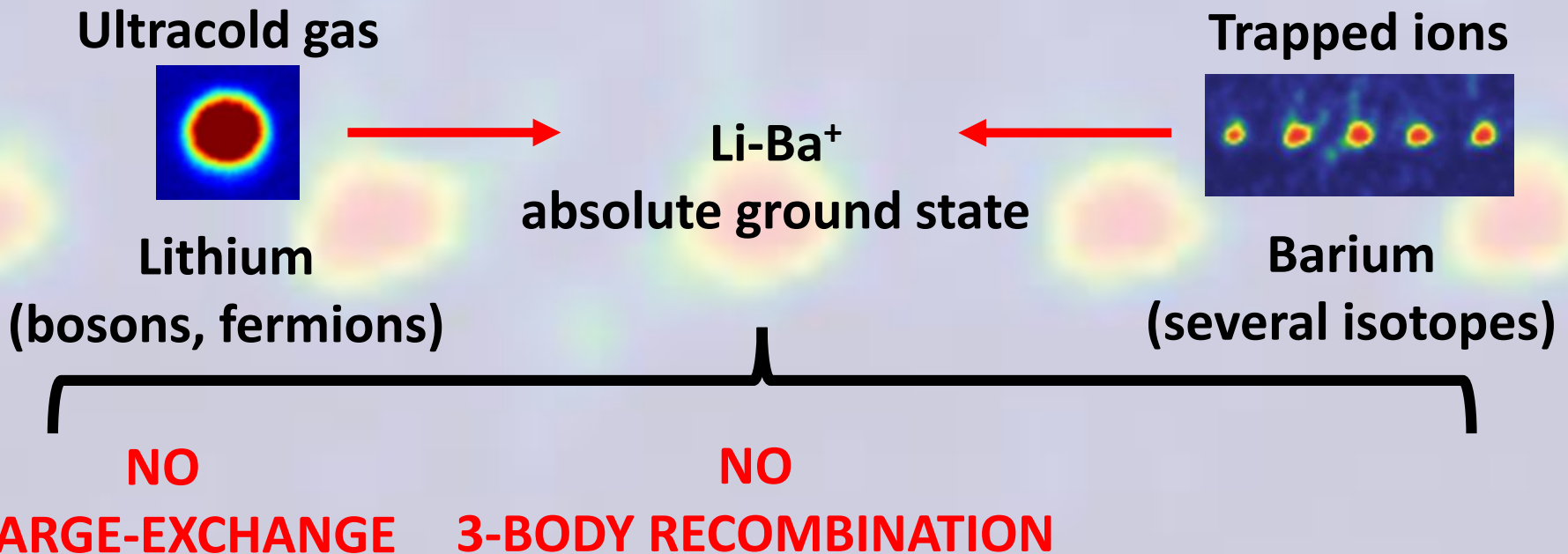
NO

CHARGE-EXCHANGE

Reactive collisions only “on demand”, i.e. only by placing the ions or the atoms in a different internal state

Our experiment

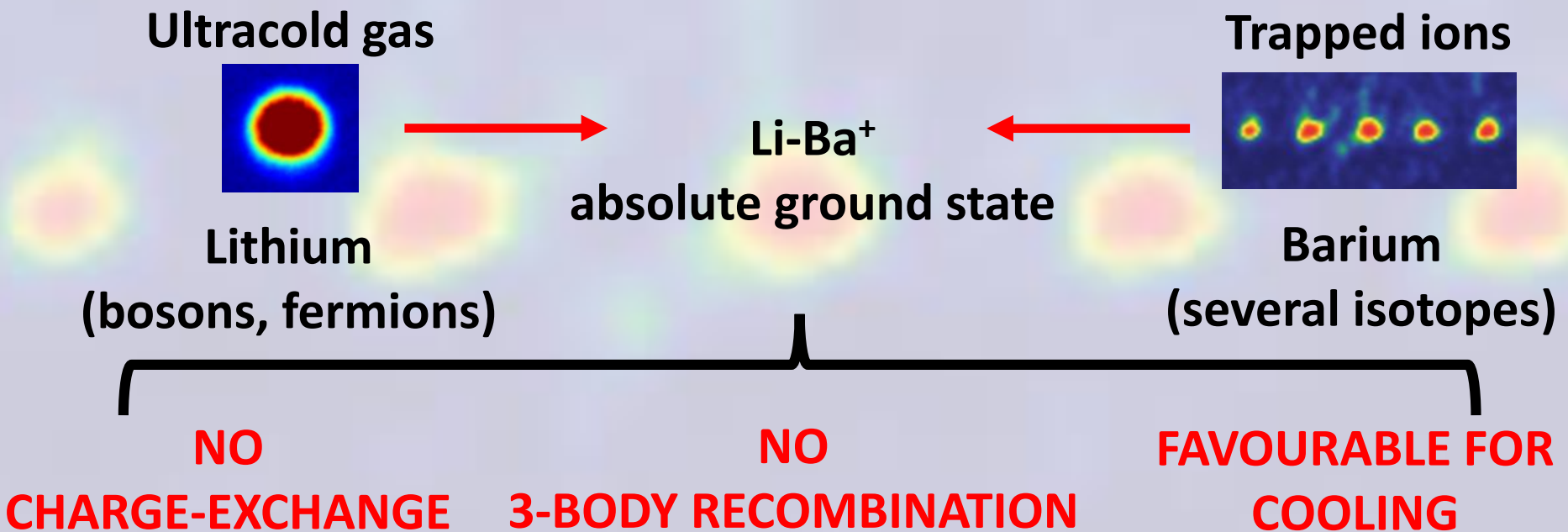
Our experimental setup: a novel atom-ion mixture



By using a spin-polarized Fermi gas, 3-body collisions (in s-wave) are suppressed by Pauli principle → less heating

Our experiment

Our experimental setup: a novel atom-ion mixture



*Large mass ratio + no spin exchange in the ground state
(Landé factor of opposite sign)*